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AFWL-TR-67-131, Vol III

AFWL-TR-  
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Vol III



## NUCLEAR EXPLOSION INTERACTION STUDIES

Volume III

The OUTPUT Code

J. R. Triplett et al.

Gulf General Atomic Incorporated  
San Diego, California 92112  
Contract No. F29601-67-C-0014

TECHNICAL REPORT NO. AFWL-TR-67-131, Vol III

April 1968

AIR FORCE WEAPONS LABORATORY  
Air Force Systems Command  
Kirtland Air Force Base  
New Mexico

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FOREWORD

This report was prepared by Gulf General Atomic Incorporated, San Diego, California, under Contract F29601-67-C-0014. The research was funded by DASA under Project 5710, Subtask 07.017, Program Element 6.16.46.01H, and by ARPA Order 313, Program Element 6.25.03.01R.

Inclusive dates of research were 29 September 1966 to 27 October 1967. The report was submitted 13 March 1968 by the Air Force Weapons Laboratory Project Officer, Major John Bode (WLRT).

This report is published in four volumes: Volume I, Laser Phenomenology (classified CONFIDENTIAL); Volume II, Two-Dimensional Code Development; Volume III, The OUTPUT Code; and Volume IV, Material Property Codes. The first volume contains a classified report on interaction of laser radiation with solid targets and a brief description of calculations done in conjunction with experiments at the Air Force Weapons Laboratory. The remaining three volumes contain reports of code development efforts in the areas of radiative transfer, hydrodynamics, radiative absorption coefficients, and equations of state.

The projects described in this report are for the most part in an incomplete state of development. This is due in part to the nature of the existing computer programs themselves, which continue in a state of development as long as they are in use, and in part to the time scale involved in bringing new programs to a state of capability for solving real problems.

Gulf General Atomic staff personnel responsible for the direction of the research include J. H. Alexander, R. Brightman, R. S. Englemore, B. E. Freeman, W. B. Lindley, L. Norris, J. T. Palmer, L. M. Schalit, J. R. Triplett, and Mrs. Chris Imes. Contractor's report number is GA-7764, Vol III.

The cooperation of Dr. P. V. Avizonis, Major J. Bode, Capt C. C. David, Major G. Spillman, and Lt L. Stoessel of AFWL is gratefully acknowledged.

Other documents produced under this contract are: GAMD-7592, "A Numerical Scheme for First-Order Compton Scattering," J. T. Palmer, December 13, 1966; GAMD-7846, "Difference Equations for Heat Flow in Two Dimensions," J. R. Triplett, March 2, 1967; GAMD-7879, "A Modified Characteristic Method for Radiative Transfer," J. R. Triplett, March 17, 1967; GAMD-7889, "R D C D. A FORTRAN Input Routine," J. H. Alexander, March 24, 1967; GAMD-8333, "Hydrodynamic Equations

for Multidimensional Problems," J. R. Triplett, October 24, 1967; GAMD-8379,  
"A Brief Study of the Thermodynamic Properties of Several Low Z Elements at Low  
Temperature," L. M. Schalit, November 22, 1967.

This technical report has been reviewed and is approved.

*John Bode*

JOHN BODE

Major, USAF

Project Officer

*Truman L Franklin*

TRUMAN L. FRANKLIN

Colonel, USAF

Chief, Theoretical Branch

*Claude K. Stambaugh*

CLAUDE K. STAMBAUGH

Colonel, USAF

Chief, Research Division

ABSTRACT

(Distribution Limitation Statement No. 2)

The OUTPUT code is designed for the analysis of early-time nuclear explosions. The equations for radiative transfer (characteristic method) and conservation of total (fluid and radiation) momentum and energy are solved in one-dimensional (plane or spherical) geometry. The radiation equations include first-order Compton scattering, and the hydrodynamic equations are treated in explicit Lagrangian form. The code is undergoing continuing development; the formulation, flow charts, glossary, and listings presented represent its status as of 27 October 1967.

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## SECTION I

### INTRODUCTION

#### PURPOSE OF CODE

The OUTPUT code, a version of the SPUTTER code, is a one-dimensional, Lagrangian, radiative hydrodynamics program written in the FORTRAN IV language. The purpose of this code is to calculate and to predict the radiative spectrum and the early-time evolution of nuclear devices. The code has been applied to several devices; the results of these applications are reported in reference 1.

#### PHYSICAL MODEL

The solution to the radiation transport equation is similar to that in the standard SPUTTER program (Ref. 2) in that the characteristic ray approach is employed for grey or multifrequency problems with mixed diffusion and transport regimes. However, OUTPUT includes as options Thomson scattering (conservative scattering) and first-order Compton scattering (nonconservative). The physics of the hydrodynamics routine have also been improved by including the radiation pressure as a tensor. This improvement has the effect of adding a term to the momentum equation and a corresponding work term to the energy equation. The energy equation also now uses the radiation pressure as calculated in the radiation routines rather than the usual equilibrium diffusion assumption  $P_r = (1/3) a \theta^4$ , where  $\theta = kT$  is the local material temperature multiplied by the Boltzmann constant, in electron volts, and  $a = 137 \text{ ergs/cm}^3/\text{eV}^4$  is the Stefan radiation constant in appropriate units.

### LOGIC OF CONSTRUCTION

The logic followed in constructing OUTPUT has been that the SPUTTER format and bookkeeping subroutines should be employed as fully as possible so as to achieve the following benefits: (1) opacities and material properties developed for SPUTTER would be immediately usable; (2) problems run with SPUTTER could be picked up (i. e., off a tape) and continued with OUTPUT or vice versa, using at most a simple transcription program; (3) the job of redeveloping the auxiliary subroutines could be avoided; and (4) personnel familiar with SPUTTER and its derivatives could easily learn the differences between the codes. The differences between the current SPUTTER program and the OUTPUT code are described in Appendix I. It should be noted that the OUTPUT code was designed for a particular class of problems; consequently, the capability of the SPUTTER program to handle certain problems, e. g., conduction, boiling, etc., has been eliminated. This has increased the efficiency of the code and released needed core storage and common variables. The user is cautioned that variables used solely in deleted sections of the SPUTTER code are now employed differently.

This report is not a complete documentation of the OUTPUT code but rather a documentation of those portions of the SPUTTER code that have been substantially altered. The SPUTTER code has been documented in reference 3. As an aid to gaining experience in using the OUTPUT code, a list of the cards used in making the successful comparison between the Gulf General Atomic and Air Force Weapons Laboratory versions of the OUTPUT code is included in Appendix II.

## SECTION II

### THEORETICAL MODEL

The equations presented in this section represent the theoretical model employed in the OUTPUT code. In general, the constraints stated with the equations limit the applicability of the model to low-temperature devices, with one-dimensional (plane or spherical) symmetry.

### RADIATION EQUATIONS

The derivation of the transport equation for the intensity  $I(\nu, \vec{\Omega}, \mathbf{r}, t)$ , i. e., the radiant energy per unit frequency  $\nu$ , per unit solid angle  $d\Omega$ , about the direction (unit vector)  $\vec{\Omega}$  per unit time  $dt$ , per unit area, including the effects of Compton scattering, is presented in reference 4. A summary of this derivation is presented below.

The contributions to the rate of change of the intensity, or to  $1/c (\partial I / \partial t) + \vec{\Omega} \cdot \nabla I$ , are assumed to be absorption, emission, and the scattering of photons by free electrons. The absorption coefficient includes all significant processes by which photons are absorbed subject to the local-thermodynamic-equilibrium (LTE) assumption. The scattering coefficient and the differential cross section for scattering are determined from the Klein-Nishina formula (Ref. 5). A series of further approximations are required to complete the hypothesis. They are:

1. The electron states before and after scattering are nondegenerate.
2. Doppler effects can be ignored.
3. Polarization is unimportant.
4.  $h\nu/m_0 c^2 \leq 0.2$ ,  $m_0 c^2$  = rest energy of an electron.
5. Stimulated scattering is negligible.

6. Retardation effects are omitted.

7. The spectrum is a smooth function of  $\nu$ .

With these assumptions, the following transport equation in plane geometry is obtained for the quantity  $I(\mu) \equiv \frac{2\pi}{c} I(\vec{\Omega})$ :

$$\mu \frac{dI}{dx} = \mu'_a (B_\nu - I) + S \quad (1)$$

where  $\mu = \vec{\Omega} \cdot \hat{x}$  is the normal direction cosine;

$$\mu'_a = \mu_a (1 - e^{-hv/\theta}) \quad (2)$$

is the linear absorption coefficient corrected for induced emission;

$$B_\nu = \frac{4\pi h\nu^3}{c^3} (e^{hv/\theta} - 1)^{-1} \quad (3)$$

is the Planck distribution function for radiation from a blackbody at temperature  $\theta$  eV, multiplied by  $\frac{2\pi}{c}$ ;

$$S = -\mu_s \left[ I - \frac{3}{16} \int_{-1}^1 d\mu' I(\mu') [3 - \mu'^2 + (3\mu'^2 - 1)(\mu')^2] \right. \\ \left. + \gamma \left\{ -2I + \frac{3}{16} \int_{-1}^1 d\mu' \left( I(\mu') - \gamma \frac{dI(\mu')}{d\gamma} \right) [3 - \mu'^2 + \mu'\mu(3\mu'^2 - 5) \right. \right. \\ \left. \left. + (3\mu'^2 - 1)(\mu')^2 + \mu(3 - 5\mu'^2)(\mu')^3] \right\} \right] + O(\gamma^2) \quad (4)$$

where

$$\mu_s = \frac{8}{3} \pi r_o^2 N_e \quad (5)$$

is the Thomson scattering coefficient valid for  $hv \ll m_o c^2$ , in which  $r_o = e^2/(m_o c^2)$  is the classical electron radius and  $N_e$  is the number of electrons per  $\text{cm}^3$ ; and

$$\gamma = \frac{hv}{m_o c^2} \quad (6)$$

Examination of the form of the scattering source term reveals that the low-energy limit, i.e.,  $\gamma \rightarrow 0$ , gives the Thomson scattering law. To determine the material heating due to Compton scattering, it is necessary to integrate over all solid angles and over the entire frequency range. The result, with sign reversed, is

$$-\nabla \cdot \vec{F} = -c \int_0^\infty \mu_a' (2B_\nu - E_\nu) d\nu + \mu_s \frac{hc}{m_e c^2} \int_0^\infty \nu E_\nu d\nu \quad (7)$$

where

$$\vec{F} = \int_0^\infty d\nu \int_{4\pi} I(\nu, \vec{\Omega}) \vec{\Omega} d\Omega$$

is the total flux and

$$E_\nu = \frac{1}{c} \int_{4\pi} I(\nu, \vec{\Omega}) d\Omega$$

is the radiant energy density spectrum. The first term on the right is the usual heating term due to absorption or emission, whereas the second term is the heating rate due to Compton scattering.

#### NUMERICAL INTEGRATION OF THE TRANSPORT EQUATION

The development of a scheme to integrate the transport equation along a ray requires that several numerical approximations be made. Rewrite the source term  $S$  by rearranging and by employing the notation

$$I_n = \int_{-1}^1 I(\mu) (\mu)^n d\mu \quad (8)$$

Thus,  $I_0$ ,  $I_1$ , and  $I_2$  are proportional to the radiation energy density, the flux, and the radiation pressure, respectively. Now,

$$\begin{aligned}
 S = & -\mu_s \left\{ I(1 - 2\gamma) - \frac{3}{16} [(3 - \mu^2) I_0 + (3\mu^2 - 1) I_2] \right. \\
 & + \frac{3}{16} \gamma [(3 - \mu^2) I_0 + \mu(3\mu^2 - 5) I_1 + (3\mu^2 - 1) I_2 + \mu(3 - 5\mu^2) I_3] \\
 & \left. - \frac{3}{16} \gamma^2 \frac{\partial}{\partial \gamma} [(3 - \mu^2) I_0 + \mu(3\mu^2 - 5) I_1 + (3\mu^2 - 1) I_2 + \mu(3 - 5\mu^2) I_3] \right\} \\
 & \quad (9)
 \end{aligned}$$

It is convenient to expand the last term as follows:

$$\gamma^2 \frac{\partial f(\gamma)}{\partial \gamma} = \frac{\partial [\gamma^2 f(\gamma)]}{\partial \gamma} - 2\gamma f(\gamma) \quad (10)$$

Using this substitution, one finds

$$\begin{aligned}
 S = & -\mu_s \left\{ I(1 - 2\gamma) - \frac{3}{16} [(3 - \mu^2) I_0 + (3\mu^2 - 1) I_2] \right. \\
 & + \frac{9}{16} \gamma [(3 - \mu^2) I_0 + \mu(3\mu^2 - 5) I_1 + (3\mu^2 - 1) I_2 + \mu(3 - 5\mu^2) I_3] \\
 & \left. - \frac{3}{16} \frac{\partial \gamma^2}{\partial \gamma} [(3 - \mu^2) I_0 + \mu(3\mu^2 - 5) I_1 + (3\mu^2 - 1) I_2 + \mu(3 - 5\mu^2) I_3] \right\} \\
 & \quad (11)
 \end{aligned}$$

In Eq. (1), the intensity  $I(\mu, \nu)$  is the monochromatic intensity. In most computer codes employing radiation transport, the quantity being calculated is the integral of the intensity over some frequency band  $(\nu_j, \nu_{j+1})$ , for which one finds

$$\mu \frac{dI_j}{dx} = \mu_{a_j} (B_j - I_j) + S_j \quad (12)$$

where

$$S_j = \int_{\nu_j}^{\nu_{j+1}} S(\nu) d\nu \quad (13)$$

$$I_j = \int_{\nu_j}^{\nu_{j+1}} I(\nu) d\nu \quad (14)$$

The notation  $I_j$  must not be confused with  $I_n$  as previously defined. However, the distinction between symbols is fairly clear. Note that single subscripts denote intensities and double subscripts represent moment quantities. Continuing, one finds

$$\begin{aligned}
 S_j = & -\mu_s \left\{ I_j(1 - 2\bar{\gamma}_j) - \frac{3}{16} [(3 - \mu^2) I_{o_j} + (3\mu^2 - 1) I_{2_j}] \right. \\
 & + \frac{9}{16} \bar{\gamma}_j [(3 - \mu^2) I_{o_j} + \mu(3\mu^2 - 5) I_{1_j} + (3\mu^2 - 1) I_{2_j} + \mu(3 - 5\mu^2) I_{3_j}] \\
 & - \frac{3}{16} \frac{h}{m_o c^2} v_{j+1}^2 [(3 - \mu^2) I_o(v_{j+1}) + \mu(3\mu^2 - 5) I_1(v_{j+1}) \\
 & \quad \left. + (3\mu^2 - 1) I_2(v_{j+1}) + \mu(3 - 5\mu^2) I_3(v_{j+1}) \right] \\
 & + \frac{3}{16} \frac{h}{m_o c^2} v_j^2 [(3 - \mu^2) I_o(v_j) + \mu(3\mu^2 - 5) I_1(v_j) + \mu(3 - 5\mu^2) I_3(v_j) \\
 & \quad \left. + (3\mu^2 - 1) I_2(v_j) \right] \} \tag{15}
 \end{aligned}$$

where the assumption has been made that  $\bar{\gamma}_j = (1/2)(\gamma_j + \gamma_{j+1})$  and

$$I_{x_j} = \int_{v_j}^{v_{j+1}} I_x(v) dv \tag{16}$$

Moreover,

$$\bar{\gamma}_j I_{x_j} = \int_{v_j}^{v_{j+1}} \gamma I_x(v) dv \tag{17}$$

One should note that in Eq. (15) the quantities  $I_o(v_j)$ ,  $I_1(v_j)$ , etc., are the moment quantities evaluated at the boundaries of the frequency group of interest, i.e.,  $v_j$  and  $v_{j+1}$ . At this point, an approximation must be made relating the boundary quantities and the frequency average quantities. If

the spectrum is flat over the various frequency groups, then

$$I_{x_j} \cong I_x (\nu_j)(\nu_{j+1} - \nu_j) \quad (18)$$

Thus,

$$\begin{aligned} S_j = -\mu_s & \left\{ I_j(1 - 2\bar{\gamma}_j) - \frac{3}{16} [(3 - \mu^2) I_{o_j} + (3\mu^2 - 1) I_{2_j}] \right. \\ & + \frac{9}{16} \bar{\gamma}_j [(3 - \mu^2) I_{o_j} + \mu(3\mu^2 - 5) I_{1_j} + (3\mu^2 - 1) I_{2_j} + \mu(3 - 5\mu^2) I_{3_j}] \\ & - \frac{3}{16} \frac{h}{m_o c^2} \frac{\nu_{j+1}^2}{(\nu_{j+2} - \nu_{j+1})} [(3 - \mu^2) I_{o_{j+1}} + \mu(3\mu^2 - 5) I_{1_{j+1}} + (3\mu^2 - 1) I_{2_{j+1}} \\ & \quad \left. + \mu(3 - 5\mu^2) I_{3_{j+1}}] \right\} \\ & + \frac{3}{16} \frac{h}{m_o c^2} \frac{\nu_j^2}{\nu_{j+1} - \nu_j} [(3 - \mu^2) I_{o_j} + \mu(3\mu^2 - 5) I_{1_j} + (3\mu^2 - 1) I_{2_j} \\ & \quad \left. + \mu(3 - 5\mu^2) I_{3_j}] \right\} \end{aligned} \quad (19)$$

Reorganizing the source term with  $\Delta\nu_j = \nu_{j+1} - \nu_j$ , one obtains

$$\begin{aligned} S_j = -\mu_s & \left\{ I_j(1 - 2\bar{\gamma}_j) - \frac{3}{16} A_2 [(3 - \mu^2) I_{o_j} + (3\mu^2 - 1) I_{2_j}] \right. \\ & + \frac{3}{16} A_1 \mu [(3\mu^2 - 5) I_{1_j} + (3 - 5\mu^2) I_{3_j}] \\ & - \frac{3}{16} A_3 [(3 - \mu^2) I_{o_{j+1}} + (3\mu^2 - 1) I_{2_{j+1}}] \\ & \quad \left. - \frac{3}{16} A_3 \mu [(3\mu^2 - 5) I_{1_{j+1}} + (3 - 5\mu^2) I_{3_{j+1}}] \right\} \end{aligned} \quad (20)$$

where

$$A_1 = 3\bar{\gamma}_j + \frac{h}{m_o c^2} \frac{v_j^2}{\Delta v_j} \quad (21)$$

$$A_2 = 1 - A_1, \quad (22)$$

$$A_3 = \frac{h}{m_o c^2} \frac{v_{j+1}^2}{\Delta v_{j+1}} \quad (23)$$

At this point, the integration of the transport equation can be performed as follows:

$$\mu \frac{dI_j}{dx} = \mu_{a_j} (B_j - I_j) - \mu_s I_j (1 - 2\bar{\gamma}_j) + \mu_s \{ \} \quad (24)$$

where

$$\{ \} = \frac{s_j}{\mu_s} + I_j (1 - 2\bar{\gamma}_j) \quad (25)$$

Using

$$e^{(x/\mu) [\mu_{a_j} + \mu_s (1 - 2\bar{\gamma}_j)]} \quad (26)$$

as an integrating factor, one finds

$$I_j(x_2) = I_j(x_1) e^{-\alpha_j(x_2 - x_1)} + \frac{\mu_{a_j}}{\mu} \int_{x_1}^{x_2} B_j(x') e^{-\alpha_j(x_2 - x')} dx' \\ + \frac{\mu_s}{\mu} \int_{x_1}^{x_2} \{ \} e^{-\alpha_j(x_2 - x')} dx' \quad (27)$$

where

$$\alpha_j = \frac{1}{\mu} [\mu_{a_j} + \mu_s (1 - 2\bar{\gamma}_j)] \quad (28)$$

The first two terms on the right-hand side have already been evaluated in radiation-transport routines currently employed by SPUTTER (Ref. 2). Here, the integral in the third term will be evaluated. The term in braces [Eq. (25)] contains such terms as  $I_{o_j}$ ,  $I_{o_{j+1}}$ , etc. At this point, some approximation must be made concerning the spatial dependence of these quantities. Following the assumption used in developing the Thomson scattering code, assume that all functions vary linearly in geometric space. Performing the required integration, one then finds

$$\begin{aligned}
 & \frac{\mu_s}{\mu} \int_{x_1}^{x_2} \{ \} e^{-\alpha_j(x_2 - x')} dx' \\
 &= \frac{\mu_s}{\mu} \frac{3}{16} A_2 \left[ \left( \frac{1 - e^{-\alpha_j \Delta}}{\alpha_j} \right) A + B \left( \frac{x_2 - x_1 e^{-\alpha_j \Delta}}{\alpha_j} - \frac{1 - e^{-\alpha_j \Delta}}{\alpha_j^2} \right) \right] \\
 & - \frac{\mu_s}{\mu} \frac{3}{16} \mu A_1 \left[ \left( \frac{1 - e^{-\alpha_j \Delta}}{\alpha_j} \right) C + D \left( \frac{x_2 - x_1 e^{-\alpha_j \Delta}}{\alpha_j} - \frac{1 - e^{-\alpha_j \Delta}}{\alpha_j^2} \right) \right] \\
 & + \frac{\mu_s}{\mu} \frac{3}{16} A_3 \left[ \left( \frac{1 - e^{-\alpha_j \Delta}}{\alpha_j} \right) E + F \left( \frac{x_2 - x_1 e^{-\alpha_j \Delta}}{\alpha_j} - \frac{1 - e^{-\alpha_j \Delta}}{\alpha_j^2} \right) \right] \\
 & + \frac{\mu_s}{\mu} \frac{3}{16} A_3 \mu \left[ \left( \frac{1 - e^{-\alpha_j \Delta}}{\alpha_j} \right) G + H \left( \frac{x_2 - x_1 e^{-\alpha_j \Delta}}{\alpha_j} - \frac{1 - e^{-\alpha_j \Delta}}{\alpha_j^2} \right) \right] \quad (29)
 \end{aligned}$$

where  $\Delta = \Delta x = x_2 - x_1$

$$A = \frac{1}{\Delta x} \{ x_2 (F I 0(J, x_1) + \mu^2 F I 2(J, x_1)) - x_1 (F I 0(J, x_2) + \mu^2 F I 2(J, x_2)) \}$$

$$B = \frac{1}{\Delta x} \{ F I 0(J, x_2) - F I 0(J, x_1) + \mu^2 (F I 2(J, x_2) - F I 2(J, x_1)) \}$$

$$\begin{aligned}
 C &= \frac{1}{\Delta x} \{x_2 FI3(J, x_1) - x_1 FI3(J, x_2) + \mu^2 (x_2 FI1(J, x_1) - x_1 FI1(J, x_2))\} \\
 D &= \frac{1}{\Delta x} \{FI3(J, x_2) - FI3(J, x_1) + \mu^2 (FI1(J, x_2) - FI1(J, x_1))\} \\
 E &= \frac{1}{\Delta x} \{x_2 (FI0(J+1, x_1) + \mu^2 FI2(J+1, x_1)) - x_1 (FI0(J+1, x_2) \\
 &\quad + \mu^2 FI2(J+1, x_2))\} \\
 F &= \frac{1}{\Delta x} \{FI0(J+1, x_2) - FI0(J+1, x_1) + \mu^2 (FI2(J+1, x_2) - FI2(J+1, x_1))\} \\
 G &= \frac{1}{\Delta x} \{x_2 FI3(J+1, x_1) - x_1 FI3(J+1, x_2) + \mu^2 (x_2 FI1(J+1, x_1) \\
 &\quad - x_1 FI1(J+1, x_2))\} \\
 H &= \frac{1}{\Delta x} \{FI3(J+1, x_2) - FI3(J+1, x_1) + \mu^2 (FI1(J+1, x_2) - FI1(J+1, x_1))\}
 \end{aligned} \tag{30}$$

$$with FI0(J, x_i) = 3I_{o_j}(x_i) - I_{2_j}(x_i)$$

$$FI2(J, x_i) = 3I_{2_j}(x_i) - I_{o_j}(x_i)$$

$$FI3(J, x_i) = 3I_{3_j}(x_i) - 5I_{1_j}(x_i)$$

$$FI1(J, x_i) = 3I_{1_j}(x_i) - 5I_{3_j}(x_i)$$

$$A_1 = 3\bar{\gamma}_j + \frac{h}{m_o c^2} \frac{v_j^2}{\Delta v_j}$$

$$A_2 = 1 - A_1$$

(31)

$$A_3 = \frac{h}{m_o c^2} \frac{v_{j+1}^2}{\Delta v_{j+1}}$$

Computationally, the solution to the integral will be inaccurate due to figure loss when  $\alpha_j \Delta < 0.01$ . Under this condition, one finds the following solution:

$$\begin{aligned}
 & \frac{\mu_s}{\mu} \int_{x_1}^{x_2} \{ \} e^{-\alpha_j(x_2 - x')} dx' \\
 &= \frac{\mu_s}{\mu} \frac{3}{16} A_2 \Delta \left[ (FI0(J, x_1) + \mu^2 FI2(J, x_1)) \left( 1 - \frac{1}{2} \alpha_j \Delta \right) + \frac{\Delta}{2} B \right] \\
 & - \frac{\mu_s}{\mu} \Delta \frac{3}{16} \mu A_1 \left[ \left( 1 - \frac{1}{2} \alpha_j \Delta \right) (FI3(J, x_1) + \mu^2 FI1(J, x_1)) + \frac{\Delta}{2} D \right] \\
 & + \frac{\mu_s}{\mu} \Delta \frac{3}{16} A_3 \left[ \left( 1 - \frac{1}{2} \alpha_j \Delta \right) (FI0(J+1, x_1) + \mu^2 FI2(J+1, x_1)) + \frac{\Delta}{2} F \right] \\
 & + \frac{\mu_s}{\mu} \Delta \frac{3}{16} \mu A_3 \left[ \left( 1 - \frac{1}{2} \alpha_j \Delta \right) (FI3(J+1, x_1) + \mu^2 FI1(J+1, x_1)) + \frac{\Delta}{2} H \right] \quad (32)
 \end{aligned}$$

#### DISCUSSION OF THE NUMERICAL APPROXIMATIONS

In Eq. (7), the material heating rate for Compton scattering was derived. It will now be shown that numerically the heating rate calculated will be consistent with the analytical result. Integrating Eq. (24) over  $\mu$ , one finds

$$\begin{aligned}
 \frac{1}{c} \nabla \cdot \mathbf{F}_j &= \mu_{a_j} (2B_j - E_j) - \mu_s E_j (1 - 2\bar{\gamma}_j) + \mu_s E_j \left( 1 - 3\bar{\gamma}_j - \frac{h}{m_o c^2} \frac{v_j^2}{\Delta v_j} \right) \\
 & + \mu_s \frac{h}{m_o c^2} \frac{v_{j+1}^2}{\Delta v_{j+1}} E_{j+1} \quad (33)
 \end{aligned}$$

Summing on  $j$ , one obtains

$$\frac{1}{c} \nabla \cdot F = \sum_{j=1}^N \mu_{a_j} (2B_j - E_j) - \mu_s \sum_{j=1}^N \bar{\gamma}_j E_j - \mu_s \frac{h}{m_o c^2} \sum_{j=1}^N \left( E_j \frac{\nu_j^2}{\Delta \nu_j} - E_{j+1} \frac{\nu_{j+1}^2}{\Delta \nu_{j+1}} \right) \quad (34)$$

For the third term on the right-hand side, note that

$$\sum \left( E_j \frac{\nu_j^2}{\Delta \nu_j} - E_{j+1} \frac{\nu_{j+1}^2}{\Delta \nu_{j+1}} \right) = E_1 \frac{\nu_1^2}{\Delta \nu_1} - E_{N+1} \frac{\nu_{N+1}^2}{\Delta \nu_{N+1}} \quad (35)$$

Now, by definition,  $\nu_1 = 0$  and  $E_{N+1} = 0$ . Thus, the calculated heating rate due to Compton scattering is

$$\mu_s c \sum_{j=1}^N \bar{\gamma}_j E_j \quad (36)$$

Comparing this with the analytical result, i. e.,

$$\mu_s \frac{hc}{m_o c^2} \int_0^\infty \nu E_\nu d\nu \quad (37)$$

one notes that the above sum is consistent with assuming a flat spectrum over the various groups. Moreover, in the limit as the number of groups increases, the sum approaches the analytical result.

A computational difficulty encountered with the present formulation is that the intensity  $I_j(\mu)$  depends on the solid-angle moments of the intensity  $I_{0j}$ ,  $I_{1j}$ , etc. However, if these quantities were known, the evaluation of the intensity would be superfluous, since one is usually concerned with finding heating rates. The solution to the transport equation gives the intensity  $I_j^{n+1}(\mu)$  at time  $n + 1$  in terms of the moment quantities at time  $n$ . The assumption made is that the moment quantities at time  $n$  are close to those at time  $n + 1$ . Computationally, there are two equivalent ways to achieve this result. The first way, and the least desirable, is to run the

calculation with very small time increments. Now, even if the moment quantities are rapidly changing, the errors involved are proportional to the time step. After a number of cycles in which the solution is constant, the moment quantities will also converge. A second way to guarantee that the moment quantities are consistent with the transport equation is to iterate on these quantities. A test can be made to see if the values at time  $n + 1$  are different from those at time  $n$ . If differences are detected, the complete transport calculations can be rerun using the new, updated, moment quantities. This scheme is presently employed in the Thomson and Compton scattering subroutines. To calculate the number of iterations required for convergence, consider the following model for the case of Thomson scattering, i.e.,  $\gamma \rightarrow 0$ . For this case, the transport equation becomes

$$\mu \frac{dI}{dx}^{n+1} = \mu_a' (B - I^{n+1}) - \mu_s I^{n+1} + \frac{3}{16} \mu_s [(1 - \mu^2) I_o^n + (3\mu^2 - 1) I_2^n] \quad (38)$$

If spatial homogeneity is assumed, then

$$I^{n+1}(\mu) = \frac{\mu_a}{\bar{\mu}} B + \frac{3}{16} \frac{\mu_s}{\bar{\mu}} [(3 - \mu^2) I_o^n + (3\mu^2 - 1) I_2^n] \quad (39)$$

where  $\bar{\mu} = \mu_a + \mu_s$ .

To find  $I_o^{n+1}$  and  $I_2^{n+1}$ , integrate over  $d\mu$ . Thus,

$$I_o^{n+1} = \frac{2\mu_a}{\bar{\mu}} B + \frac{\mu_s}{\bar{\mu}} I_o^n \quad (40)$$

and

$$I_2^{n+1} = \frac{2}{3} \frac{\mu_a}{\bar{\mu}} B + \frac{3}{10} \frac{\mu_s}{\bar{\mu}} I_o^n + \frac{1}{10} \frac{\mu_s}{\bar{\mu}} I_2^n \quad (41)$$

Solving these recursion equations, one finds that

$$I_o^n = 2B \left[ 1 - \left( \frac{\mu_s}{\bar{\mu}} \right)^n \right] \quad (42)$$

and

$$I_2^n = \frac{2B}{3} \left[ 1 - \left( \frac{\mu_s}{\bar{\mu}} \right)^n - \left( \frac{\mu_s}{10\bar{\mu}} \right)^2 \right] \quad (43)$$

From the form of these equations, one can conclude that the convergence rate is controlled by the ratio  $\mu_s / \mu_a + \mu_s$ . Thus, in the interesting situation in which scattering is the dominant mechanism, the number of iterations required to achieve a given accuracy increases as the amount of absorption decreases. This fact should be kept in mind when employing the scattering code in regions with dominant scattering.

For the case of Compton scattering, consider the homogeneous transport equation, or

$$\begin{aligned} 0 = \mu_a (B - I^{n+1}) - \mu_s I^{n+1} (1 - 2\gamma) + \frac{3}{16} \mu_s [(3 - \mu^2) I_o^n + (3\mu^2 - 1) I_2^n] \\ - \frac{9}{16} \mu_s \gamma [(3 - \mu^2) I_o^n + \mu(3\mu^2 - 5) I_1^n + (3\mu^2 - 1) I_2^n + \mu(3 - 5\mu^2) I_3^n] \\ + \frac{3}{16} \mu_s \frac{\partial}{\partial \gamma} \gamma^2 [(3 - \mu^2) I_o^n + \mu(3\mu^2 - 5) I_1^n + (3\mu^2 - 1) I_2^n + \mu(3 - 5\mu^2) I_3^n] \end{aligned} \quad (44)$$

If one now integrates over  $\int_{-1}^1 d\mu$ , one finds

$$\bar{\mu} I_o^{n+1} = 2\mu_a B + \mu_s (1 - 3\gamma) I_o^n + \mu_s \frac{\partial(\gamma^2 I_o^n)}{\partial \gamma} \quad (45)$$

where  $\bar{\mu} = \mu_a + \mu_s (1 - 2\gamma)$ .

Integrating over frequency, one finds

$$I_o^{n+1} = \frac{2\mu_a}{\bar{\mu}_j} B_j + \frac{\mu_s}{\bar{\mu}_j} (1 - 3\bar{\gamma}_j) I_o^n + \frac{\mu_s}{\bar{\mu}_j} \frac{h}{m_o c^2} \left[ \frac{\nu_{j+1}^2}{\Delta \nu_{j+1}} I_o^n_{j+1} - \frac{\nu_j^2}{\Delta \nu_j} I_o^n_j \right] \quad (46)$$

where

$$\bar{\mu}_j = \mu_a + \mu_s (1 - 2\bar{\gamma}_j) \quad (47)$$

Since  $I_{0k+1} = 0$ , the solution to the last equation is

$$I_{0k}^n = B_k \left[ \frac{2\mu_{a_k}}{\mu_{a_k} + \mu_s \left( \bar{\gamma}_k + \frac{h}{m_o c^2} \frac{v_k^2}{\Delta v_k} \right)} \right] \left\{ 1 - \left[ \frac{\mu_s}{\mu_k} \left( 1 - 3\bar{\gamma}_k - \frac{h}{m_o c^2} \frac{v_k^2}{\Delta v_k} \right) \right]^n \right\} \quad (48)$$

The remaining equations can be solved in reverse order. What is significant is that the coefficient that determines the convergence is frequency-dependent, and, moreover, the higher the frequency, the faster the convergence, provided  $\gamma < 0.2$ .

Computationally, the iteration scheme has been developed for both the Thomson and Compton scattering subroutines. Presently, the options available for the iteration scheme are the number of iterations, an accuracy criterion to determine convergence, and an extrapolation switch allows one to use the rate of convergence to extrapolate to find the desired function.

### SPHERICAL GEOMETRY

In spherical geometry, the transport equation (1) is replaced by

$$\frac{dI}{dx} = \mu'_a (B_v - I) + S \quad (49)$$

where  $x$  is now a coordinate measuring distance along a characteristic ray. In the treatment of the scattering terms, the only difference is that the  $\mu$ 's appearing in  $S(x)$  must be replaced by an average  $\bar{\mu} = (\mu_1 + \mu_2)/2$ , where  $\mu_1$  and  $\mu_2$  are the cosines of the angles of the characteristic ray with respect to the normal at  $x_1$  and  $x_2$ , respectively.

### EQUATIONS OF MOTION

The same independent coordinates are used in the OUTPUT code and SPUTTER code: mass  $m$  and time  $t$ . The Lagrangian mass

coordinate  $m$  is defined by

$$m = \int_0^r \alpha r'^{\alpha-1} \rho(r', t) dr' \quad (50)$$

where  $\rho$  is the density and

$$\alpha = \begin{cases} 1 & \text{plane} \\ 3 & \text{sphere} \end{cases}$$

This coordinate has the physical units of mass/cm<sup>2</sup> in plane geometry; whereas in spherical geometry, it is the mass interior to radius  $r$  divided by  $4\pi/3$ . If  $\tau$  is the specific volume, the equation of continuity is given by

$$\frac{d\tau}{dt} = \tau r^{-(\alpha-1)} \frac{\partial}{\partial r} \left( r^{\alpha-1} \frac{dr}{dt} \right) \quad (51)$$

The momentum equation can be written as

$$\rho \frac{du}{dt} = \frac{\partial \sigma_{11}}{\partial r} + \frac{(\alpha - 1)}{2r} (2\sigma_{11} - \sigma_{22} - \sigma_{33}) \quad (52)$$

where  $u = dr/dt$  and  $(\sigma_{ij})$  is the symmetric stress tensor. The stress tensor is given by

$$\sigma_{ij} = -P_m \delta_{ij} - P_{ij} \quad (53)$$

where  $P_m$  is the material pressure including artificial viscosity, and  $P_{ij}$  is the radiation pressure tensor. In the case of plane geometry,  $\alpha = 1$  and the equation becomes

$$\rho \frac{du}{dt} = - \frac{\partial(P_m + P_r)}{\partial r} \quad (54)$$

where  $P_r$  is the radiation pressure in the  $r$  direction.

This is the standard equation solved in the SPUTTER code. However, the SPUTTER program assumes that  $P_r = (1/3) a \theta^4$ , the equilibrium diffusion value. This assumption is not made in the OUTPUT program; instead, the

radiation pressure for a zone is computed as the second angular moment of the intensity. Since the intensity and its moments are computed on zone boundaries  $r_i$ , the radiation pressure in the zone  $r_i \leq r < r_{i+1}$  is defined as the arithmetic mean

$$P_{r_i+\frac{1}{2}} = \frac{1}{2} (P_{r_i} + P_{r_{i+1}}) \quad (55)$$

In the case of spherical symmetry, one must evaluate  $\sigma_{22}$  and  $\sigma_{33}$ . The radiation pressure tensor is given by

$$(P) = \frac{1}{2\pi} \int_{4\pi} d\Omega \vec{\Omega} \vec{\Omega} I(\mu) \quad (56)$$

where  $\mu$  is the cosine of the angle  $\theta$  between the ray  $\vec{\Omega}$  and the radius vector  $\vec{r}$ . If  $\vec{r}$  is taken as a polar axis, an azimuthal angle  $\phi$  can be introduced and  $\vec{\Omega}$  can be expressed as a column vector

$$\vec{\Omega} = \begin{bmatrix} \Omega_1 \\ \Omega_2 \\ \Omega_3 \end{bmatrix} = \begin{bmatrix} \mu \\ \sqrt{1 - \mu^2} \cos \phi \\ \sqrt{1 - \mu^2} \sin \phi \end{bmatrix} \quad (57)$$

The dyadic  $\vec{\Omega} \vec{\Omega}$  is then obtained by postmultiplying this column by its transpose, which gives

$$(P) = \frac{1}{2\pi} \int_{-1}^1 d\mu \int_0^{2\pi} d\phi I(\mu) \begin{bmatrix} \mu^2 & \mu \sqrt{1 - \mu^2} \cos \phi & \mu \sqrt{1 - \mu^2} \sin \phi \\ \mu \sqrt{1 - \mu^2} \cos \phi (1 - \mu^2) \cos^2 \phi & (1 - \mu^2) \cos \phi \sin \phi & \mu \sqrt{1 - \mu^2} \sin \phi (1 - \mu^2) \sin \phi \cos \phi (1 - \mu^2) \sin^2 \phi \end{bmatrix} \quad (58)$$

On evaluating the integrals, one obtains

$$P_{ij} = 0 \quad i \neq j \quad (59)$$

$$P_{11} = \int_{-1}^1 \mu^2 I(\mu) d\mu \quad (60)$$

$$P_{22} = P_{33} = \frac{1}{2} \int_{-1}^1 (1 - \mu^2) I(\mu) d\mu \quad (61)$$

Since  $P_{22} = P_{33} = 1/2 [E_r - P_{11}]$ , where  $E_r$  is the radiation energy density,

$$E_r = \int_{-1}^1 I d\mu \quad (62)$$

Eq. (52) for the spherical case is

$$\rho \frac{du}{dt} = - \frac{\partial(P_m + P_r)}{\partial r} - \frac{1}{r} (3P_r - E_r) \quad (63)$$

where  $P_r = P_{11}$ . The last term in Eq. (63) is in some cases a source of numerical noise, particularly at small radius  $r$ . Hence, a parameter  $S4$  is employed in the code such that in the "diffusion" case,  $P_r/E_r < S4$ , a difference representation of Eq. (63) is used. However, in the "streaming" case,  $P_r/E_r > S4$ , one can rewrite Eq. (63) as

$$\rho \frac{\partial u}{\partial t} = - \frac{\partial P_m}{\partial r} - \frac{1}{2} \frac{\partial}{\partial r} (r^2 P_r) - \frac{1}{r} (P_r - E_r) \quad (64)$$

and a difference equation based on this form is used.

$S4$  should be assigned on the basis of the characteristics of the problem being solved; a typical value is  $1/2$ .

### ENERGY EQUATION

The equation for conservation of energy is given by

$$\frac{dE_m}{dt} + \tau \left[ \frac{\partial E_r}{\partial t} + \nabla \cdot \vec{F} \right] = -(P_m + P_r) \dot{r} + \frac{(\alpha - 1)^\tau}{2r} (3P_r - E_r) u + \dot{e}_s \quad (65)$$

where  $E_m$  is the material specific internal energy,  $\vec{F}$  is the flux of radiation,

and  $\dot{e}_s$  is an external energy source rate per unit mass. The standard SPUTTER code takes the radiative heating rate per unit mass,  $\dot{e}_r$ , as

$$\dot{e}_r = -(\nabla \cdot \vec{F}) \tau. \quad (66)$$

However, the OUTPUT code includes the radiative streaming contribution, and  $\dot{e}_r$  is calculated in the spherical case as

$$\dot{e}_r = \left[ -\nabla \cdot \vec{F} + \frac{1}{r} (3P_r - E_r) u \right] \tau. \quad (67)$$

The form  $\tau(\partial E_r / \partial t)$  is combined with  $dE_m / dt$ , so that Eq. (65) is treated as an equation for the rate of change of total specific energy. The difference equations are then developed as in the standard SPUTTER code.

The energy source rates  $\dot{e}_s$  are determined by the source routines QUE8, QUE9, QUE10. These source routines are described briefly in Appendix I.

After determining the change in total specific energy during a time interval due to  $\dot{e}_r$ ,  $\dot{e}_s$ , and the radiative and material work terms as shown in Eq. (65), the material temperature is advanced by inverting the equation of state to find the temperature corresponding to the new values of  $E_m$  and  $\tau$ .

In order to avoid excessive restriction of the time step during the early portion of a calculation, it has been found advisable to utilize the equilibrium diffusion approximation

$$E_r = a \theta^4$$

for the energy density of the radiation field in the energy conservation equation. This approximation is valid in the core, where most of the energy exchange between the field and the material occurs.

### SECTION III

#### OUTPUT CODE PROGRAMMING

##### LOGIC OF THE CODE

The two basic independent variables in the OUTPUT code are mass and time. Each of these continuous variables is divided into discrete elements, the mass into an array of zones of variable size  $M_i$  and the time into a series of time steps  $\Delta t^n$ . The properties of the system are calculated for all zones at time  $t^{n+1}$ , from given values at time  $t^n$ . There are essentially two kinds of properties which define the system: kinematic properties, such as the positions  $R_i$  and velocities  $\dot{R}_i$  of the zone boundaries; and thermodynamic properties, such as the specific volumes  $\tau_i$ , the temperatures  $\theta_i$ , the specific internal energies  $E_i$ , the material pressures  $P_i$ , and the specific heats at constant volume  $C_v$ . A specification of  $M_i$ ,  $R_i$ ,  $\tau_i$ , and  $\theta_i$  completely defines the system at any given time, assuming that local thermodynamic equilibrium prevails throughout the system. All kinematic quantities are functions of the first two variables, and all thermodynamic properties are functions of the last two. The heating rates in each zone--specifically, heating from an external source ( $\dot{e}_{si}$ ) and heating by radiation transport within the system ( $\dot{e}_{ri}$ )--will generally depend on all four variables and the time.

The program is divided into two segments (MP1 and MP2) to reduce computer storage requirements. Segment MP1 sets up the initial conditions. The temperatures  $\theta_i$ , masses  $G_i$ , velocities  $\dot{R}_i$ , and interface positions  $R_i$  for each zone are provided by card input. All other quantities necessary for complete problem specification are calculated within the MP1 segment. This segment is also utilized to redefine new zones during the course of a

calculation if, for reasons of economy or greater definition, a rezoning of the system is desired. MP1 performs the required peripheral operations, such as reading or changing dump tapes on restarts of a calculation.

Segment MP2 consists of two loops (see Fig. 1). The main loop represents one cycle, representing an advance in time, and the secondary loop represents a radiation subcycle. The division of the main loop into separate boxes corresponds to the way various phases of the calculation are parceled out to subroutines, called in turn by the main program (MP2). The labels next to the boxes are in some cases the names of the subroutines; in two cases, EOS and TEMPIT, they are names of "sub-subroutines" called by the subroutines. Unlabeled boxes signify computations (or logical decisions) made in the main program itself.

The cycle begins with a computation of the time increment  $\Delta t^{n+1/2}$ . The time increment is set equal to the minimum of (1) DTMAX1, DTMAX2, and DTMAX3, which are external time controls, (2) Courant stability conditions  $\Delta t_c^{n+1/2}$ , and (3) nine-tenths of the radiation time step  $\Delta t_R^{n-1/2}$ . The nine-tenths used for the radiation time calculation is introduced to prevent radiation subcycling due to small changes in  $\Delta t_R$  from one cycle to the next. DTR, the time increment used by those subroutines involved in the subcycling loop, is set to  $\Delta t^{n+1/2}$ . The hydrodynamics portion of the cycle is entered next, where the kinematic quantities are updated. The artificial viscosity term  $P_{2,i}$ , which enters the calculation as a pressure to be added to the material pressure, acts as a shock-smoothing term, spreading the shock structure over three zones. The next two phases of the calculation compute the heat term  $\Delta Q$  in the equation expressing the first law of thermodynamics,  $\Delta Q = \Delta E + P\Delta V$ . The heating rate due to an external source ( $\dot{e}_{si}^{n+1/2}$ ) is calculated in one of several choices of sub-subroutines, depending on the nature of the source.

The heating rate due to radiation ( $\dot{e}_{ri}^{n+1/2}$ ) is calculated in the RADTN section. The main subroutine of this section calls an auxiliary subroutine,

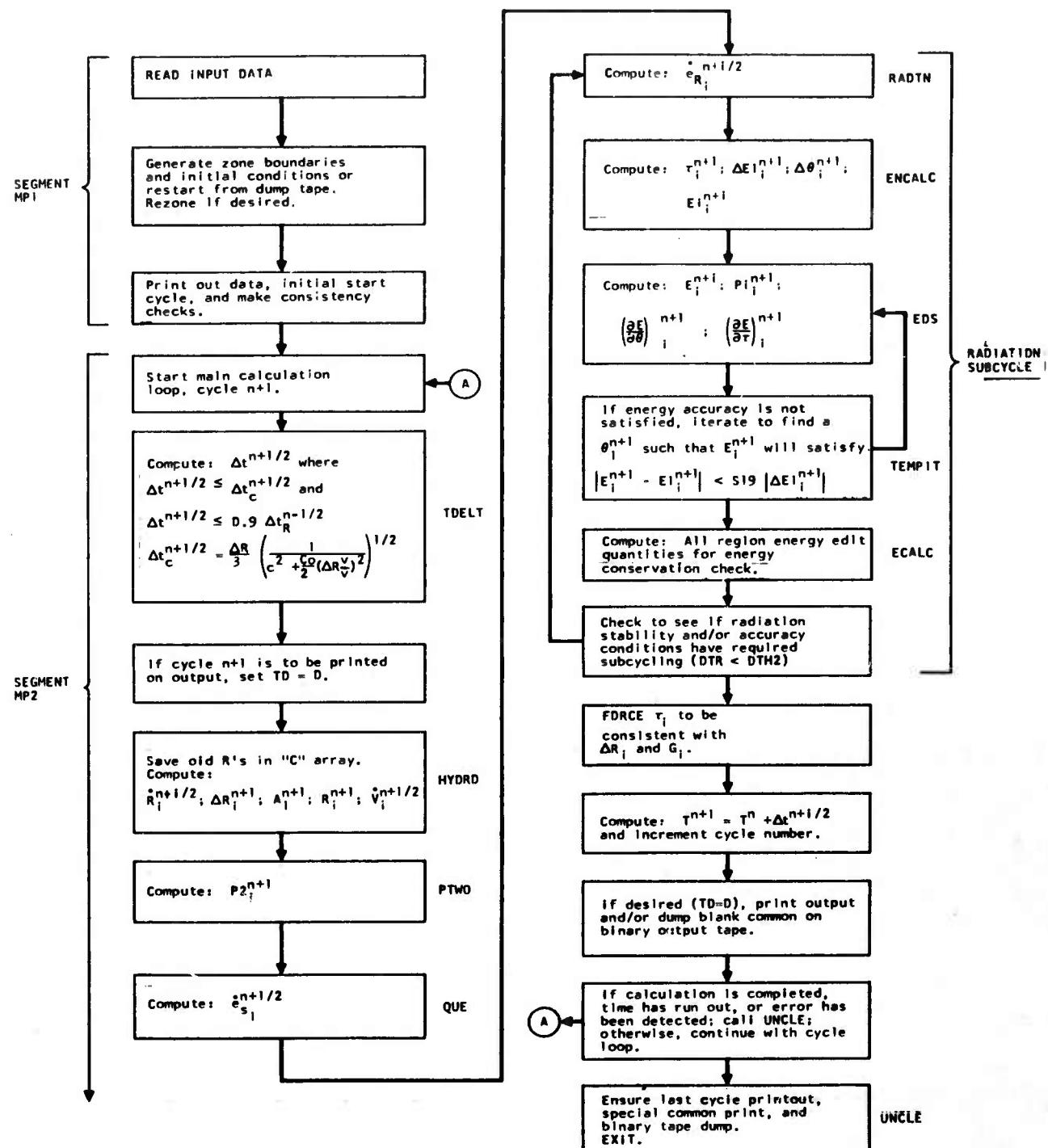


Figure 1. The OUTPUT code: Summary Flow Logic

KAPPA, for absorption coefficients. The calculation may be either frequency-independent (grey) or frequency-dependent. In the latter case, the absorption coefficients and radiation fluxes are computed in each of a desired set of frequency groups, and the fluxes are summed up to obtain the total flux at each boundary. Absorption coefficients in the grey case are Rosseland-averaged opacities and are computed from analytic fits to the numerical calculations of Stewart and Pyatt (Ref. 6). For multifrequency calculations, the program uses a table of group-averaged opacities, which are read from a data tape prepared by an auxiliary program.

ENCALC, where the next calculations are performed, gathers the heating rates, computes the work done by each zone from the pressure and the rate of volume change, and finds the increment of internal energy from the first-law equation. One then obtains the corresponding increment of temperature in each zone. After the temperature and specific volume are updated, the equation of state is used to find the internal energy  $E$ , the material pressure  $P_1$ , and the two partial derivatives of  $E$ ,  $\partial E / \partial \theta$  and  $\partial E / \partial r$ . The difference between the first-law internal energy  $E_{I_i}$  and the equation-of-state energy  $E_i$  is compared to the change in internal energy  $\Delta E_{I_i}$  as a check on the accuracy of the calculation. If the comparison shows they do not agree to within a certain specified value S19, the subroutine TEMPIT uses the regula-falsi and interval halving iteration methods to find a temperature for which the equation of state will return an acceptable  $E_i$ . At the conclusion of the energy checks in ENCALC, all quantities describing the system have been incremented from time  $t^n$  to  $t^{n+1}$ . The main program (MP2) checks the time step DTR to determine if the radiation subroutines have induced subcycling. If radiation stability requires subcycling, DTR will have been set such that  $\Delta t^{n+1/2}$  is an even multiple of DTR. Subcycling continues until the sum of all subcycling time steps is equal to  $\Delta t^{n+1/2}$ .

The last act in the cycle is the updating of the time  $t^n$  to  $t^{n+1}$  and increasing of the cycle number by unity. After some decisions have been

made regarding whether to print the information for this cycle and/or to dump common on a binary dump tape for future restarts, control is passed to the beginning of the loop and a new cycle begins.

In table I, a brief definition of terms used in figure 1 is given and the variables used for storage in the code itself are specified.

#### FLOW CHART OF RADIATION

A flow chart of the radiation hierarchy is presented in figure 2 through 6. The figures that show the various subroutines are as follows:

RAD . . . . . Figure 2  
STRANS . . . . . Figure 3  
PTRANS . . . . . Figure 4  
SCAT . . . . . Figure 5  
STEP . . . . . Figure 6

#### OUTPUT CODE GLOSSARY

This section contains a complete list of the FORTRAN variables appearing in several subroutines of the OUTPUT code. For each variable, a brief definition or description is given. A number in parentheses superscript to a variable indicates that the variable has been defined in the list of another subroutine, as follows:

- (1) SCAT
- (2) STEP
- (3) STRANS
- (4) PTRANS
- (5) RAD

Variables appear in alphabetical order within a subroutine list. The storage allocation--Blank Common, a name common, or private storage, which is used only within the subroutine--is also given.

TABLE I

## DEFINITION OF TERMS USED IN FIGURE 1

OUTPUT		
Term	Code	Definition
$\Delta t^{n+1/2}$	= DTH2	= time increment for cycle loop
$\Delta t_c^{n+1/2}$	= RDIA	= maximum time step for Courant stability
$\Delta t_R^{n+1/2}$	= DTRMIN	= maximum time step for radiation stability for current cycle
$\Delta t_R^{n-1/2}$	= DTRMIN	= maximum time step for radiation stability for previous cycle
$R_i^{n+1}$	= R(i)	= zone boundary positions at time TH
$R_i^{n+1/2}$	= RD(i)	= velocity of zone boundaries
$\Delta R_i^{n+1}$	= DELTAR(i)	= $R(i + \Delta t) - R(i)$
$A_i^{n+1}$	= A(i)	= area of zone ( $\alpha r_i^{\alpha-1}$ )
$V_i^{n+1/2}$	= VD(i)	= rate of change of zone volumes
$P_2i^{n+1}$	= P2(i)	= artificial viscosity pressure
$\dot{e}_{si}^{n+1/2}$	= SMLQ(i)	= rate of energy deposition by source
$\dot{e}_{ri}^{n+1/2}$	= ER(i)	= rate of energy deposition by radiation
$\tau_i^{n+1}$	= SV(i)	= specific volume of zone
$\Delta EI_i^{n+1}$	= PB(i)	= first-law increment of energy change
$\Delta \theta_i^{n+1}$	= W(i)	= change in temperature during cycle
$\theta_i^{n+1}$	= THETA(i)	= temperature of zone
$EI_i^{n+1}$	= EI(i)	= first-law internal energy
$E_i^{n+1}$	= E(i)	= equation-of-state internal energy
$P_1i^{n+1}$	= P1(i)	= equation-of-state material pressure
$(\partial E / \partial \theta)_i^{n+1}$	= CV(i)	= specific heat at constant volume
$(\partial E / \partial \tau)_i^{n+1}$	= PB1(i)	= self-explanatory
t	= TH	= time

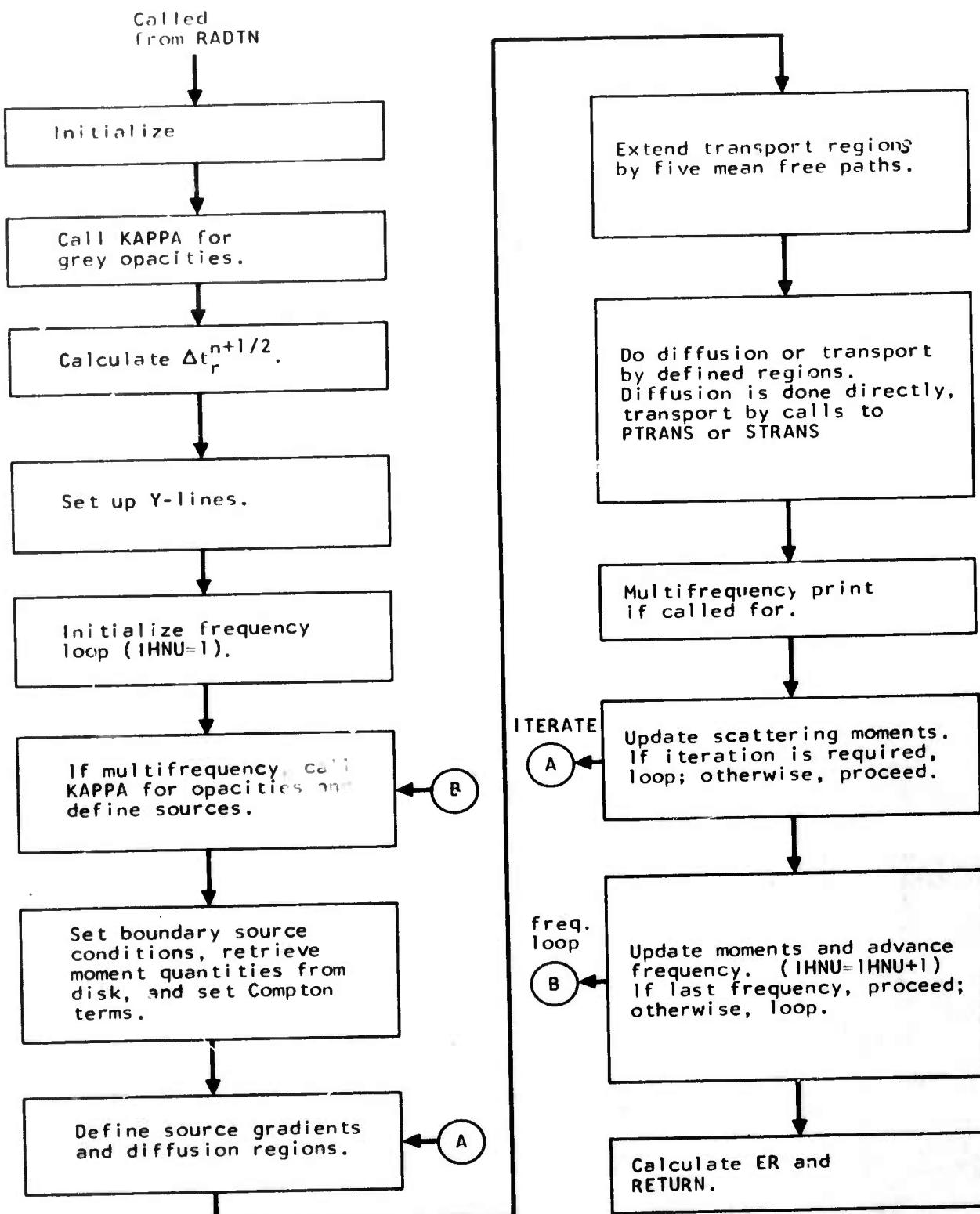


Figure 2. RAD

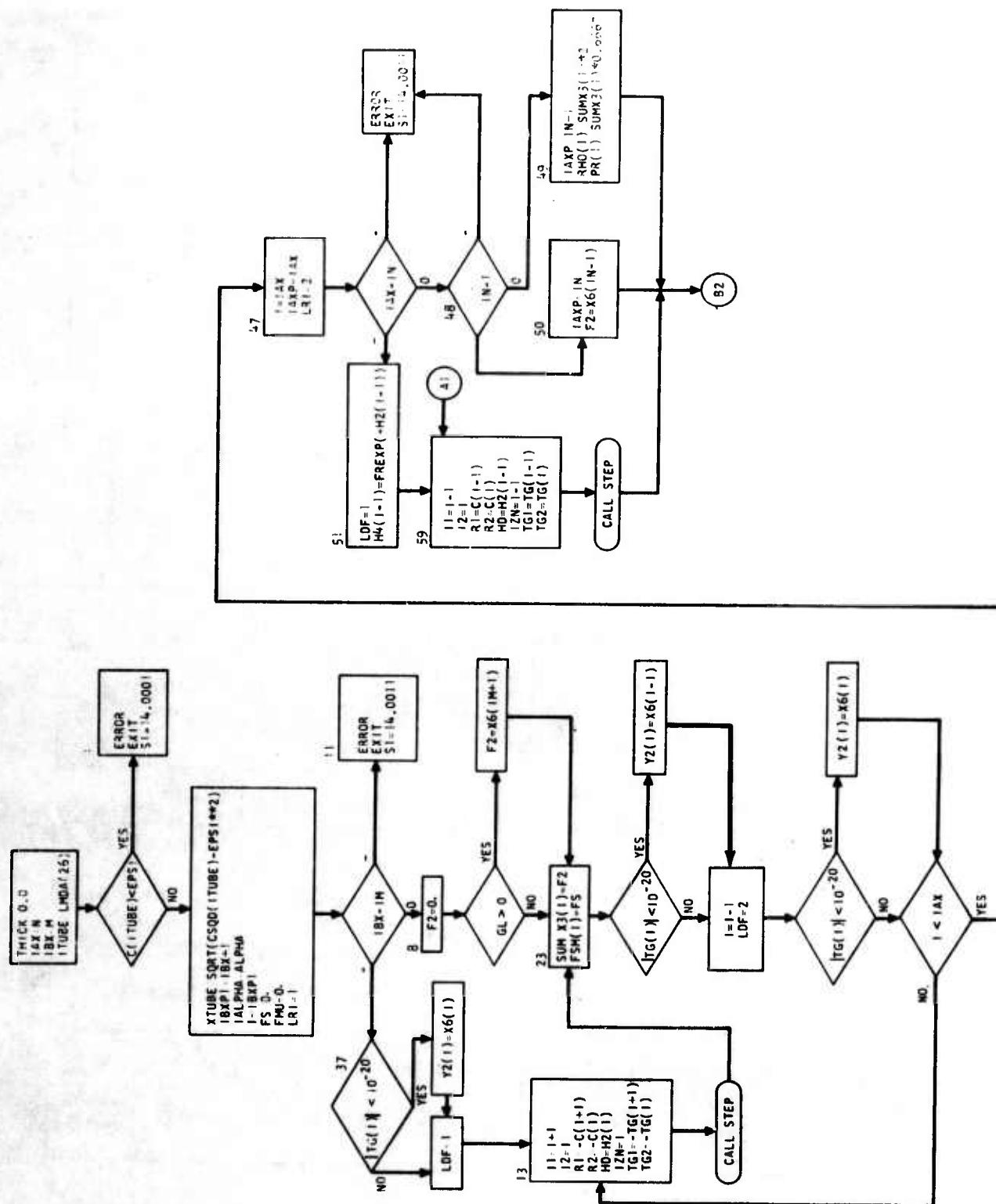


Figure 3. STRANS(N, M)

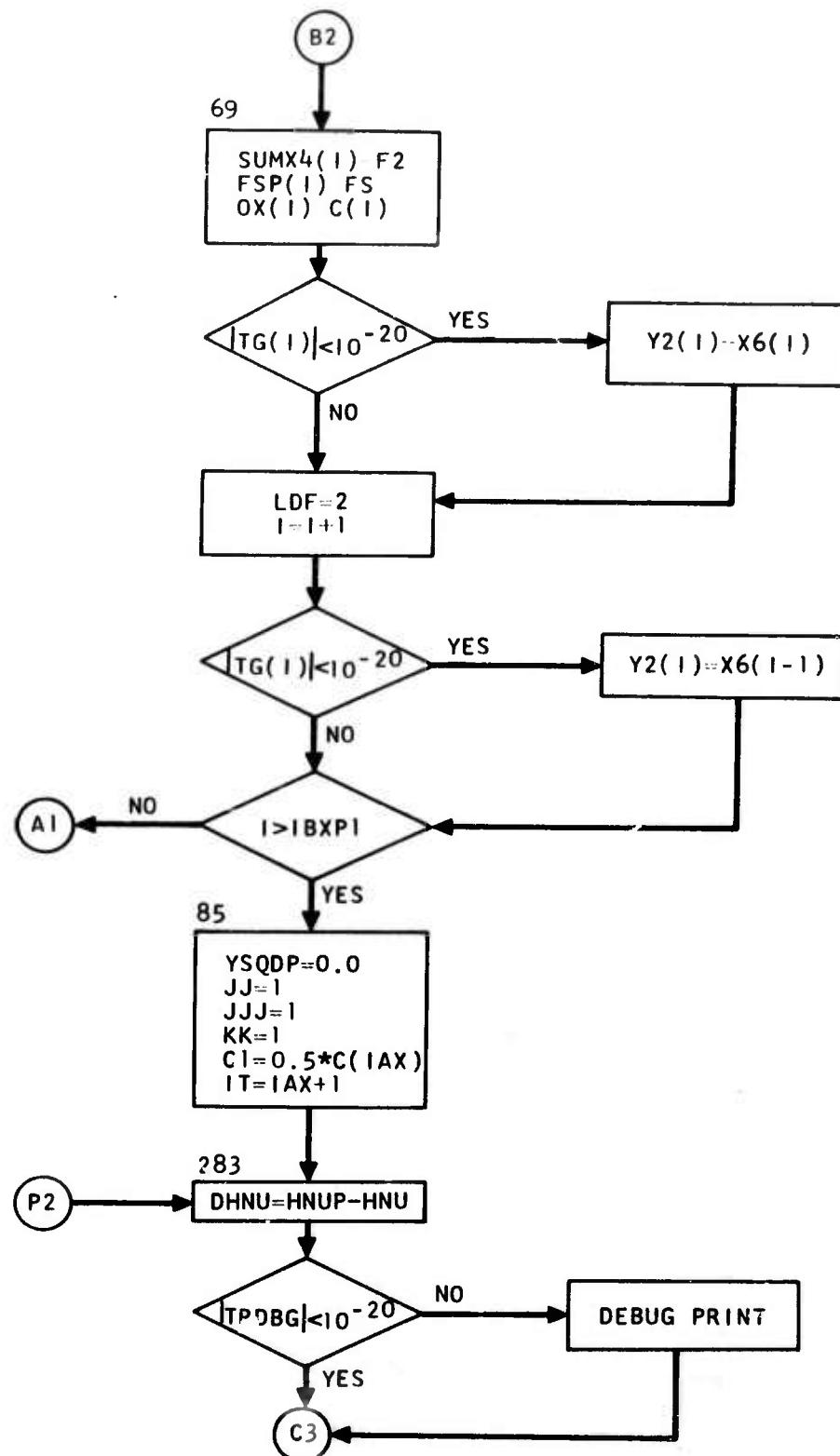


Figure 3 (continued). STRANS(N, M)

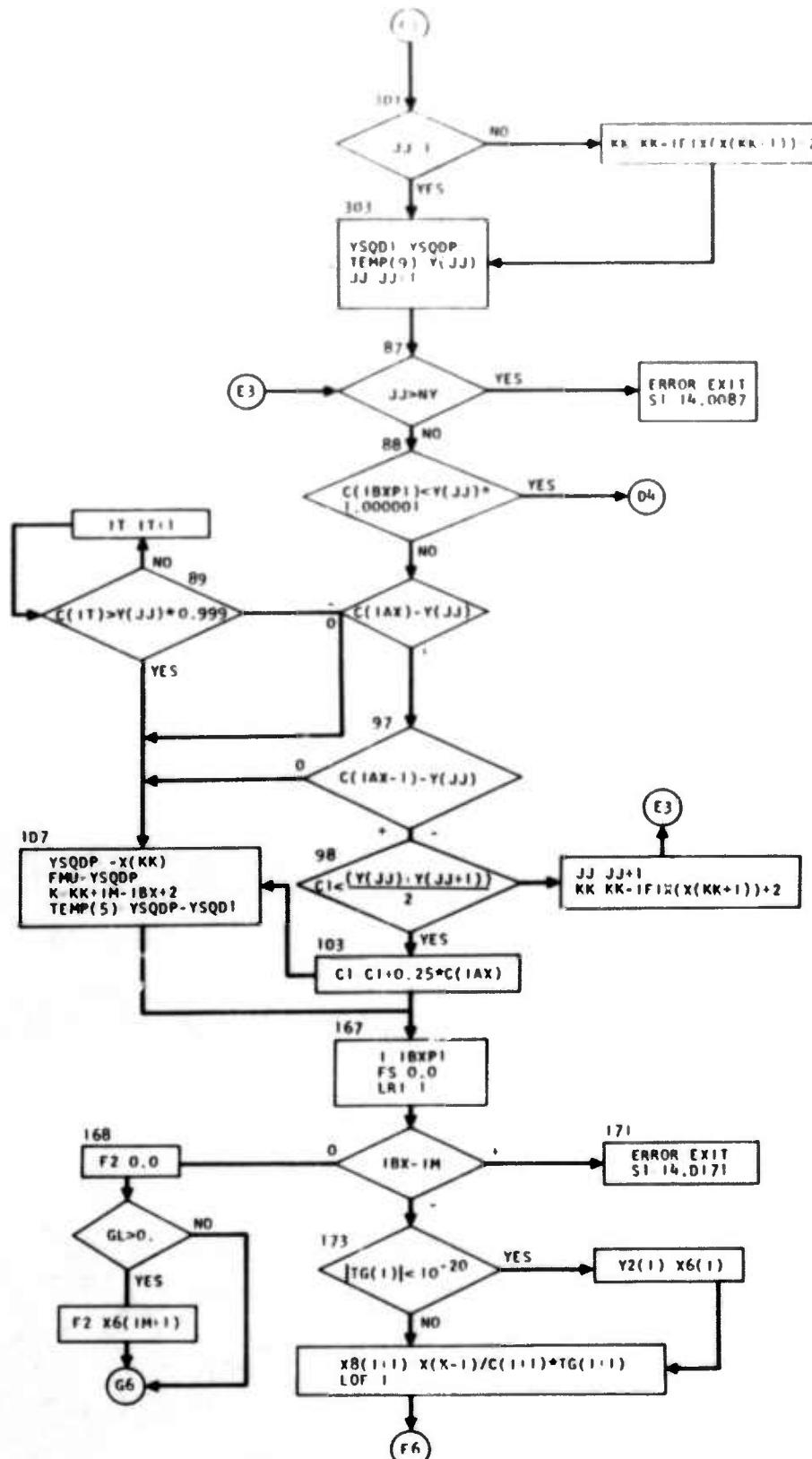


Figure 3 (continued). STRANS(N, M)

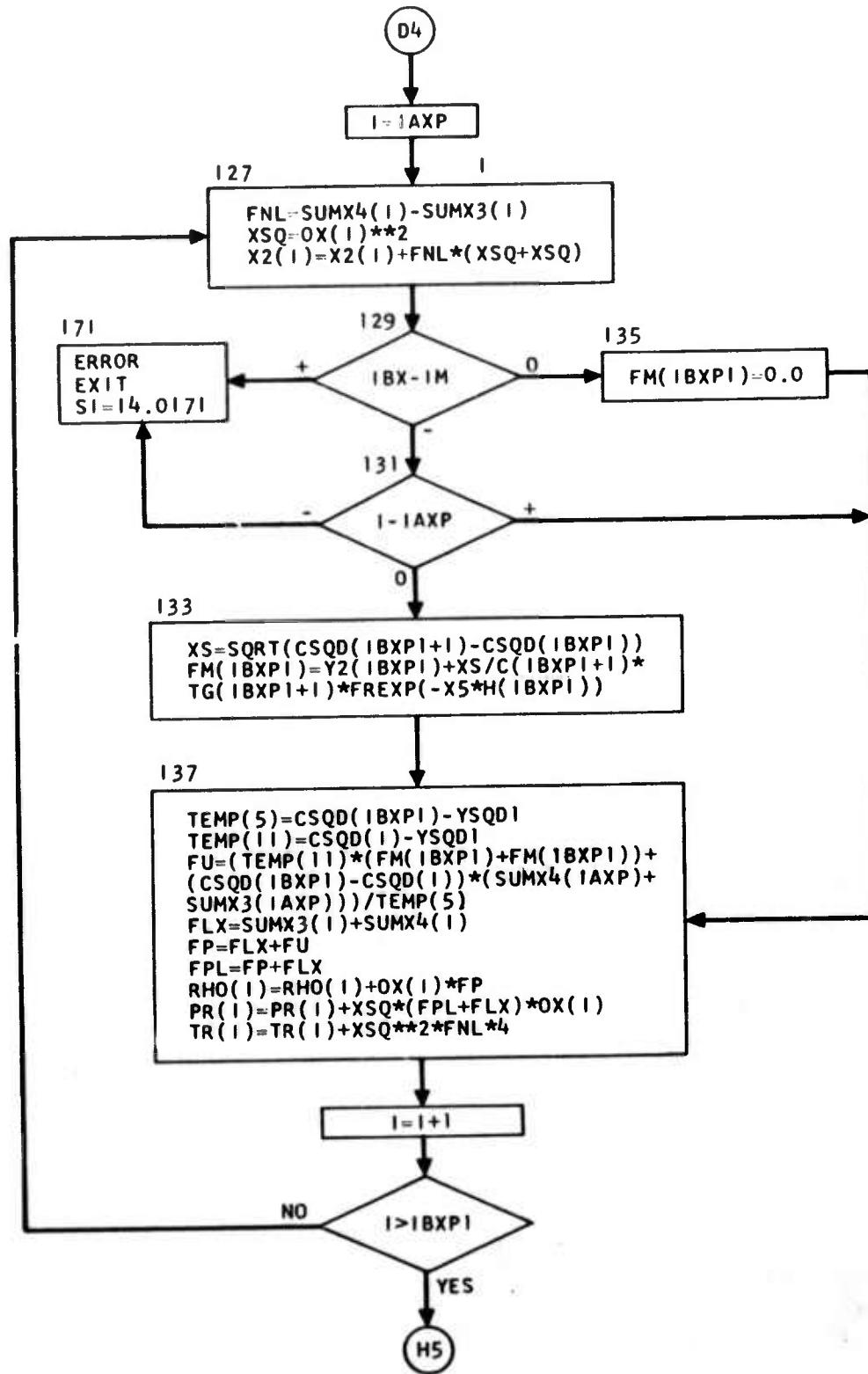


Figure 3 (continued). STRANS(N, M)

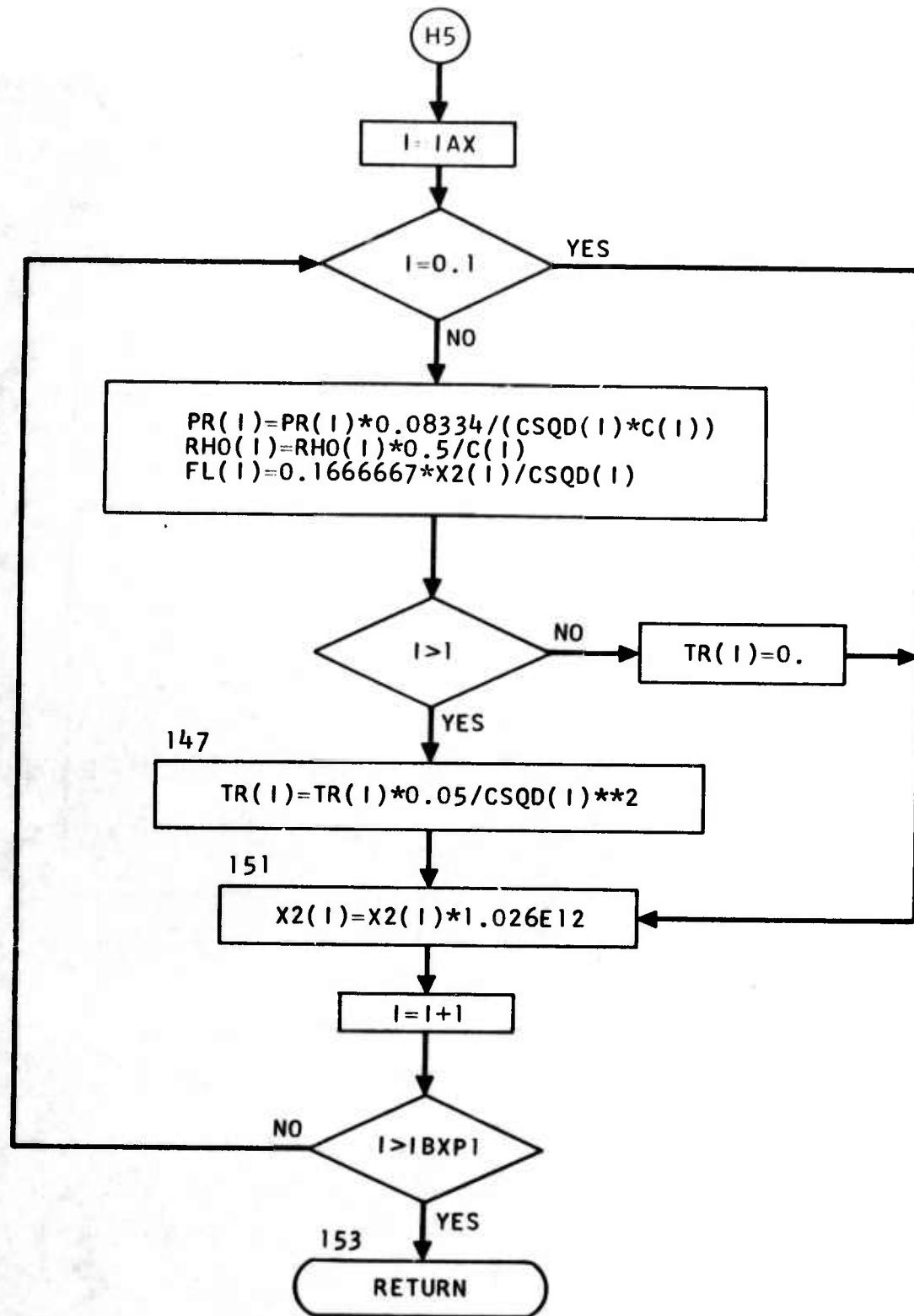


Figure 3 (continued). STRANS(N, M)

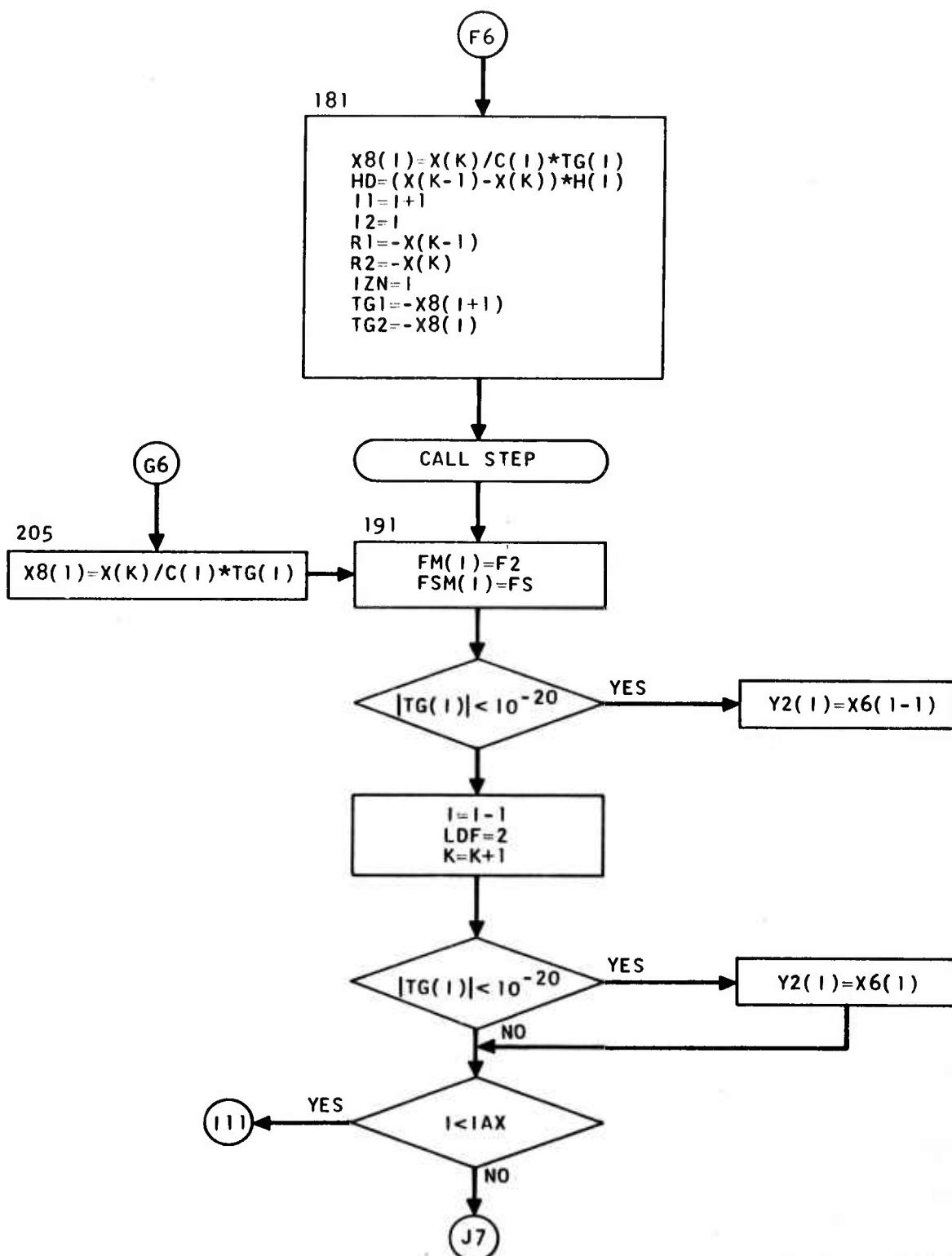


Figure 3 (continued). STRANS(N, M)

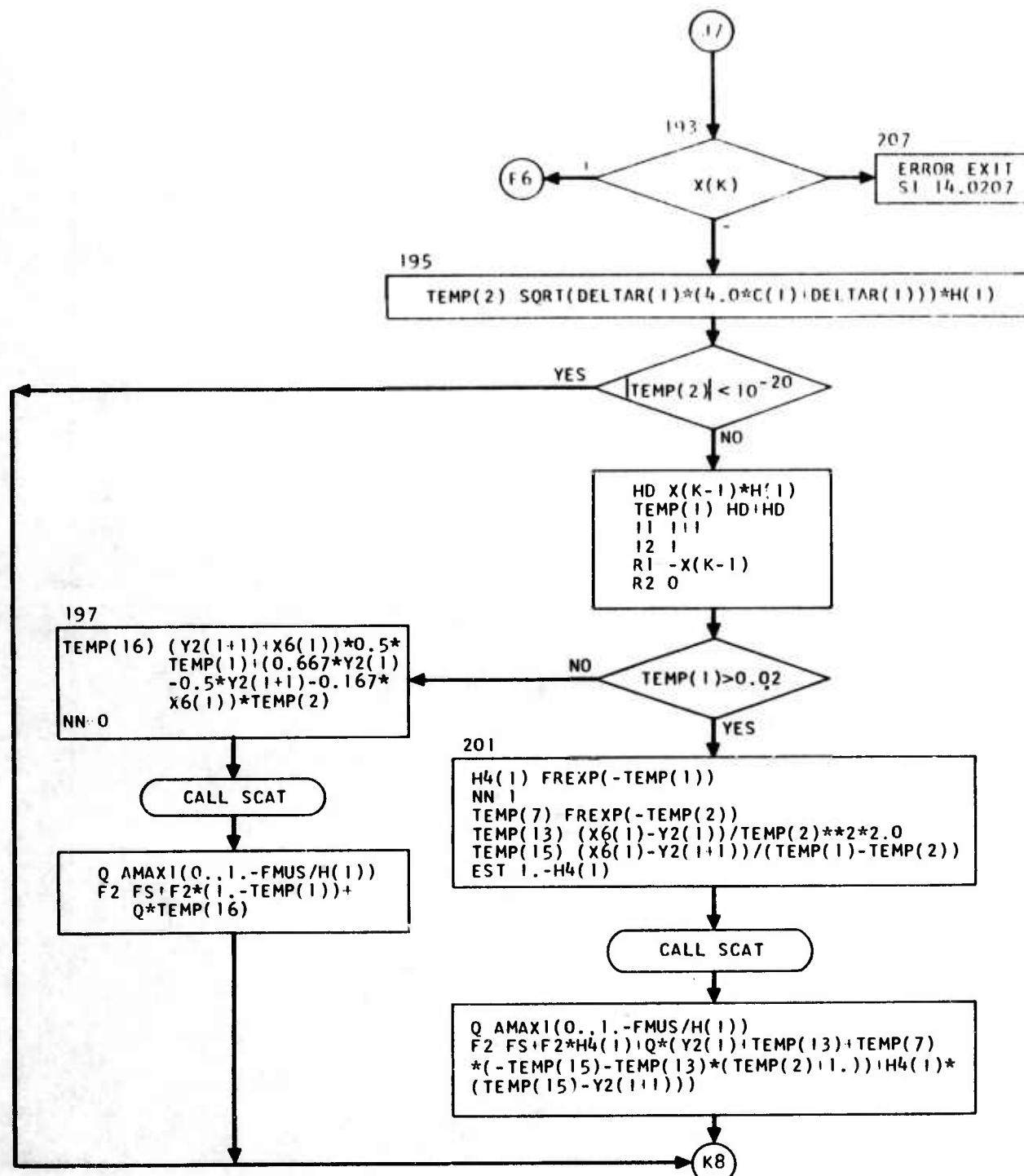


Figure 3 (continued). STRANS(N, M)

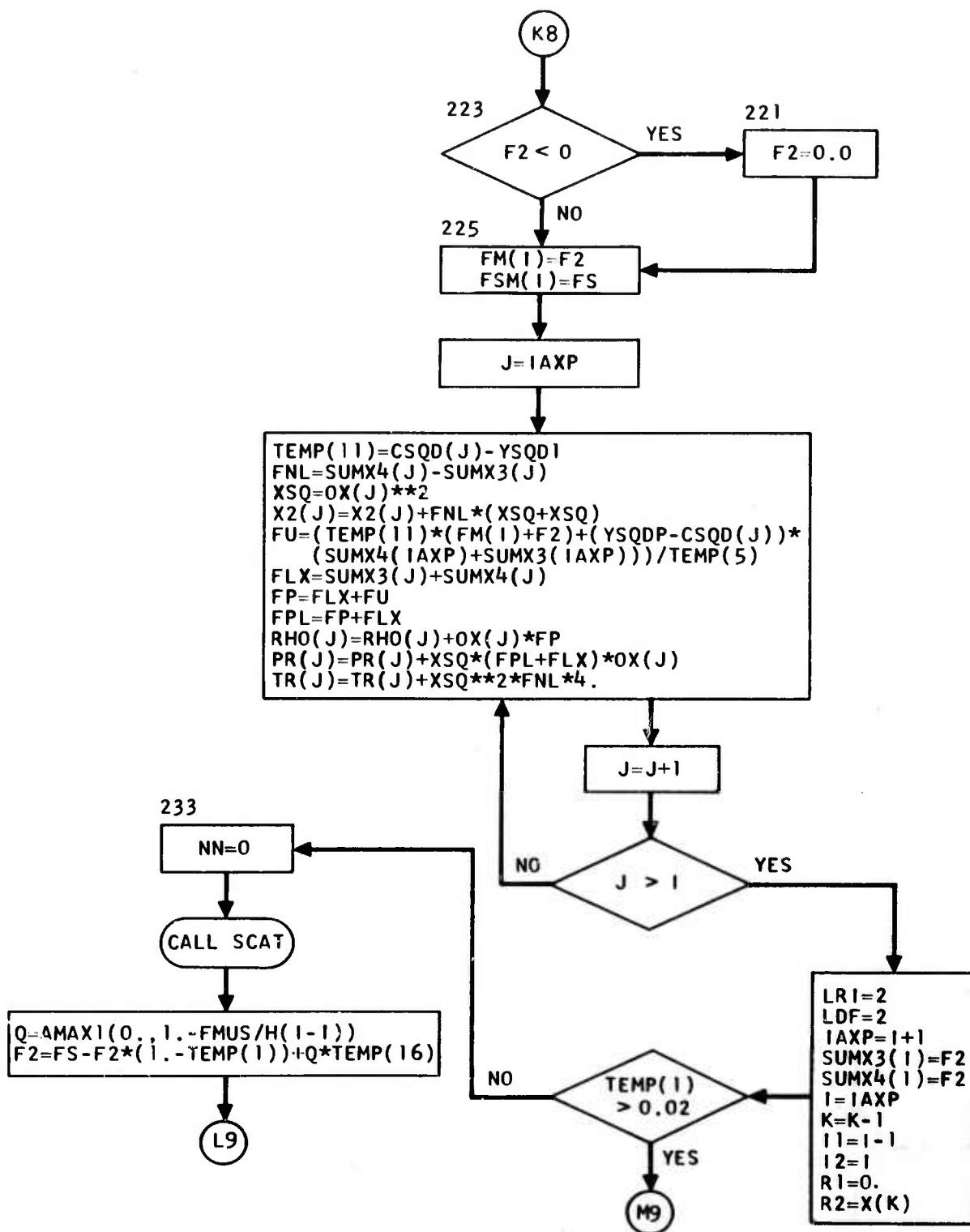


Figure 3 (continued). STRANS(N, M)

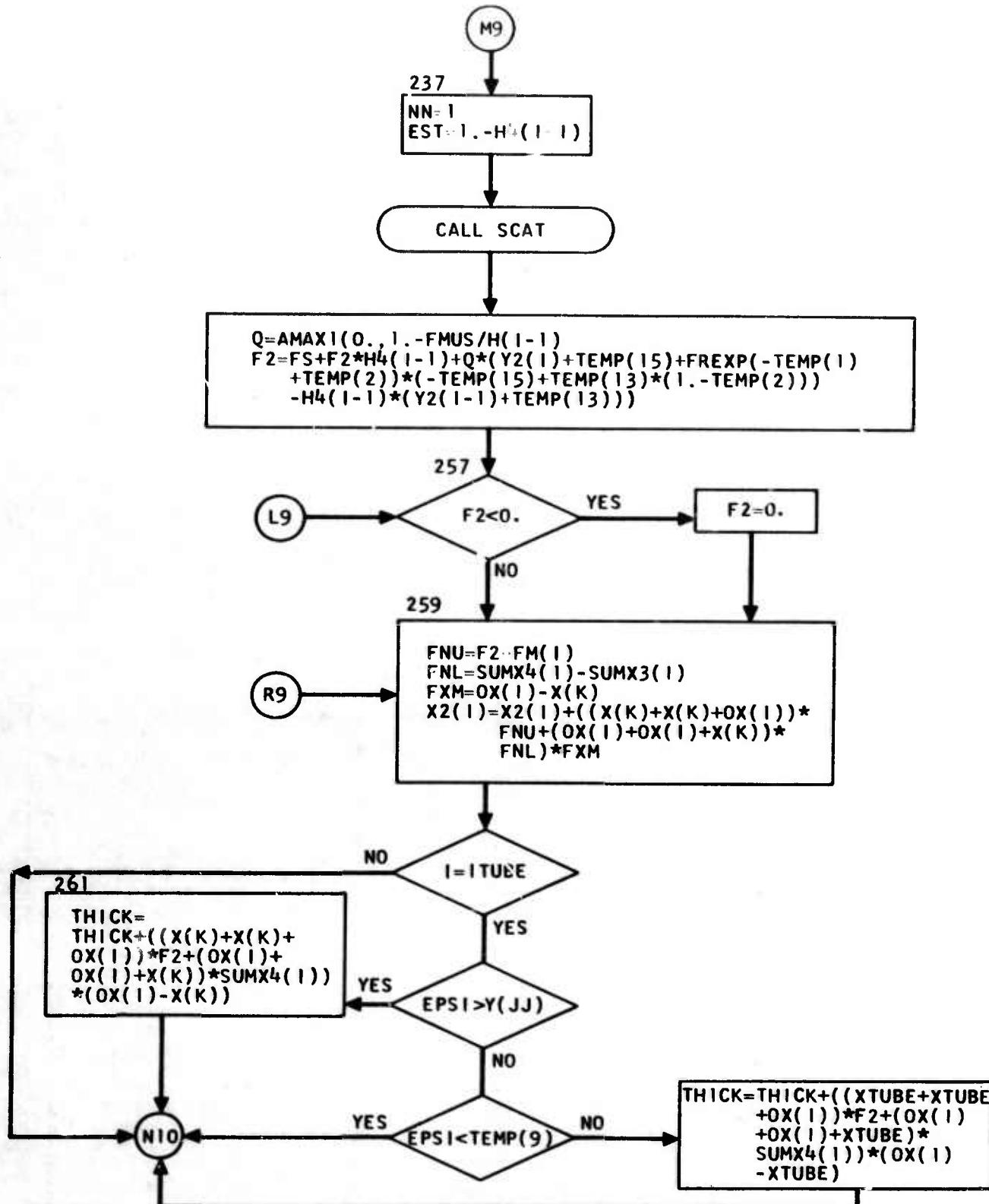


Figure 3 (continued). STRANS(N, M)

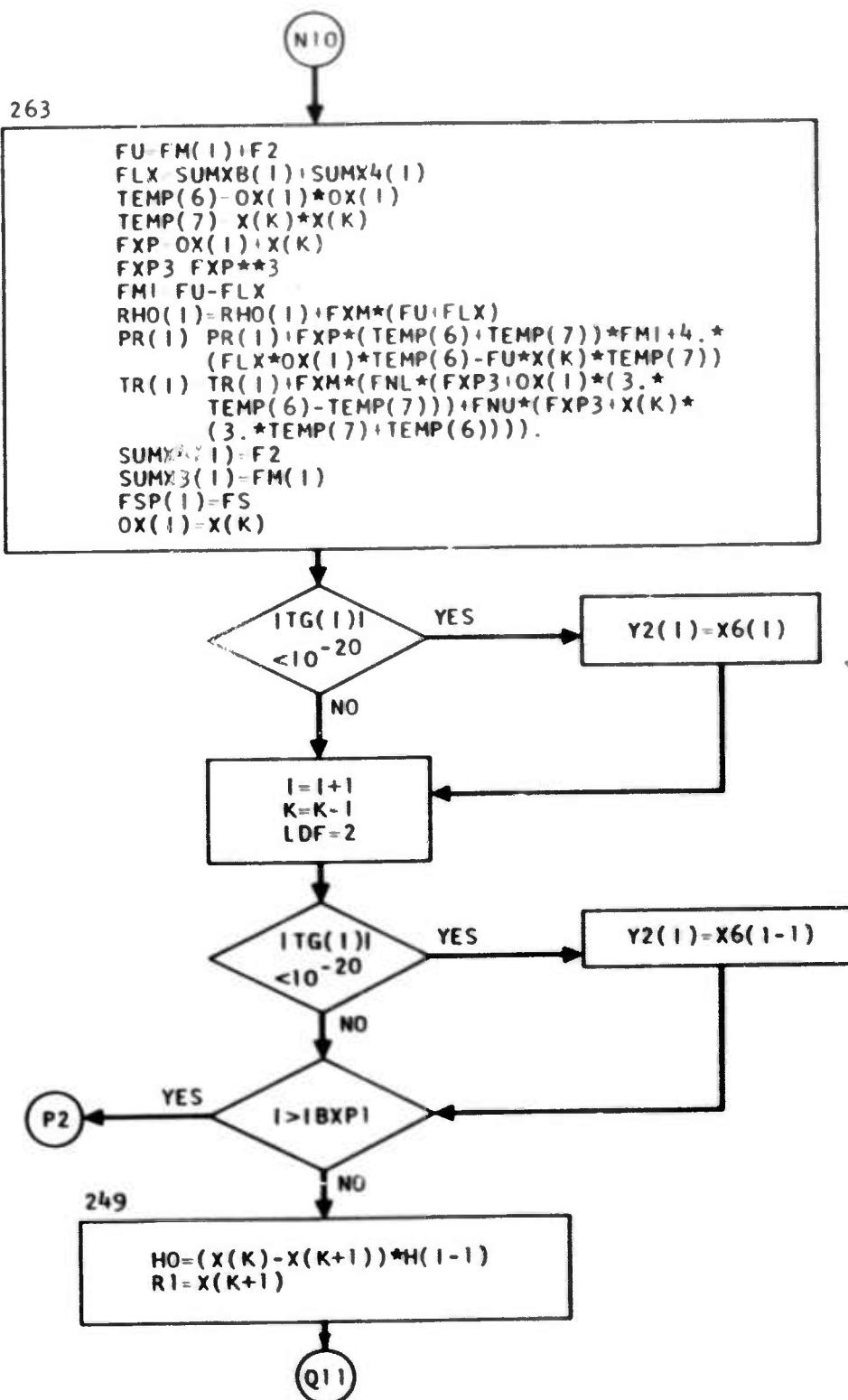


Figure 3 (continued). STRANS(N, M)

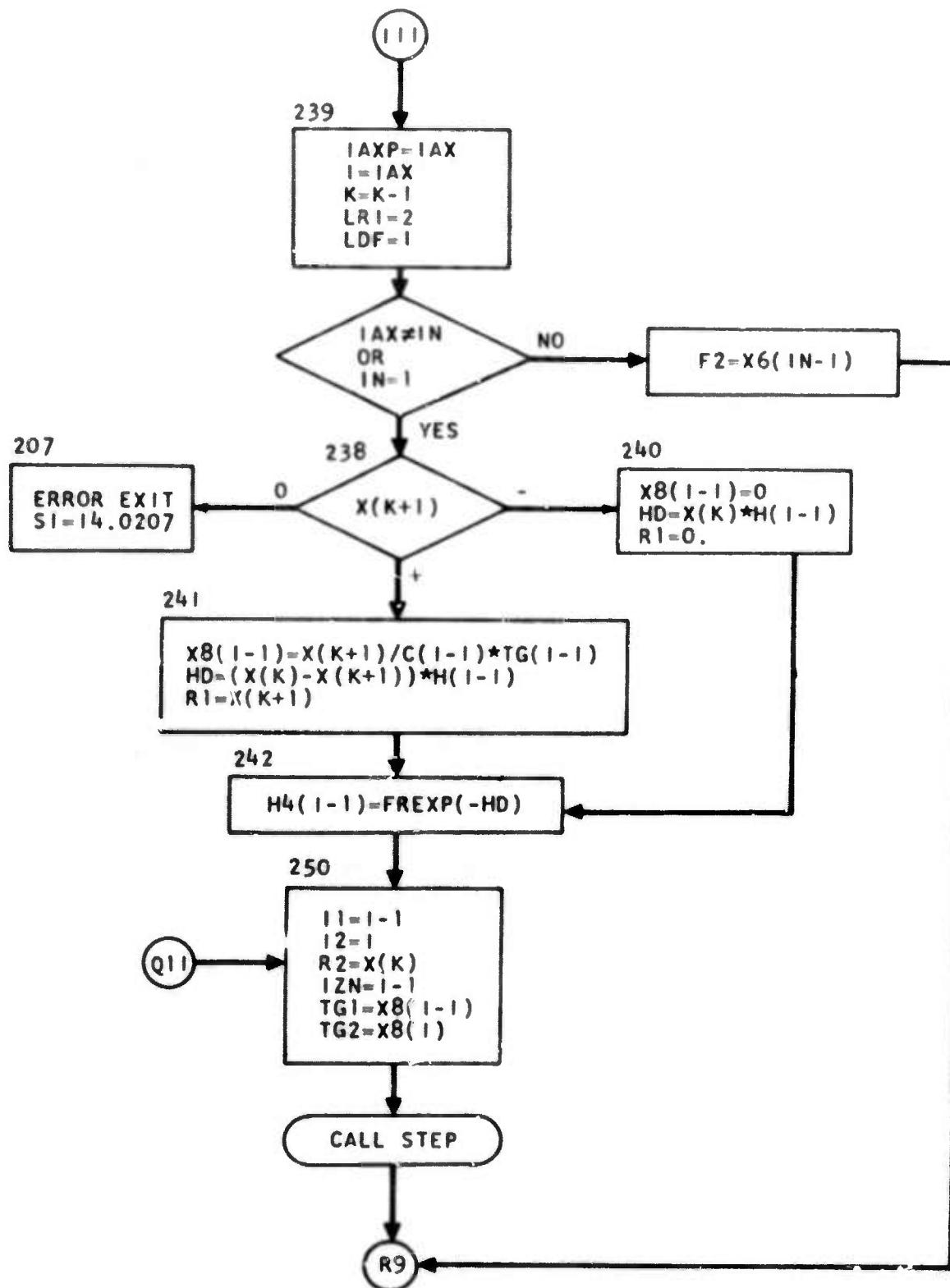


Figure 3 (concluded). STRANS(N, M)

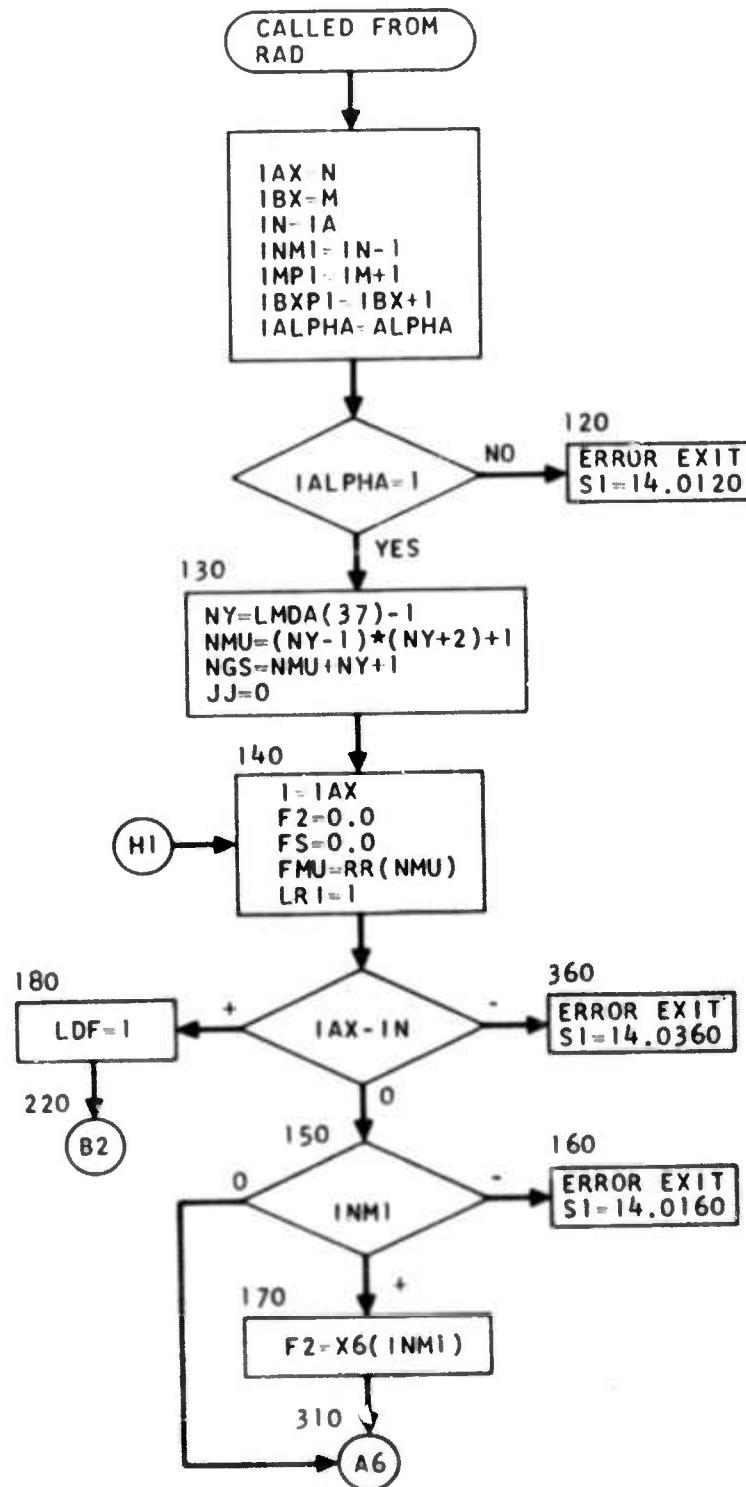


Figure 4. PTRANS

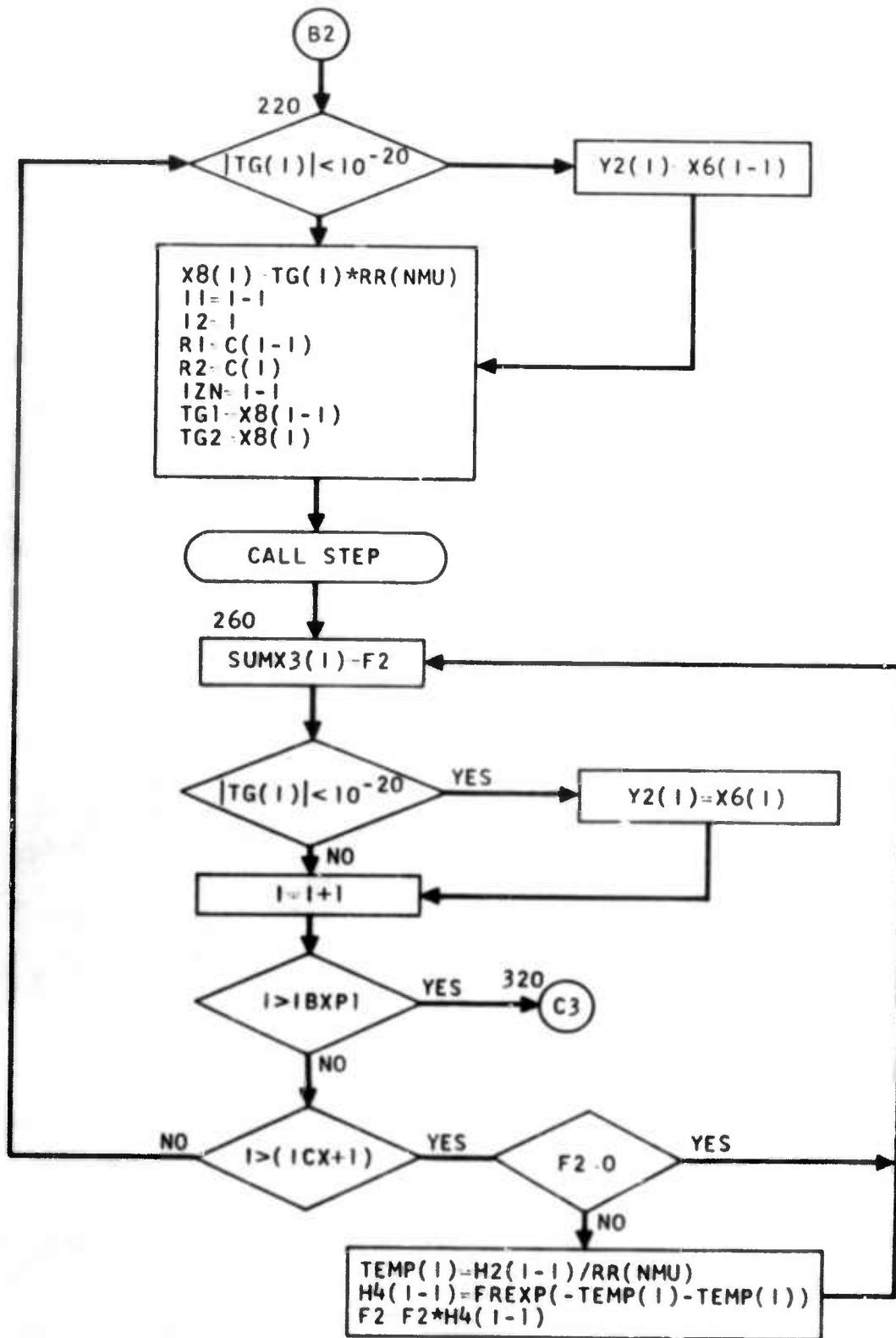


Figure 4 (continued). PTRANS

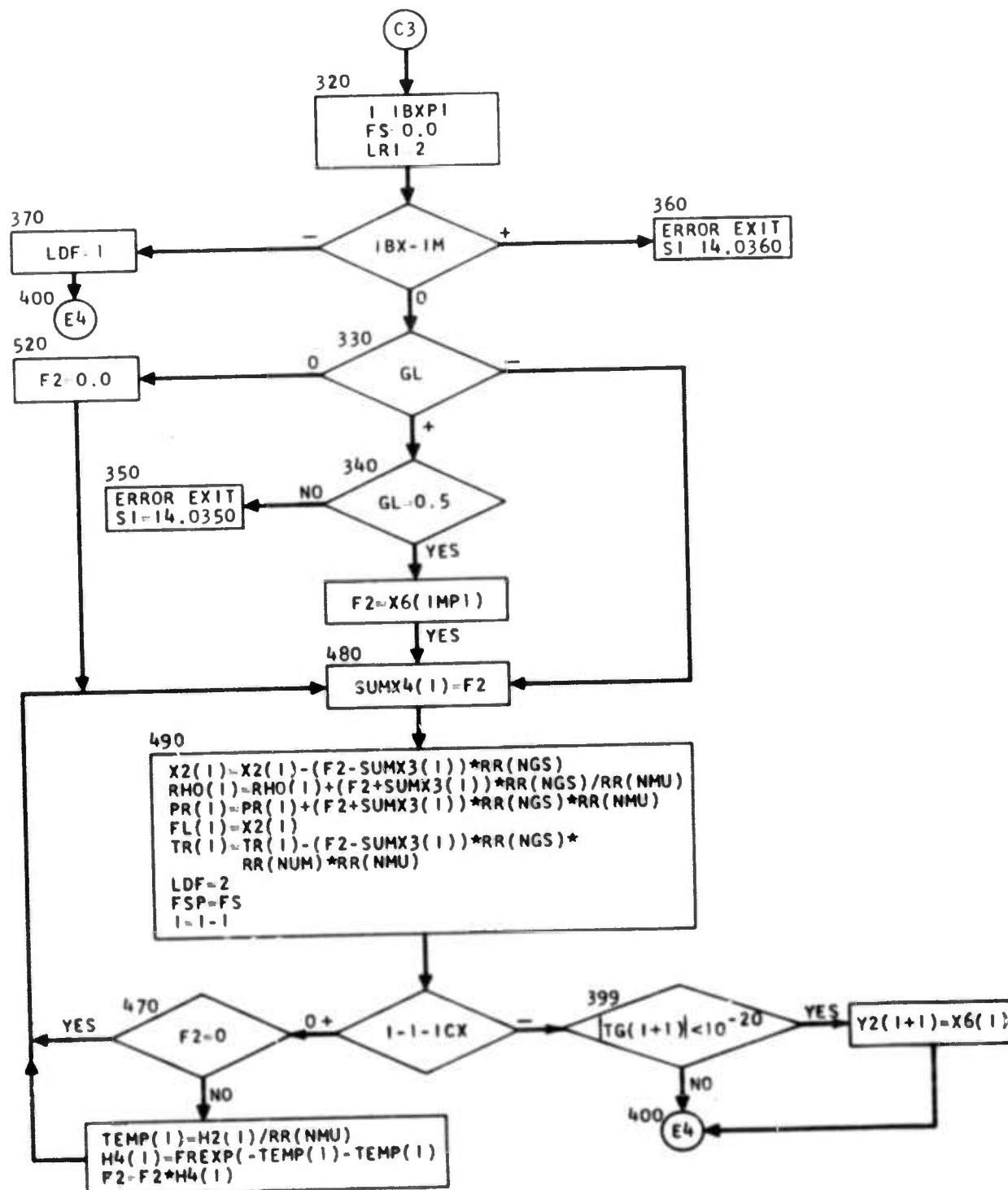


Figure 4 (continued). PTRANS

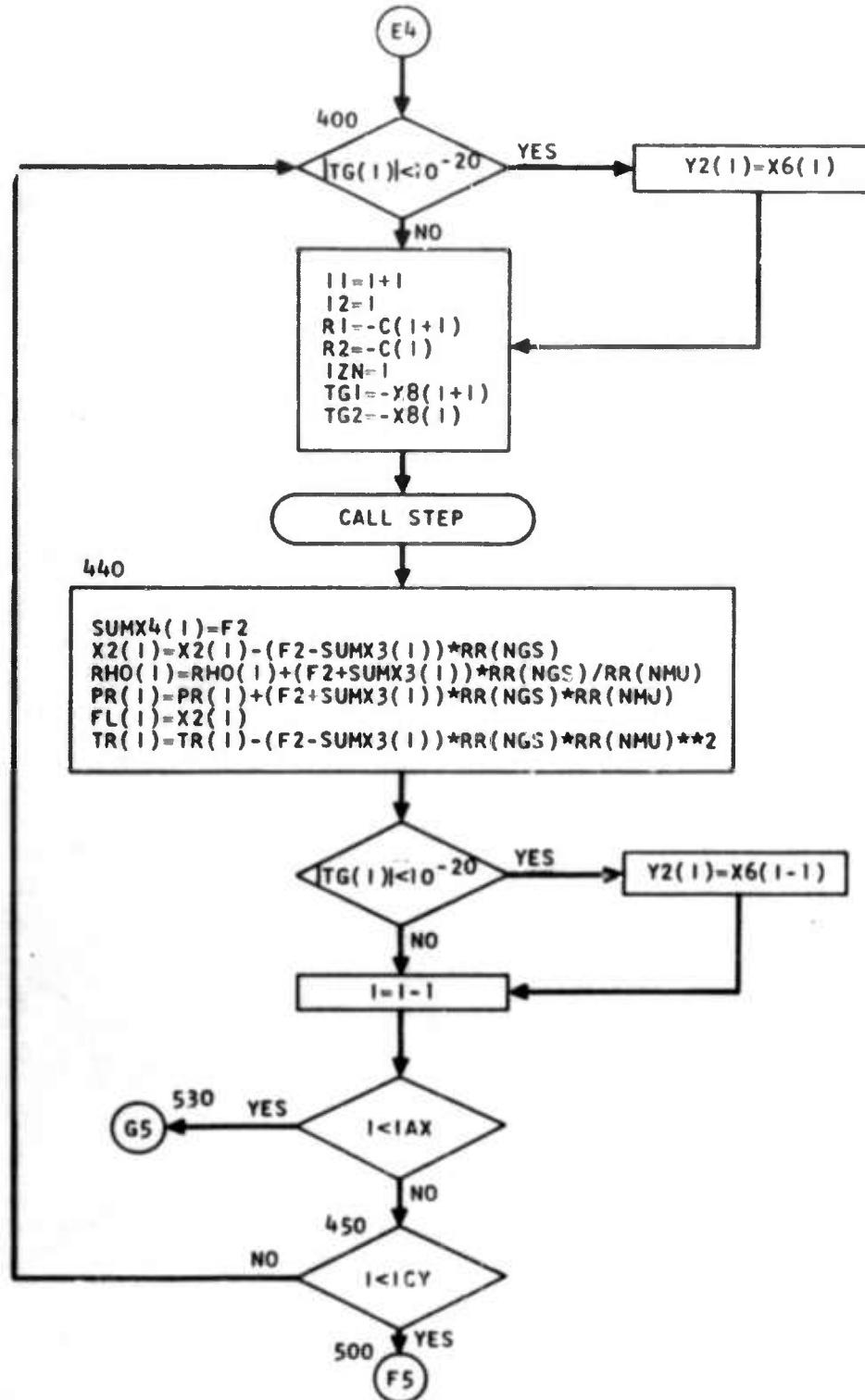


Figure 4 (continued). PTRANS

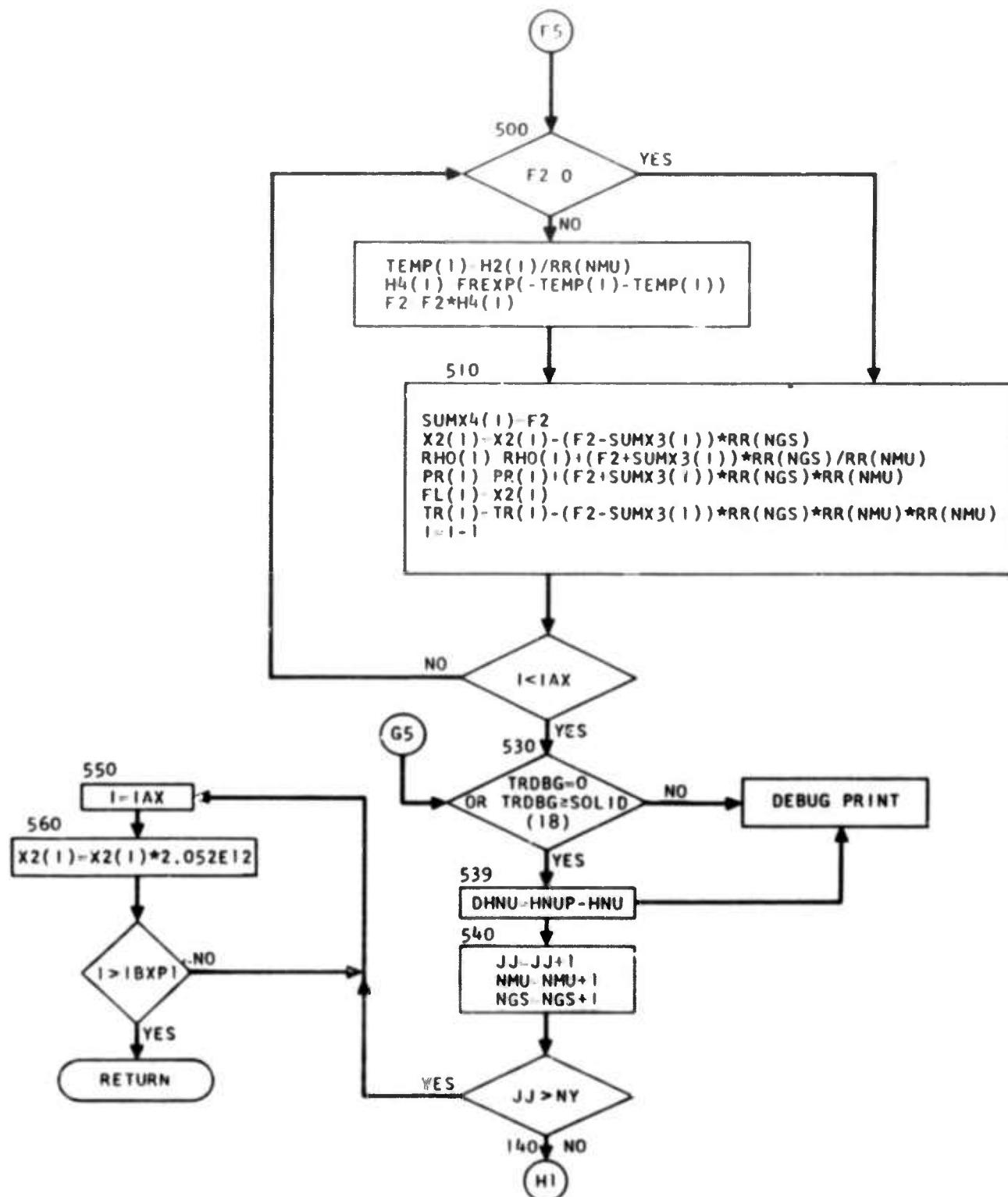


Figure 4 (continued). PTRANS

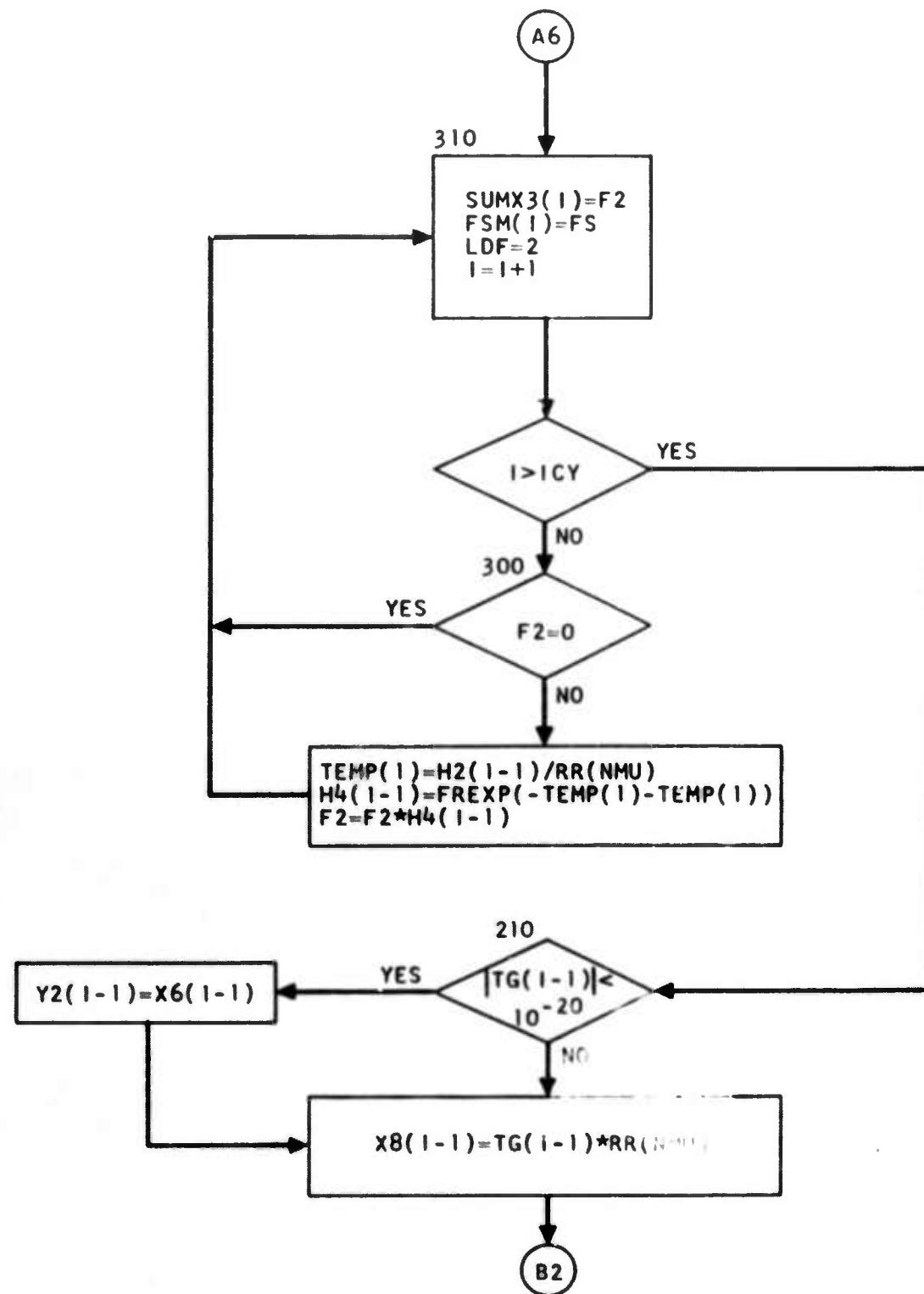


Figure 4 (concluded). PTRANS

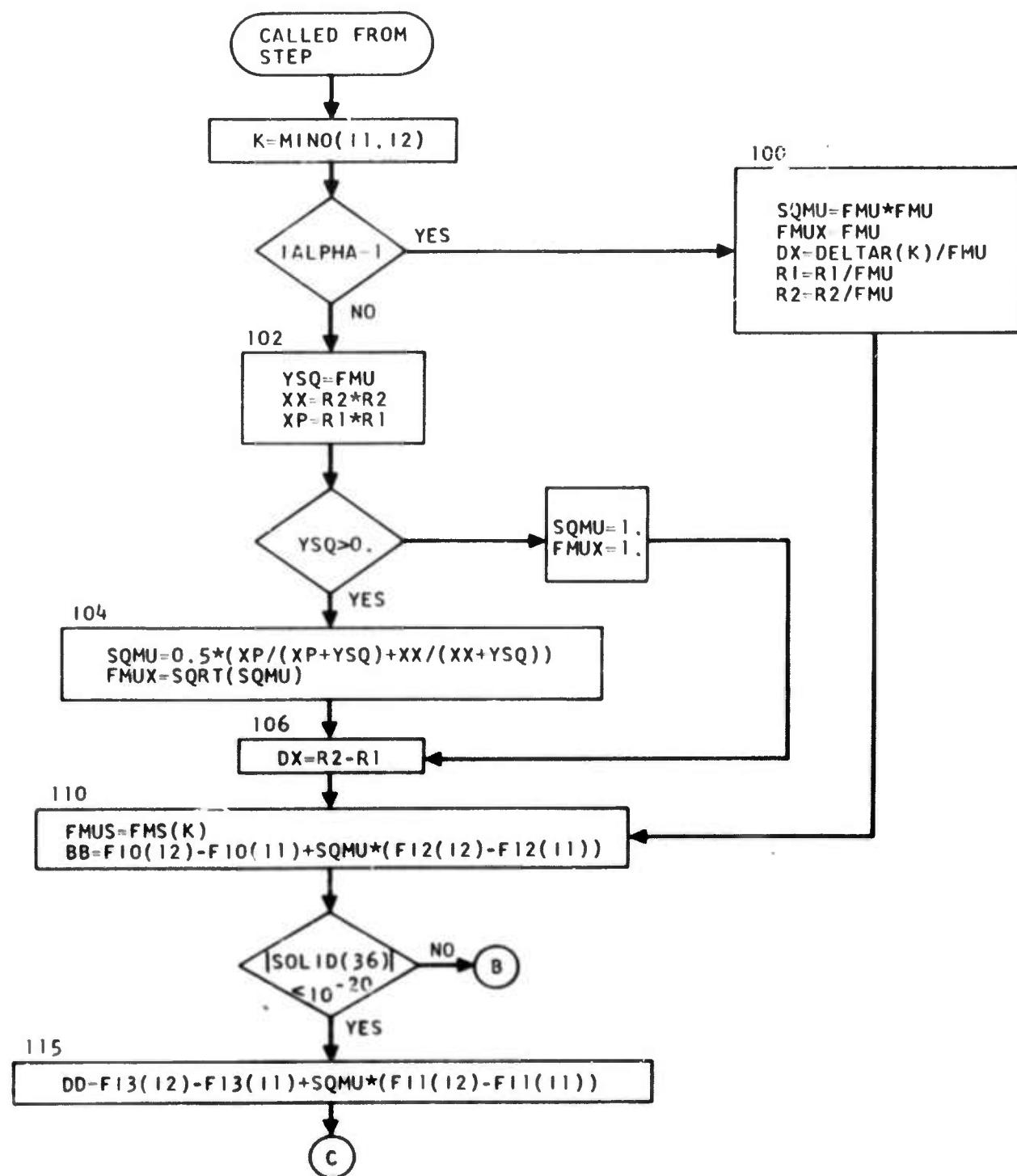


Figure 5. SCAT

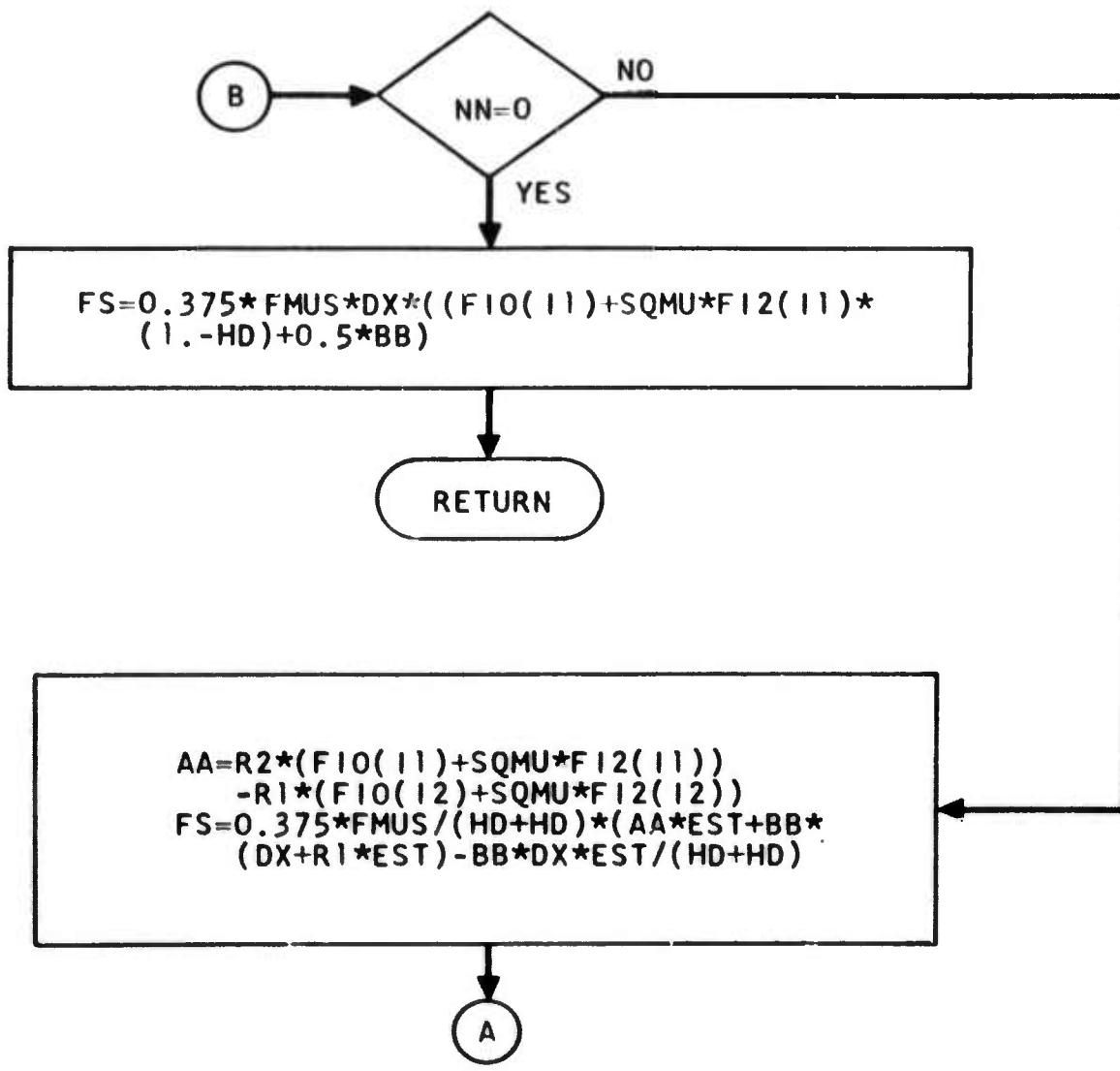


Figure 5 (continued). SCAT

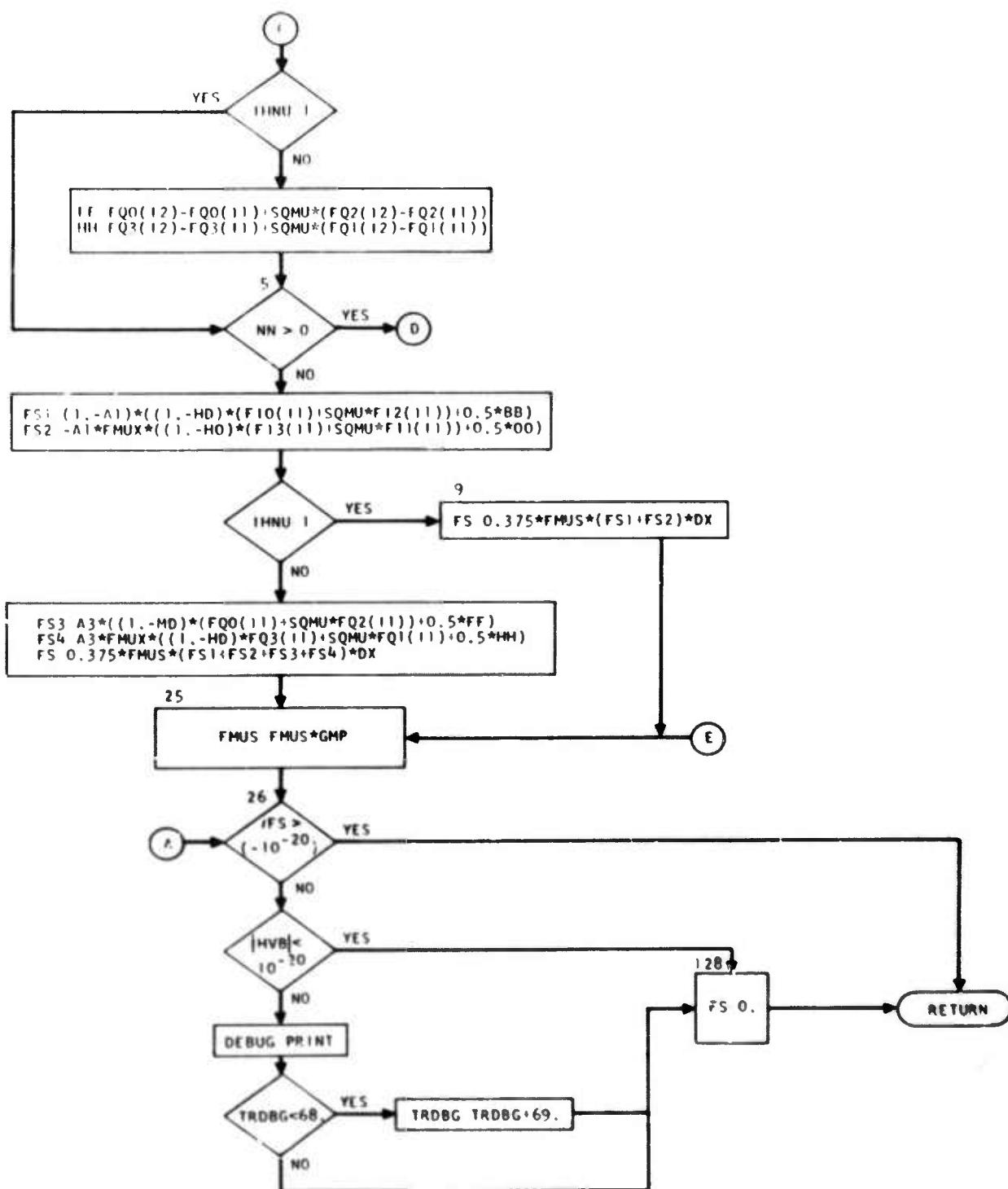


Figure 5 (continued). SCAT

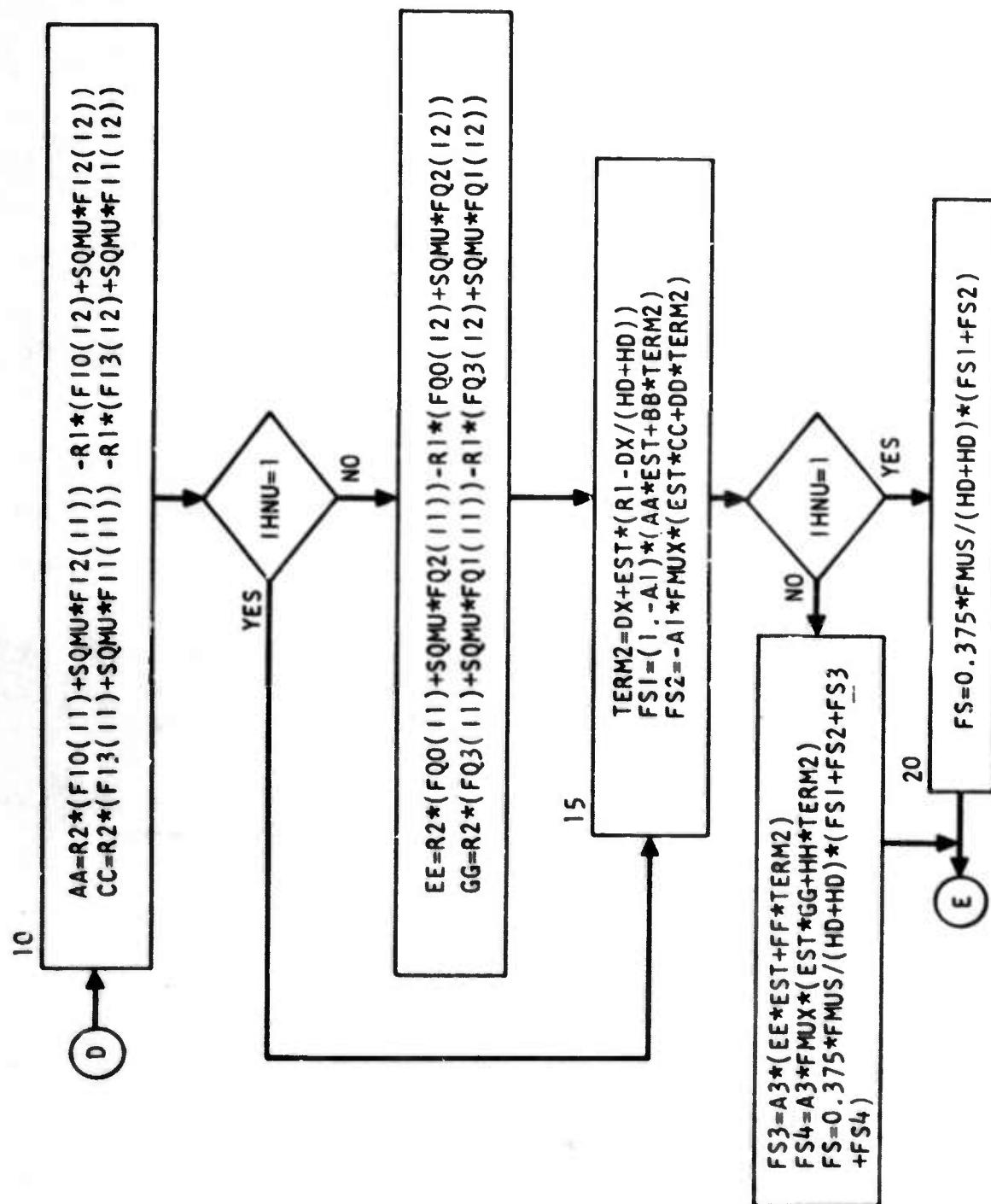


Figure 5 (concluded). SCAT

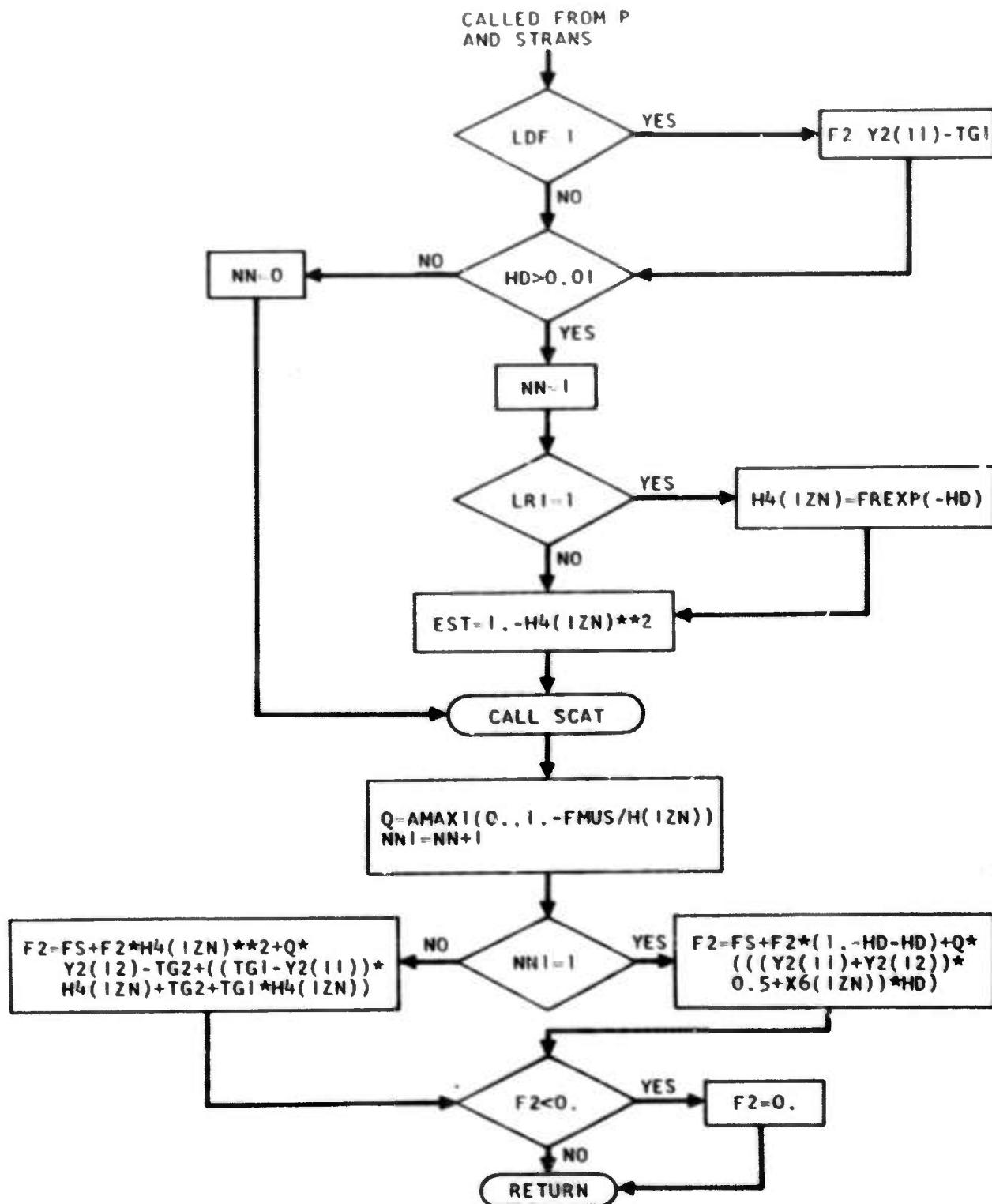


Figure 6. STEP

SCAT

AA	Following the notation of Eq. (30), AA is $A \cdot \Delta x$ as defined there. Private.
A1	This is defined in Eq. (31). It is evaluated in RAD. JIM Common.
A3	This is defined in Eq. (31). It is evaluated in RAD. JIM Common.
BB	This is $B \cdot \Delta x$ as defined in Eq. (30). Private.
CC	This is $C \cdot \Delta x$ as defined in Eq. (30). Private.
CVB	This input quantity, the negative intensity abort flag, is described on p. 87 of this report. Blank Common.
DD	This is $D \cdot \Delta x$ as defined in Eq. (30). Private.
DELTAR	Used in definition of DX in the plane case only. This variable is $r_{i+1}^{n+1} - r_i^n$ , evaluated in HYDRO. Since $r^n$ , the SPUTTER variable C, is used elsewhere in the radiation routines to define coordinates, it is recommended that the statement two lines below statement 100 be deleted and the GO TO 110 three lines below that be replaced by GO TO 106. This would give a more consistent, cheaper definition of DX. Blank Common.
DX	The distance from the initial to the final point of the current step along the characteristic ray, given in Eqs. (29) and (30), both as $\Delta$ and $\Delta x$ . Private.
EE	This is $E \cdot \Delta x$ as defined in Eq. (30). Private.
EST	This is $1 - e^{-\Delta\tau}$ , where $\Delta\tau$ is the optical depth of the current step. Evaluated in STEP and kept to avoid redundant calculation. JIM Common.
FF	This is $F \cdot \Delta x$ as defined in Eq. (30). Private
FI0	This is the mixture of moment quantities $3I_0 - I_2$ , as defined in Eq. (31). The last character of the variable is a zero. Although doubly indexed (frequency, zone) in Eq. (31), it is singly indexed (zone) in the OUTPUT code. Since the code calculates downward in frequency, the variables for the next upper (previous)

FI0 (continued)	frequency are available in FQ0, etc. The full set of FI0, etc., over frequency is stored on drum or disc, since there is no longer room in core storage for the doubly indexed array. PALMER Common.
FI1	This is the mixture of moment quantities $3I_1 - 5I_3$ , as defined in Eq. (31). Used only in Compton scattering. PALMER Common.
FI2	This is the mixture of moment quantities $3I_2 - I_0$ , as defined in Eq. (31). PALMER Common.
FI3	This is the mixture of moment quantities $3I_3 - 5I_1$ , as defined in Eq. (31). FI0, FI1, FI2, and FI3 are all evaluated in RAD. PALMER Common.
FMS	This zone array is $1/2 \mu_s$ , defined in Eq. (5), evaluated in RAD, using SOLID(37), an input quantity, for $\kappa_s$ . Equivalenced to SMLA in Blank Common.
FMU	This is a linkage variable that provides information about the characteristic line. In plane geometry, it is simply $ \mu $ , the absolute value of the cosine of the angle with respect to the normal. In spherical geometry, however, it is $y^2$ , the square of the impact parameter, evaluated once each y-line in STRANS. An approximate average value for $\mu$ is calculated in SCAT in this case. In both plane and spherical geometry, the derived parameter is the absolute value of the cosine. JIM Common.
FMUS	This variable, set to FMS(K) early in SCAT, is multiplied by GMP before being used in STEP to calculate an approximate $\mu_a / (\mu_a + \mu_s)$ . The multiplication by GMP is done only in the case of Compton scattering. It is not clear that this variable is needed as linkage. A slight speed gain is achieved at the cost of clarity. JIM Common.
FMUX	This is the angular variable $\mu$ defined beneath Eq. (1). It is unfortunate that the angular variable $\mu$ can be so easily confused with the photon absorption and scattering coefficients $\mu_a$ and $\mu_s$ , but it appears to be the common notation and is reflected in this computer program. Private.

- FQ0 This is the mixture of moment quantities  $3I_0 \dots I_2$  and corresponds to the FI0 ( $J + 1, x$ ) used in Eq. (30). That is, it is the FI0 zone array of the next higher (previously treated) frequency group. For the first group, these quantities are set zero in RAD and are not used in SCAT. It should be noted that the zero assumption is a particularly poor one as discussed on p. 82 of this report. The last character of FQ0 is a zero, and the array is in PALMER Common.
- FQ1 The array corresponding to FI1. PALMER Common.
- FQ2 The array corresponding to FI2. PALMER Common.
- FQ3 The array corresponding to FI3. PALMER Common.
- FS This is the scattering intensity, the "result" of executing the SCAT routine. It is the right-hand side of Eq. (29). JIM Common.
- FS1 The right-hand side of Eq. (29), which gives the detailed formulation of the scattering intensity, has four lines. Except for the factor  $\mu_g/\mu \cdot 3/16$ , FS1 is an intermediate term representing the first line. Private.
- FS2 Represents the second line of the same equation. Private.
- FS3 Represents the third line of the same equation. Private.
- FS4 Represents the fourth line of the same equation. Private.
- GG This is  $G \cdot \Delta x$  as defined in Eq. (30). Private
- GMP This is  $1 - 2y$ , mentioned on p. 82 of this report. It appears in Eqs. (19) ff. JIM Common.
- HD This is  $1/2 \kappa \rho \Delta x$ , where  $\Delta x$  is the geometrical length of the current step along the characteristic ray,  $\rho$  is the density, and  $\kappa$  is the photon coefficient  $\kappa_a + \kappa_g$ . Thus, HD, evaluated in PTRANS or STRANS and used in STEP and SCAT, is one-half the optical depth of the step. JIM Common.

HH	This is $H \cdot \Delta x$ as defined in Eq. (30). Private.
HVB	This input quantity, the negative scattering intensity debug print flag, is described on p. 87 of this report. Blank Common.
IALPHA	The geometry flag (1 = plane, 3 = sphere). Used in SCAT to obtain FMUX and SQMU, given FMU. Blank Common.
IHNU	The frequency group index, used in Compton scattering to branch to simpler coding for the first frequency group. LINDLY Common.
IMHAD	A divide check abort flag. Private.
I1	The index of the left-hand boundary of the current step. Evaluated in PTRANS or STRANS, used in STEP and SCAT. JIM Common.
I2	The corresponding index of the right-hand boundary. JIM Common.
K	The zone index of the current step, naturally the lesser of I1 and I2 provided that the characteristic ray at closest approach always is tangent to a zone boundary, as it is for the current standard SPUTTER and the OUTPUT code. Replacing K with IZN (see STEP) would remove this as a necessary condition, rendering SCAT safely generalizable. Private.
NN	This is a thick-thin flag set in STEP to allow SCAT to execute either Eq. (29) or Eq. (32), whichever is appropriate. JIM Common.
R1	The x-position of the left side of the step, not the radial position in spherical geometry. In STRANS, R1 and R2 are evaluated directly. In PTRANS, R1 and R2 are set to the slab coordinates and then adjusted for slant angle in SCAT. JIM Common.
R2	The x-position of the right side of the step. JIM Common.
SOLID(36)	The Compton switch described on p. 86 of this report. Blank Common.
SQMU	$\mu^2$ , where $\mu$ is given by the variable FMUX. Private.

S1            The SPUTTER error flag. See pp. 88-91 of this report. Blank Common.

TERM2        This is an intermediate quantity

$$\frac{x_2 - x_1 e^{-\alpha_i \Delta}}{\alpha_i} - \frac{1 - e^{-\alpha_i \Delta}}{\alpha_i^2}$$

appearing four times in Eq. (29). Private

TRDBG        This quantity, equivalenced to ACO3T4 in Blank Common, is the transport debug print flag. Zero for normal operation and no print, non-zero for transport debug print. It is changed in SCAT to trigger a print and then turned back in PTRANS or STRANS to forestall more prints, if HVB is set and a negative intensity is encountered.

XP            This is  $R1^{**2}$ , used to find a  $\mu$  in the spherical case. Private.

XX            This is  $R2^{**2}$ , used to find a  $\mu$  in the spherical case. Private.

YSQ           This is the square of the impact parameter, used to find  $\mu$  in the spherical case. Private.

### STEP

EST<sup>(1)</sup>

FMUS<sup>(1)</sup>

FS<sup>(1)</sup>

F2

This is the intensity on the right side of the step, the "result" of executing the STEP routine. It appears as  $I_j(x_2)$  on the left side of Eq. (27). JIM Common.

H

This is a zone array of  $\kappa\rho = \mu_a' + \mu_s$ , equivalenced to BIGB, evaluated in RAD, and used in the formation of Q. Blank Common.

HD<sup>(1)</sup>

H4	This is a zone array of $e^{-1/2\Delta\tau}$ , where HD (see SCAT) is $1/2\Delta\tau$ . It is formed in STEP if one is proceeding inward in spherical geometry or forward in slab geometry and is available in the reverse case. It is equivalenced to SMLH in Blank Common.
IZN	This is the index of the zone being traversed in the current step along the characteristic ray. Evaluated in PTRANS or STRANS, used in STEP. JIM Common.
$I_1^{(1)}$	
$I_2^{(1)}$	
LDF	In subsection 2.1.5 of reference 2, the three initial boundary conditions for $I_{i-1}$ (initial value of F2) are described. For the first two of these, F2 is defined in PTRANS or STRANS, and for the general case, F2 is left over from the previous step. However, for the third case, diffusion, the initial intensity is given by Eq. (2.19) of reference 2, and this is executed in STEP. If LDF is 1, the diffusion boundary condition is applied. If LDF is 2, F2 is assumed to be properly initialized. LDF, in JIM Common, is evaluated in PTRANS or STRANS.
LRI	This is the left-right index, used to decide whether to evaluate H4. It is set in STRANS or PTRANS. JIM Common.
NN <sup>(1)</sup>	
NN1	This is NN + 1, set to allow a computed GO TO on the thick-thin flag. A code reform would increase NN by 1 in the several places in PTRANS and STRANS where it is set, and would eliminate the need for a second variable. Private.
Q	This is, or should be, $\mu_a / (\mu_a + \mu_s)$ . In the current, rather crude calculation of this quantity, Q is set zero in case the value should go negative. Private.
TG1	This is the slant source gradient along the characteristic ray, evaluated at the left side of the current step, R1. This quantity appears as $\mu_{i-1} (\partial B / \partial h) _{i-1}$ in Eq. (2.19) of reference 2. JIM Common.
TG2	This is the corresponding quantity evaluated at the right side, R2. JIM Common.

- X6 This is a zone array of the source  $B_j$  in Eq. (24), evaluated at the zone center. Calculated in RAD. Blank Common.
- Y2 This is a zone array of the source evaluated at the zone boundary according to criteria discussed in subsections 3.2 and 5.1.3 of reference 2 (see also the model in subsection 2.1.1 of reference 2). Equivalenced to X5 in Blank Common, evaluated in RAD, and adjusted in STRANS if TG (see STRANS) is zero.

STRANS

Variables defined in the STEP or SCAT lists are given first.

DELTAR <sup>(1)</sup>	I2 <sup>(1)</sup>
EST <sup>(1)</sup>	LDF <sup>(2)</sup>
FMU <sup>(1)</sup>	LRI <sup>(2)</sup>
FMUS <sup>(1)</sup>	NN <sup>(1)</sup>
FS <sup>(1)</sup>	R1 <sup>(1)</sup>
F2 <sup>(2)</sup>	R2 <sup>(1)</sup>
H <sup>(2)</sup>	S1 <sup>(1)</sup>
HD <sup>(1)</sup>	TG1 <sup>(2)</sup>
H4 <sup>(2)</sup>	TG2 <sup>(2)</sup>
IALPHA <sup>(1)</sup>	TRDBG <sup>(1)</sup>
IHNU <sup>(1)</sup>	X6 <sup>(2)</sup>
IZN <sup>(2)</sup>	Y2 <sup>(2)</sup>
I1 <sup>(1)</sup>	
ALPHA	The geometry indicator in real form (see IALPHA <sup>(1)</sup> ). SPUTTER has both real and integer forms for this variable. Blank Common.
C	The "old" space variable, $r_i^n$ , used throughout the radiation codes. Blank Common.
CNT1	The updated cycle number, evaluated only if a transport debug print is called. Private.

CSQD	Evaluated in RAD and used there and in STRANS. This is a zone array of the square of the coordinate, $C^* * 2$ , and is provided to save multiplications within frequently executed loops. It probably saves relatively little time and can be considered one of the more expendable arrays. Equivalenced to CRTC in Blank Common.
C1	For $IAX > 1$ , there is an interior solid or diffusion region. Even if this consists of many zones, one requires, nevertheless, only a few y-lines. As stated in subsection 2.2.2 of reference 2, "If a diffusion region having an outer boundary $r_D$ exists inside the transport region, y-lines are placed as near as possible to $0.5 r_D$ , $0.75 r_D$ , and $r_D$ penetrating the diffusion region." C1 is set first to $0.5 r_D$ and then to $0.75 r_D$ to provide this placement. Private.
DHNU	The width, in eV, of the current frequency group. Set in STRANS and it is probably not used. LINDLY Common.
EPSI	Radius of OUTPUT sample tube. Blank Common.
FL	A separate array of the first moment of the intensity, $\int_{-1}^1 I \mu d\mu$ , normalized differently from X2 and used in the evaluation of the scattering moment quantities. Equivalenced to SMLB in Blank Common.
FLX	An intermediate quantity, $I_- + I_+$ along the previous y-line, used in calculating the intensity moments. Private.
FM	The array of $I_-$ along the current y-line. Four intensities are required to calculate the moment quantities for a particular zone boundary and y-band. One can be used immediately after being calculated; the others have to be available in zone array storage. Equivalenced to ER in Blank Common.
FM1	The intermediate quantity $(I_+ + I_-)$ along the current y-line minus $(I_+ + I_-)$ along the previous y-line. Private.
FNL	The intermediate quantity $(I_+ - I_-)$ along the previous y-line. Private.
FNU	The intermediate quantity $(I_+ - I_-)$ along the current y-line. Private.

FP	Evaluated and used in the "top slice" section. this intermediate quantity is $I_- + I_+ + I_T$ , where $I_T$ is an interpolated value at the top of the arc. See subsections 2.3.2 and 5.2.5 of reference 2. Private.
FPL	This is $2(I_- + I_+) + I_T$ . See discussion of FP. Private.
FSM	The zone array of scattering intensity on the inward sweep. Stored for display in the debug print. Equivalenced to SMLD in Blank Common.
FSP	The corresponding array of scattering intensity for the outward sweep. Equivalenced to SMLE in Blank Common.
FU	$I_T$ . See discussion of FP. Private.
FXM	On the zone boundary, the x of the previous y-line minus the x on the current y-line. Private.
FXP	On the zone boundary, the x of the previous y-line plus the x on the current y-line. Private.
FXP3	The cube of FXP, used twice in evaluating the third moment of the intensity. Private.
GL	If positive, there is a blackbody exterior to the radiation region. The initial intensity is set accordingly. Blank Common.
HNU	The frequency (eV) of the lower limit of the current frequency group. Set to $10^{-3}$ for a grey problem. LINDLY Common.
HNUP	The frequency (eV) at the upper limit of the current frequency group. LINDLY Common.
H2	Depending on SOLID(10), an input quantity, H2 is a zone array of one-half either the Planck or Rosseland optical depth. Planck if SOLID(10) is zero; otherwise, Rosseland. Used to define optical depth in the transport routines. Evaluated in RAD. Equivalenced to EC in Blank Common.
I	Generally used in STRANS to denote the zone or zone-boundary index. Private.
IAX	The left (interior) zone limit of the current transport subregion, inside of which may lie a diffusion region, an interior blackbody, or the center. Private, but linked to RAD by the first argument of the subroutine.

IAXP	An index denoting the lower limit of zone boundaries for which a top-slice calculation is to be made. Private.
IAXP1	The lower limit of zone boundaries to be treated in the transport debug print for a given y-line; this is the maximum of IAX, IAXP-1. Private.
IBX	The right (exterior) zone limit of the current transport subregion. Private, but linked to RAD by the second argument.
IBXP1	This is IBX+1 and is the upper zone-boundary limit of the current transport subregion. Private.
IM	The right (exterior) zone limit of the whole radiation region, set in RAD. LINDLY Common.
IN	The left (interior) zone limit of the whole radiation region, set in RAD. LINDLY Common.
IT	A running index used in setting up y-lines; its function is probably obsolete. Private.
ITUBE	Index of zone boundary at which OUTPUT sample tube is affixed. Set to LMDA(26), the input quantity, near the beginning of STRANS. Used only because subscripted subscripts are illegal in FORTRAN. Private.
J	Used as the running index of the top-slice DO loop, where I is a limit of the loop. Private.
JJ	The y-line index. Private.
JJJ	Set to JJ throughout, its function as distinct from JJ is to be found in the coding given in reference 2 (since deleted). It was involved in a y-line skipping procedure, now prohibited by scattering considerations. A code cleanup would remove JJJ. Private.
K	This is the index for the variable X. Private.
KK	A counter updated from the X array which gives a "master index" for a given y-line. Private.
KKK	An index for X used in the transport debug print. Private.

KOOOFX	A dummy argument used in CALL DVCHK, analogous to IMHAD <sup>(1)</sup> . When codes were processed from FORTRAN II to FORTRAN IV by SIFT, the dummy variable KOOOFX was invented, where the middle characters were zeros. By historical accident, the characters in this variable are letter O's. A code cleanup would change this to KX. Private.
KX	A dummy argument used in CALL DVCHK. Private.
LMDA(26)	The input quantity that defines ITUBE. Blank Common.
NY	The number of y-lines. LINDLY Common.
OX	A zone array storing X for the previous y-line. Since all the X are available, OX could be dispensed with if an index analogous to K but for the previous y-line were defined. Equivalenced to W in Blank Common.
PR	A zone array for the second moment of intensity, $\int_{-1}^1 I\mu^2 d\mu$ . Used in calculating the scattering moment quantities. Equivalenced to X7 in Blank Common.
Q	Same meaning as Q <sup>(2)</sup> , but different storage. Private.
RHO	A zone array for the zeroth moment of intensity, $\int_{-1}^1 I d\mu$ . Used in calculating the scattering moment quantities and, in RAD and elsewhere, as radiation energy density. Blank Common.
SOLID(18)	The current cycle number. Blank Common.
SUMX3	A zone array for I <sub>-</sub> along the previous y-line. See discussion of FM. Equivalenced to CHIC in Blank Common.
SUMX4	A zone array for I <sub>+</sub> along the previous y-line. See discussion of FM. Equivalenced to BC in Blank Common.
TEMP(1)	In SPUTTER, the TEMP array is used by many subroutines for scratch storage. Liberal use of it was made in STRANS, especially for "top-slice" coding. (See subsections 2.1.2, 2.3.2, and 5.2.5 of reference 2.) All the TEMP variables are in Blank Common. TEMP(1) is used to give half the optical depth of a top slice, or the full optical depth of one side of it.
TEMP(2)	This is $\Delta_{1/2}$ (defined in subsection 2.1.2 of reference 2).

TEMP(5)	This is $r_{IBXP1}^2 - y_j^2$ , where $j$ is the index of the last y-line used. Intermediate quantity in calculating intensity moments in a top-slice calculation.
TEMP(6)	$x^2$ of a point on the previous y-line.
TEMP(7)	The corresponding $x^2$ for the current y-line. Also used for an entirely different purpose, $e^{-TEM(2)}$ in a top-slice calculation.
TEMP(9)	Used as the y-value of the last y-line treated, which need not be $Y(JJ-1)$ , since lines may be skipped.
TEMP(11)	This is $r_i^2 - y_j^2$ , used in calculating FU in a top slice.
TEMP(13)	A special top-slice source gradient, F, defined in Eq. (2.13), subsection 2.1.2, reference 2.
TEMP(15)	G defined in Eq. (2.13), subsection 2.1.2, reference 2.
TEMP(16)	In the top-slice thin approximation, this is the last term of Eq. (2.15), subsection 2.1.2, reference 2.
TG	The source gradient, defined in RAD. It may be set zero depending on conditions discussed in subsection 5.1.3 of reference 2. See also table II in this report. This is a zone array equivalenced to V in Blank Common.
THICK	This is the "OUTPUT-output," in gross outward flux down a tube of specified radius at specified zone boundary. LINDLY Common.
TR	A zone array for the third moment of intensity, $\int_{-1}^1 I\mu^3 d\mu$ . Used in calculating the scattering moment quantities. Equivalenced to SMLC in Blank Common.
X	This is an array of all the values of $\sqrt{r_i^2 - y_j^2}$ for all y-lines set up early in RAD. The X array is evaluated in RAD, outside the frequency loop, in order to save taking thousands of square roots within the frequency loop, say, in STRANS. For each y-line, in addition to the set of x values, there are stored in the X array the number of x values and the negative of $y^2$ . DAVIS Common.
XS	In the final top-slice calculation, this variable is $r_{IBXP1+1}^2 - r_{IBXP1}^2$ . Used to calculate an initial diffusion intensity. Private.
XSQ	For a given cell boundary, this is $x^2$ for the previous y-line. Used as an intermediate quantity in evaluating moments. Private.

X TUBE	This is $\sqrt{r_{ITUBE}^2 - EPSI^2}$ , the x-value of the intersection of the OUTPUT sample tube with the zone boundary at which the tube is affixed. Private.
X2	This is the first moment of the intensity $\int_{-1}^1 I\mu d\mu$ , normalized before exiting STRANS to serve as the radiation flux for the calculation of heating rates. Blank Common.
X8	This is a zone array of the slant source gradient, $TG \cdot x/r$ . Equivalenced to X4 in Blank Common.
Y	This is the array of values of y, the impact parameter, for y-lines. Equivalenced to BIGA in Blank Common.
YSQDP	This is convenient storage for $y^2$ of the current y-line. Private.
YSQD1	This is convenient storage for $y^2$ of the previous y-line. Private.

PTRANS

Variables defined in the STEP, SCAT, or STRANS lists are given first.

ALPHA <sup>(3)</sup>	H2 <sup>(3)</sup>	KOOOFX <sup>(3)</sup>	TG1 <sup>(2)</sup>
C <sup>(3)</sup>	H4 <sup>(2)</sup>	LDF <sup>(2)</sup>	TG2 <sup>(2)</sup>
DHNU <sup>(3)</sup>	I <sup>(3)</sup>	LRI <sup>(2)</sup>	TR <sup>(3)</sup>
FL <sup>(3)</sup>	IALPHA <sup>(1)</sup>	PR <sup>(3)</sup>	TRDBG <sup>(1)</sup>
FMU <sup>(1)</sup>	IAX <sup>(3)</sup>	RHO <sup>(3)</sup>	X2 <sup>(3)</sup>
FS <sup>(1)</sup>	IBX <sup>(3)</sup>	R1 <sup>(1)</sup>	X6 <sup>(2)</sup>
FSM <sup>(3)</sup>	IBXP1 <sup>(3)</sup>	R2 <sup>(1)</sup>	X8 <sup>(3)</sup>
FSP <sup>(3)</sup>	IM <sup>(3)</sup>	SOLID <sup>(18)</sup> <sup>(3)</sup>	Y2 <sup>(2)</sup>
F2 <sup>(2)</sup>	IN <sup>(3)</sup>	SUMX3 <sup>(3)</sup>	
GL <sup>(3)</sup>	IZN <sup>(2)</sup>	SUMX4 <sup>(3)</sup>	
HNU <sup>(3)</sup>	I1 <sup>(1)</sup>	S1 <sup>(1)</sup>	
HNUP <sup>(3)</sup>	I2 <sup>(1)</sup>	TG <sup>(3)</sup>	

IA	The left-hand (interior) zone index of the SPUTTER vapor region. Since this may not correspond to the left limit of the radiation region, IN (defined in RAD) should be used and not redefined (perhaps erroneously) in PTRANS. Code revision is needed here. Blank Common.
ICX	The right-hand zone limit of region with source. Set to IM in RAD, it is retained solely to procrastinate on cleanup in PTRANS. DAVIS Common.
ICY	The left-hand zone limit of region with source. See ICX. DAVIS Common.
IMP1	IM+1. See IM <sup>(3)</sup> . Private.
INM1	IN-1. See IN <sup>(3)</sup> . Private.
JJ	The index of a characteristic line, analogous to the y-line index JJ <sup>(3)</sup> . Maximum legitimate value is 5. Private.
JJJ	JJ+1, used in the transport debug print. Private.
LMDA(37)	The input quantity specifying the number of characteristic lines to be used (a maximum of 6). Blank Common.
NGS	The index of Gauss weights. See RR. Private.
NMU	The index of angles. See RR. Private.
NY	The upper limit to JJ, set to LMDA(37)-1. LINDLY Common.
RR	The plane transport calculation is a scheme of evaluating intensities along characteristic lines and integrating these (forming the moments) by the "double Gaussian" method, where special weighting quantities corresponding to the angular intervals are stored. A table of average angles and corresponding Gauss weights for 2, 3, 4, 5, and 6 angular intervals is stored by DATA statement in RR and used in the finite sum formulation in the code. There are 40 entries. Private.
TEMP(1)	This is $1/2 \rho \kappa \Delta x$ , one-half the slant optical depth, calculated in the case of no source, an option now eliminated. See ICX, ICY, and p. 79 of this report. A code cleanup would delete reference to TEMP(1). Blank Common.

RAD

Variables defined in SCAT, STEP, STRANS, or PTRANS are given first.

$A_1^{(1)}$	$FQ0^{(1)}$	$IBX^{(3)}$	$SOLID(36)^{(1)}$
$A_3^{(1)}$	$FQ1^{(1)}$	$ICX^{(4)}$	$S_1^{(1)}$
$C^{(3)}$	$FQ2^{(1)}$	$ICY^{(4)}$	$TG^{(3)}$
$CNT_1^{(3)}$	$FQ3^{(1)}$	$IHNU^{(1)}$	$THICK^{(3)}$
$CSQD^{(3)}$	$GMP^{(1)}$	$IM^{(3)}$	$TR^{(3)}$
$DHNU^{(3)}$	$H^{(2)}$	$IMP_1^{(4)}$	$X^{(3)}$
$FI0^{(1)}$	$HNU^{(3)}$	$IN^{(3)}$	$X_2^{(3)}$
$FI1^{(1)}$	$HNUP^{(3)}$	$INM_1^{(4)}$	$X_6^{(2)}$
$FI2^{(1)}$	$H_2^{(3)}$	$KX^{(3)}$	$Y^{(3)}$
$FI3^{(1)}$	$I^{(3)}$	$PR^{(3)}$	$Y_2^{(2)}$
$FL^{(3)}$	$IALPHA^{(1)}$	$RHO^{(3)}$	
$FMS^{(1)}$	$IAX^{(3)}$	$SOLID(18)^{(3)}$	

<b>A</b>	The SPUTTER area term, $\alpha r^{\alpha-1}$ , a zone array, used in calculating the radiation flux at the external boundary for either diffusion or no vapor. Blank Common.
<b>A1P</b>	This is the intermediate quantity $h\nu_j^2/m_e C^2 \Delta\nu_j$ in Eq. (31). Used in forming A1. Private.
<b>BETA</b>	The quantity $h\nu/\theta$ , often represented by the variable u. In this case, $h\nu$ is the minimum photon energy for the current frequency group. Private.
<b>BETAP</b>	The quantity $h\nu/\theta$ , where $h\nu$ is the maximum photon energy for the current frequency group. Private.
<b>BLANK3</b>	Used in radiation supercycling. If the time step calculation in RAD calls for supercycling, BLANK3 is set to the current time plus whatever is allowed for the supercycle. Blank Common.

CAPAC	$\kappa_{\text{Planck}}$ , evaluated in KAPPA or one of its subroutines and used in RAD. CAPAC(150-152) is reserved for special input quantities. See p. 86 of this report. Blank Common.
CAPAR	$\kappa_{\text{Rosseland}}$ , evaluated in KAPPA or one of its subroutines and used in RAD. CAPAR(121-150) used for up to 15 frequencies of "OUTPUT-output" flux rates and fluxes accumulated over time. This rather clumsy storage mechanism sets an upper limit of 120 zones and 15 frequencies for problems run on the OUTPUT code.
CNTMAX	This is the SPUTTER cycle limit and is tested so that, assuming one wants a multifrequency print with the regular SPUTTER print, a print is obtained for the last cycle. Blank Common.
CPA	This is the Rosseland $\kappa$ with scattering, adjusted for Compton scattering if appropriate. Used in forming the optical depth arrays. Private.
CPB	This is the Planck $\kappa$ with scattering, used in forming the H(2) and H2(3) arrays if SOLID(10) is zero. Private.
CPC	This is the Rosseland $\kappa$ with scattering, not adjusted for Compton scattering. Private.
CV	This is the array of $C_v$ , specific heat, calculated elsewhere in SPUTTER and used in RAD to formulate the stability time step. Blank Common.
DFB	The portion of the area of the Planck function, normalized to 1, occupied by the current frequency group at a given temperature. Private.
DHNUP	The width of the previous (next higher) frequency group, used in calculating A3 <sup>(1)</sup> . Private.
DTH2	The current SPUTTER time step, set in the TDELT routine. Blank Common.
DTR	The time step that governs RAD. This is usually set to DTH2 in TDELT, but is set smaller in RAD if subcycling is called for, such that the SPUTTER master time step is an integral multiple of DTR. Blank Common.

DTRMIN	The actual minimum radiation time step calculated in RAD. This is then used to determine whether subcycling or supercycling is necessary. Blank Common.
DTR1	A "running minimum" time step, set so that at the end of the loop, the smallest and the next-to-smallest time step can be saved. Private.
DTR2	A "running second-smallest" time step, corresponding to DTR1 above. Private.
E	The SPUTTER zone array for internal energy per unit mass, used in the time step energy accuracy calculation. Blank Common.
EC	This is the heating rate due to conduction, zeroed out at the end of RAD to avoid later trouble. Blank Common.
EDITMF	The multifrequency edit flag, equivalenced to S12 in Blank Common. It should be noted that one sometimes wishes to get "multifrequency" prints when running grey problems. Some of the intermediate quantities printed out are of general interest and value in troubleshooting.
EK	This array is used to accumulate radiation energy density over frequency, serving the same purpose as the old variable SUMRHO (see Ref. 2), now deleted. RHO <sup>(3)</sup> is set to EK after exit from the frequency loop. Blank Common.
ELM	This is a special array of internal energy in a "lambda region" (see Ref. 3). It is used in RAD to form quickly the total internal energy in the radiation region for use in the time-step calculation. See WSB, WSSB. Blank Common.
EO	A special storage for the index of the zone which has the smallest radiation time step. Used in PRINT. Blank Common.
ER	The heating rate due to radiation transfer. While the radiation energy density and radiation pressure are used elsewhere in the code, and the flux is edited and viewed with interest, it is safe to say that calculating ER is the basic purpose of RAD and its subroutines. Blank Common.

FIOSV	Temporary storage for FIO while testing whether to iterate. Unfortunately, the third character is a letter "O," not a zero. Private.
FI1SV	Temporary storage for FI1 while testing whether to iterate. Private.
FI3SV	Temporary storage for FI3 while testing whether to iterate. Private.
G	This SPUTTER array of mass per volume factor is used in converting energy per unit mass to energy preparatory to calculating the energy accuracy time step. Blank Common.
GAMMA	This is the intermediate quantity $\bar{\gamma}$ defined beneath Eq. (15). Private.
GL	This is the right (exterior) boundary flag, specifying blackbody exterior if $GL = 1/2$ , vacuum otherwise. STRANS and PTRANS initialize a blackbody intensity for any positive value of GL (see GL <sup>(3)</sup> ). There are historical reasons for the inconsistency, but it is not justified. Blank Common.
GR	Maximum allowed number of radiation supercycles. Blank Common.
HCB	See p. 87 of this report. Blank Common.
HHTAX	This is $2 \kappa_s \bar{\gamma}$ , used to adjust the absorption coefficient for Compton scattering. See Eq. (33). Private.
HNUP4	The fourth power of the upper limit of photon energy for the current frequency group. Perhaps it is not used in the current version of RAD. Private.
HNUX	Set to the minimum of HNUP <sup>(3)</sup> and $10^5$ to avoid execution of Compton scattering at invalid frequencies. Used in forming A3 <sup>(1)</sup> . Private.
HNU4	The fourth power of the lower limit of photon energy for the current frequency group. This, like HNUP4, may be expendable. Private.
H3	This is $1/2 \rho \kappa_R \Delta r$ , where $\kappa_R$ is the local Rosseland mean absorption coefficient adjusted for scattering. Used only in RAD. The array H2 <sup>(3)</sup> is set to H3 if SOLID(10) is nonzero. Equivalenced to BR in Blank Common.

IA	The left-hand (interior) zone limit to the vapor region. IN <sup>(3)</sup> is set to it unconditionally in this version of RAD. IN can take on other values, if the programmer wishes to confine the radiation region to only part of the vapor (e.g., the part in local thermodynamic equilibrium), but this would require a small programming change. Blank Common.
IB	The right-hand (exterior) zone boundary limit to the vapor region. For generality, this should be replaced by IMP1 <sup>(4)</sup> . Blank Common.
IBM1	This is IB-1, the right-hand (exterior) zone limit to the vapor region. IM <sup>(3)</sup> set unconditionally to IBM1 in this version of RAD. Blank Common.
ICXM1	This is ICX <sup>(4)</sup> -1, evaluated as a separate variable solely for use as the upper limit of a DO statement. Private.
IDMX	The dimension limit for the X array, currently 4000. When testing an index against a dimension limit in the body of a computer program, it is good practice to have this a variable set at the beginning of the code to facilitate later changes in the size of the array. Private.
IJZILC	Purpose unknown. Private.
IMN1	This is the index of the zone with the smallest radiation time step. EO is set to IMN1. Private.
IMN2	This is the index of the zone with the second-smallest radiation time step. Private.
INP1	This is IN+1 defined solely to serve as the lower limit of a DO loop. Private.
IQEM	The index of the zone boundary suffering the worst change in FIO. Used in monitoring the scattering iteration. Private.
IR	The index of the rightmost (exterior) zone for which one wishes to calculate radiation. In reference 2, IR is set and used, but not discussed. In the fireball configuration, it was desired not to do radiation for zones colder than, say, 0.05 eV. IR was set to the index of the outermost zone warmer than that. The code is now in the rather hazardous situation that IR is set to IM if one does scattering or has a blackbody exterior. Otherwise, it is not evaluated at all. This should be repaired. Private.

J	A y-line index while the X <sup>(3)</sup> array is being formed. Also used as a zone-boundary index while enlarging the transport subregions by five mean free paths on each side. Private.
JDRUM	The logical unit, either 25 or 26, of the drum being read for moment quantities. PALMER Common.
JDRUMI	The logical unit, either 26 or 25, of the drum on which moment quantities are being written. Private.
JK	The y-line index advanced while y-lines are being formed. Private.
K	A counter, used as other than zero only if more than 4000 entries in the X array are required to place a y-line every zone. The code will successively try to place a y-line every (K+1) <sup>st</sup> zone until K reaches 10, at which point the code aborts. Private.
KK	A zone-boundary counter advanced as the X array is formed for a given y-line. Private.
KKK	A special index to insert the number of X entries for the current y-line into the X array. Private.
KMAX	The multifrequency flag. Zero, grey; nonzero, multifrequency. This is a standard SPUTTER input quantity. Blank Common.
K000FX	This divide check flag has zeros, contrasting with KOOOFX <sup>(3)</sup> . A genuine holdover from the days of SIFT. Private.
K1	A counter to skip zones so that a y-line is drawn every (K+1) <sup>st</sup> zone. See K above. Private.
LMDA(37)	Same as LMDA(37) <sup>(4)</sup> , but used in RAD for definition of external input intensities, a feature deleted from the OUTPUT code. (See page 81 of this report.) Reference to LMDA(37) should be removed from RAD. Blank Common.
MAXLM	The number of lambda regions. Used in calculating WSB from ELM. Blank Common.
NHNU	The number of frequency groups, specified in the OUTPUT code by LMDA(36). See p. 86 of this report. LINDLY Common.
NMU	The number of characteristic rays in plane geometry. Set to LMDA(37). The FORTRAN statement that does this should be removed. Private.

NRAD	The number of radiation subcycles, set at the end of the time-step calculation if appropriate. Abort if NRAD > 50. Blank Common.
NSMLR	Set to 1 at the beginning of RAD, it appears never to be used or reset. The statement should be removed. Blank Common.
NY	The number of y-lines if spherical geometry. LINDLY Comr. on.
Q	A dummy variable, it is $\Delta\tau$ (optical depth) in defining source gradients for forced diffusion and serves as the normalized "OUTPUT-output" flux. Private.
QD	The denominator in the relative difference expression for FIO, the iteration test. Private.
QE	The relative difference of FIO, tested against CAPAC(150). See p. 86 of this report. Private.
QEM	The largest of the QE's. Used along with IQEM to monitor the iteration. Private.
QN	The numerator in the relative difference expression for FIO, the iteration test. Private.
QQ1	This is $\rho\Delta r$ , the formulation of which is geometry-dependent. Used in forming optical depths. Private.
Q1	A zone array of $\theta^4$ , used to avoid calculating $\theta^4$ within the frequency loop. Very dispensable if one is tight for storage. Equivalenced to PB in Blank Common.
Q3	Integrated optical depth, a zone array used only in the section expanding transport subregions by five mean free paths on each side. See subsection 3.2 of reference 2 and p. 80 of this report. Equivalenced to GOFR in Blank Common.
Q31	A running integrated optical depth. Quite dispensable. Private.
Q37	Zone array for $\log \theta$ , used in linear-in-log interpolation for multifrequency absorption coefficients in DIANA. Equivalenced to CAR in Blank Common.
Q38	Zone array for $-\log \rho$ , used similarly. Equivalenced to CHIR in Blank Common.
RD	The SPUTTER array of velocity, used in computing the radiation work term portion of the heating rate. Blank Common.

RDD	Used to store radiation energy density calculated on the previous cycle. Blank Common.
SLUG	The fraction of the internal energy of a zone allowed to leak out (or be added; the latter is causing difficulties) of a zone in the time step calculated in the energy accuracy criterion. An input quantity, usually set to 0.1. Blank Common.
SMLR	Radiation pressure, summed over frequency, formed in RAD, used in HYDRO. Blank Common.
SOLID(10)	The Planck-Rosseland switch. Zero, Rosseland; nonzero, Planck. A standard SPUTTER input quantity. Blank Common.
SOLID(37)	See p. 86 of this report. Blank Common.
SUMX2	The sum over frequency groups of $X_2^{(3)}$ . $X_2$ is set to SUMX2 at the end of RAD. Equivalenced to CRTR in Blank Common.
SUMX3	Same as SUMX3 <sup>(3)</sup> . Zeroed in RAD for no good reason.
SUMX4	Same as SUMX4 <sup>(3)</sup> . Zeroed in RAD for no good reason.
SV	The specific volume, $1/\rho$ , a regular SPUTTER zone array. Used in forming optical depths. Blank Common.
TAUX	See the discussion of SOLID(37) on p. 86 of this report. TAUX is added to the absorption coefficient, so it is zero if SOLID(37) is negative; otherwise it assumes the value of SOLID(37). Private.
TAX	This is $2\kappa_s/m_e c^2$ . See HHTAX. Private.
TD	A variable set in MP2 as a print flag. Used in RAD to make the multifrequency print coincide with the regular SPUTTER print. Blank Common.
TELM(25)	An input quantity to adjust the time step by a constant factor. Blank Common.
TELM(26)	Set to DTR1. Blank Common.
TELM(27)	Set to IMN1. Blank Common.
TELM(28)	Set to DTR2. Blank Common.
TELM(29)	Set to IMN2. Blank Common.

TELM(30)	The advanced cycle number, making CNT1 <sup>(3)</sup> redundant. Blank Common.
TEMP(1, 2, 3)	Used as intermediate quantities in the time-step calculation. Blank Common.
TH	The SPUTTER variable indicating time. Used in evaluating BLANK3. Blank Common.
THETA	The zone array of temperature in eV. Blank Common.
THETAK(103)T4	See p. 86 of this report. Blank Common.
THTAMX	The highest temperature in the problem, formerly used in evaluating IR. Its sole purpose now is to bypass the radiation calculation for problems colder than 0.05 eV throughout. Private.
TSi	Used in calculating X values to indicate the point of closest approach if negative, or x if positive. Private.
T4	The scattering moment iteration counter. Compared with CAPAC(152). See p. 86 of this report. Private.
WSB	The total internal energy in the problem, used in the energy accuracy time-step criterion to ignore zones whose internal energy is less than 0.001 WSB. Private.
WSBB	The internal energy of a zone in the energy accuracy time-step calculation. Private.
X3	An array of flags to define the transport and diffusion subregions. If for a zone X3 is -1, the zone is in a diffusion region. If X3 = 0., the zone is in a transport region. Other values of X3 are nonsense. Blank Common.
X4	Used to denote $y^2$ for each y-line formed. Later zeroed and used by equivalence in PTRANS and STRANS. Blank Common.
ZP1(18)	Another place for the SPUTTER time step, used in calculating NRAD and DTR if there is to be radiation subcycling. Blank Common.
ZZ	A flag in the scattering moment iteration test. If any zone requires iteration, the flag is turned on. Off: ZZ = 0. On: ZZ = 1. Private.

APPENDIX I

THE OUTPUT CODE AS A MODIFICATION OF SPUTTER

INTRODUCTION AND SUMMARY

The OUTPUT code is a version of SPUTTER (Ref. 3) that is adapted to a special class of problems. Radiation transfer is especially important in determining the behavior of such problems, and the OUTPUT code differs from SPUTTER primarily in its more sophisticated treatment of radiation transfer and in the deletion of space-consuming portions of SPUTTER that deal with matter in the solid state. It should be borne in mind that the OUTPUT code is under development and is continually being changed. Any description of it will therefore rapidly lose accuracy of detail, and major changes made in the near future may go unreported for some time. However, the present configuration is a convenient one to describe thoroughly. Since a large fraction of the content of the OUTPUT code has been described in earlier reports (Refs. 2, 3), much of the material presented in this appendix pertains to changes made in the older programs. Three areas of the OUTPUT code are discussed: (1) changes in SPUTTER subroutines, other than radiation subroutines, (2) changes in the radiation routines, and (3) progress in treating Compton scattering. Changes now in progress include input quantities whose use differs from their use in standard SPUTTER and the abort indicators.

CHANGES IN SPUTTER SUBROUTINES OTHER THAN  
RADIATION SUBROUTINES

Altered Routines

MP2

This routine has been changed to eliminate references to CNDCTN and the boil codes, which have been deleted. Another change eliminates division by BLANK1 when it is zero. This change has also been incorporated in the standard SPUTTER code.

HYDRO

A small section that does "dummy hydro" (sets C and DELTAR to their appropriate values, but makes no changes in radii or velocities) has been added. This section is executed if S4 is negative. The option is very useful when one wishes to observe the effects of radiation transfer in a static configuration. Changes to delete references to solid material regions and permit the use of more accurate radiation pressures available from the OUTPUT radiation routines are now under development.

EOS

This routine has been changed to incorporate the actual radiation pressure, rather than an equilibrium diffusion approximation. Reference to ZPART, an array of 760 words not used by the CUTPUT code, has been deleted.

ECALC

Changes to eliminate treatment of solids and improve the precision of the source treatment are under development.

RTAPE

This routine, which picks up the desired configuration from the SPUTTER dump tape, has been modified to pick up additional data (moment quantities and, in the future, intensities) and set up the drum storage for them.

WTAPE

This routine does the write operations corresponding to RTAPE. Also, on initial starts, where there are no data to pick up, the drum storage is set up and appropriate starting values are written.

KAPPA

A change has been made to catch division by zero before exit. This may become standard.

KAP6

This routine, which is valuable in test problems, uses the THETAK array to define the frequency table and a corresponding set of absorption coefficients (independent of temperature and density) without using a DIANE tape. It is not available in the standard SPUTTER code.

KAP12

This routine gives an analytic approximation for the grey absorption coefficient of uranium as a function of temperature and density. It is not available in the standard SPUTTER code.

QUE8

This source routine deposits energy (evaluates SMLQ) into a predetermined set of contiguous zones, the energy being distributed uniformly over "mass space." The rate varies stepwise with time. The RDK array is used as input for the zone limits, the time cuts, and the energy deposition rates.

QUE9

This source routine is similar to QUE8, except that it allows for several (up to seven) regions or sets of contiguous zones, each of which has its own set of time cuts and energy deposition rates. RDK storage limits require that no region have more than six distinct time sections.

QUE10

This source routine is like QUE8 in that it allows only one region. But instead of the energy source being distributed uniformly, it is distributed as some power of the space variable, the exponent being determined by values of the energy deposition rate specified at the limits of the region.

DIVCHK

This is a machine-language routine which makes a report each time a division by zero is made. It is handy to have if one still wishes to abort the run when such a division occurs; in addition, there is a report precisely indicating where the division occurred. This routine, which was written by the Gulf General Atomic systems group, is not generally available or widely advertised, because it carries with it the hazard of uncontrolled computer behavior should the routine be overlaid by other coding subsequent to a call to it. The OUTPUT code precludes this possibility.

The MAP

All large programs on the UNIVAC-1108 make use of the MAP (Memory Allocation Processor) to fit their coding into the available storage with appropriate overlays. The SPUTTER and OUTPUT codes are no exception. The OUTPUT MAP differs from the SPUTTER MAP as follows:

1. DIVCHK is explicitly represented at the independent level with no overlay possible. (SPUTTER does not have DIVCHK at all.)
2. RADTN is segmented along with MP2, allowing an overlay with MP1, rather than being independent.
3. NAMEC (a name common block used only in MP1 and its subroutines) is segmented with MP1, so as to allow overlay with MP2, etc.
4. The overlay of the KAP routines has been deleted because, for reasons that are not now clear, it does not function properly. The storage penalty this causes has been minimized.

5. Several USE cards are provided for those routines bearing the same name as standard, undeleted SPUTTER routines.

#### Deleted Routines

The standard SPUTTER subroutines CNDCTN, BOIL, and CBOIL have been deleted, as have references to them in MP2. NONEQ and some other subroutines may soon also be deleted.

#### Dummy Routines

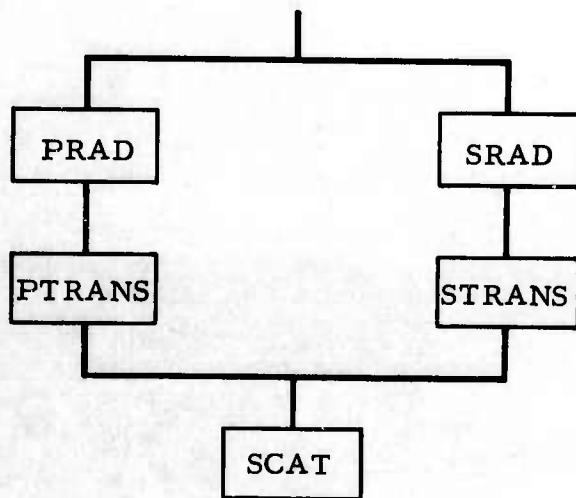
It has been convenient in several instances to write dummy subroutines with the same name as standard SPUTTER subroutines, since outright deletion of the latter would require extensive changes elsewhere in the code. The standard subroutines and the reasons why they are undesirable are as follows:

1. DRAD, ERAD, QUE4. These refer to a 642-word array not needed by the OUTPUT code.
2. QUE16, NONEQ. These refer to a 760-word array not used by the OUTPUT code. NONEQ, furthermore, has nearly 300 words of private data storage.
3. CMOL, KAP5. Both have considerable private data storage.

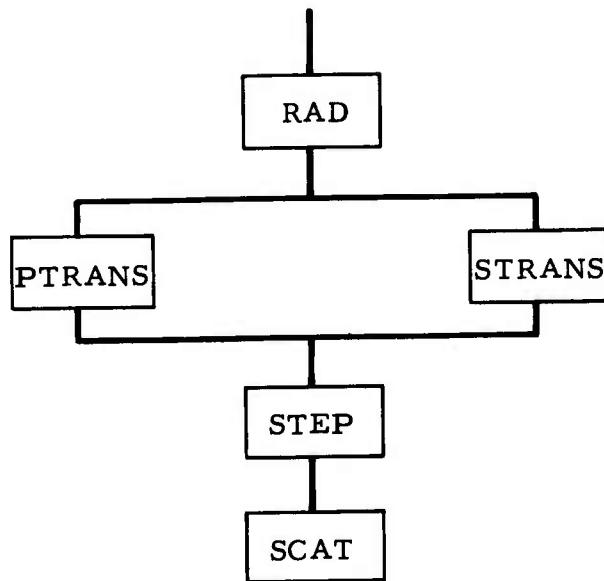
#### CHANGES IN RADIATION ROUTINES

Two substantial changes in program organization have been made in the past few months. (1) The codes PRAD and SRAD, executed for plane and spherical geometry, respectively, have been replaced by one code, RAD. The two codes were over 90 percent the same, line for line, and considerable duplication had taken place in maintaining them. The necessary allowances for geometry were incorporated, and the codes were unified. This change has been made both in the standard SPUTTER and the OUTPUT codes. (2) In PTRANS and STRANS, the coding to solve the transport equation was repeated many times and in several forms to allow for different

configurations. (See the STRANS program listing in reference 2.) Considerable duplication of effort has resulted because scattering led to considerable changes in these codes, other minor changes are made from time to time, and more sophisticated solutions of the transport equation are being coded. The advantages of solving the transport equation in only one place in the code became more apparent. The new subroutine, STEP, solves the transport equation in general form for one increment along a characteristic line. Calls to STEP have been incorporated in STRANS and PTRANS. The data block COMMON/JIM/, originally provided for linkage between the transport routines and SCAT, has been expanded by six words to accommodate STEP. The call hierarchy of the radiation routines was formerly as follows:



The new hierarchy, which results in considerable reduction of program storage and much more understandable codes, is as follows:



STEP has been added only in the OUTPUT code.

Another change made in both the standard SPUTTER and the OUTPUT versions of RAD is the elimination of several input parameters. The deleted parameters are described briefly in table II and in detail in reference 2. Some of the parameters listed in reference 2 were assigned different SPUTTER variable names after the reference was issued. The more recent names are also given in table II. In some cases, the option involving the parameter has been removed altogether, and for these no "built-in" value is given. Instances where built-in values differ from the suggestions in reference 2 reflect experience in running problems.

Two rather substantial deletions have been made in the OUTPUT version of RAD. (1) ICX and ICY have been eliminated as special indices indicating sourceless subregions of the radiation region. This eliminates quite a bit of complicated branching, and the time saved by doing simpler calculations in sourceless regions, although substantial in some configurations run on standard SPUTTER, would be insignificant in most applications of the OUTPUT code. (2) The multifrequency merge procedure

TABLE II  
DELETED VARIABLES

Name Used in reference 2	Recent name(s)	Value suggested in reference 2	Built-in value	Purpose
CB*	DELPRT	0	---	Brightness print (option deleted)
GA	GA	0.33	0.333	Source, optical depth gradient criterion (TG)
GL	GA, GR, (0.333)	0.3	---	Source, MFP gradient criterion (y-lines, option deleted)
CMIN	AC	0.3	--( $\infty$ )	Minimum depth for TG criterion (option deleted)
ACO3T4	TA	0.1	0.01	Half-optical depth for thick-thin transition
S15	S15	1	---	Restart on grey calculation, not needed with current DIANA
TELM(37)	TELM(37)	0.005	0.001	Fraction of total energy in zone for time-step criteria
CVB	CVB	0	0.5	Select y-lines
HVB ..	HVB	5	5	Buffer of transport region in number of mean free paths
HCB	HCB	0.1	0.1	Diffusion criterion
Not mentioned	BOILB = NTIMES	--	50	Subcycle limit

\* CB is currently used in standard SPUTTER as a multifrequency merge criterion, replacing the constant 10.0 in reference 2. The OUTPUT code has no frequency merging.

has been taken out. Although it may save time in some applications, no provision for merging has been made in the scattering coding. The variable CB, of course, is now not used in the OUTPUT code. Another, less significant feature deleted in the OUTPUT code is the provision for external nonthermal radiation intensities as an energy source in plane geometry.

A major addition to RAD is the calculation of the moment quantities, used as scattering source terms in the SCAT routine, first reported in reference 7. An iteration procedure to obtain better values of the moment quantities at cycle  $(n + 1)$  (the scattering equation is implicit in this respect), described in section 1 has been incorporated. The FQ0, FQ1, FQ2, and FQ3 arrays should not be updated inside the iteration loop, since they are the moment quantities for the previous frequency. This error would cause trouble only in Compton scattering. The updating has been moved outside the iteration loop, and the scheme now works well for both Thomson and Compton scattering. The choice of three words at the end of the CAPAC array as input parameters to control the iteration has been retained. These are described in "Input Quantities" of this appendix. However, the size of the array is the standard SPUTTER value of 152.

The X array, the set of all values  $x_{ij} = \sqrt{r_i^2 - y_j^2}$ , is calculated early in RAD to avoid taking many square roots within the frequency loop. In reference 2, this array was given 2400 words of essentially private data storage. (This would allow 66 zones if a y-line were drawn at each boundary.) In the standard SPUTTER version of RAD, the X array is stored in 1064 words of SPUTTER Blank Common, allowing for 41 zones with y-line every zone. Since OUTPUT calculations will have many zones as a result of complex configurations, and since accuracy in the scattering calculation requires a y-line at every boundary, the private storage has been reinstated in the OUTPUT code; this storage is 4000 words, allowing for 86 zones. The

scratch area in SPUTTER Blank Common made available by this change has been taken up by intermediate moment quantities and scattering parameters.

Three variables used in the Compton scattering calculation are dependent only on frequency and not on position:

$$1. \quad GMP = 1 - 2\gamma, \text{ where } \gamma = \frac{1/2(h\nu_- + h\nu_+)}{m_e c^2}$$

$$2. \quad A1 = \frac{h\nu_-^2}{m_e c^2 (h\nu_+ - h\nu_-)} + 3\gamma$$

$$3. \quad A3 = \frac{h\nu_+^2}{m_e c^2 (h\nu_{++} - h\nu_+)}$$

Although evaluated in SCAT, they should be calculated only once per frequency in order to save time. This is done in RAD.

#### PROGRESS IN TREATING COMPTON SCATTERING

Satisfactory results have been obtained for a test problem with spherical geometry analogous to the test problem with plane geometry described in appendix IV. However, some interesting difficulties arose, which have not yet been resolved. At high frequency, for sufficiently large negative  $dI/d\nu$ , where  $I$  may be intensity or some moment quantity (these usually vary similarly with  $\nu$ ), the code will calculate a negative scattering intensity. At present, the code sets the negative intensity to zero and goes on. (There are input parameters to trigger a debug print or abort, or both.) The large negative  $dI/d\nu$ , with the concomitant negative intensities, has appeared in two different situations.

First, in the calculation of the highest-frequency group, the starting value for  $I_x(\nu_j)$  is zero. This results in the appearance of negative intensities, but at least some sort of radiation field description is achieved. Two schemes

have been proposed for improving the formulation of  $I_x(v_j)$ . The better one will soon be chosen and incorporated in the code. However, this case of the first frequency is a rather artificial way of arriving at large negative  $dI/dv$ .

The other situation arose when an attempt was made to refine the frequency interval. The result was, for one frequency group, a bistable oscillation in the iteration procedure for obtaining updated moment quantities. It is believed that it arose in the following manner. First, given zero starting values for the moment quantities, the code had a positive  $dI/dv$  and calculated scattering intensities, generating moment quantities appropriate for that frequency. Then the code noted that an iteration was necessary and set the "old" moment quantities to the "new" ones. (See "Input Quantities" in this appendix. CAPAC(151) was 0.) It proceeded to do Compton scattering, found large negative  $dI/dv$ , calculated negative intensities, set them to zero, generated zero moment quantities as a result, iterated again, and at the third pass was right where it started. While this may not be an exact description of what happened, it is probably close enough to indicate the trouble unambiguously. The "extrapolation switch" described in section II of this volume (CAPAC(151) in the code) could easily be used as an interpolation switch by assigning to it a value between 0 and -1. Interpolation is suggested by the fact that the first "crack" at a solution, resulting in large negative  $dI/dv$ , may have been an overshoot. This idea was tried out with CAPAC(151) = -0.5 and proved successful. There was no convergence difficulty in areas of negative  $dI/dv$ . However, at lower frequencies the convergence, while monotonic, was discouragingly slow. This suggests a possible frequency-dependent extrapolation switch.

In summary, Compton scattering in the OUTPUT code is at present a laboratory curiosity. It appears to work well, if not reliably, in an interesting test problem. (It takes about three times as long to calculate as Thomson scattering.) However, this tentative state should change shortly as experience is acquired in using Compton scattering in production

problems. It should be emphasized that Thomson scattering has been used for over six months to calculate complex problems, and a number of shortcomings, which appear only in unforeseeable configurations, have been found and overcome.

#### CHANGES IN PROGRESS

Several areas of current code development relevant to the OUTPUT code are not discussed in detail in this report. These include improvements in the nonequilibrium radiation diffusion code DRAD, incorporation of Compton scattering in DRAD, changes in the energy bookkeeping (routines ECHK and PRINT), and modifications to provide free drum input-output. Three other changes in the program are as follows:

1. Improved source interpolation in spherical geometry. The current code follows the original STRANS (Ref. 2) logic of step-or-linear interpolation of the source, with a linear-quadratic fit across the "top slice." This is being replaced with a linear-in- $r^2$  interpolation scheme developed by J. R. Triplett (Ref. 3). The option of a piecewise constant source, as described in reference 2, is retained. Each "step" along a y-line is divided into two intervals equal in  $r^2$  to conform with the new interpolation scheme. Although this change has been essentially completely checked out, it is not included in this report.
2. The Sampson approximation of the average absorption coefficient, where  $\bar{\kappa} = (b + \kappa_R)/(b + \kappa_p) \cdot \kappa_p$  and b varies as the integrated optical depth, appears to be a desirable first step toward a more sophisticated transmission function treatment than the one now being used.
3. Saving intensities. Future code development in several areas, e.g., retardation, more accurate calculation of moment quantities, and improved Compton scattering recipes, will require the storage

of the entire set of intensities calculated at cell-ray intersections.

This change also was completed too late for inclusion in the present report.

For each frequency, given that the "x" storage ( $x_{ij} = \sqrt{r_i^2 - y_j^2}$ ) for all values on one side of the plane of symmetry is 4000 words, the intensity store will require 8000 words of random access storage, plus 8000 times an upper limit to the number of frequencies for drum storage. This requirement can be met, but not economically on the UNIVAC-1108. (The UNIVAC-1108 has the odd feature of a  $3/8\text{-}\mu\text{sec}$  fetch from opposite bank versus a  $3/4\text{-}\mu\text{sec}$  fetch from same bank, so that it pays to put coding and data in opposite halves of the "almost-random access" fast storage.) The speed penalty can probably be minimized so as to be negligible on long, expensive problems by an adroit choice of what data are to be kept in the wrong half of the store. Other problems are definition of the appropriate set of intensities for a cell boundary treated by diffusion and definition of intensities for y-lines where a specific calculation is not made. At present, the OUTPUT code sets up a y-line at every zone boundary and calculates intensities along all of these. However, one might wish to run a problem with more than 86 zones, the maximum possible for less than 8000 intensities, or one might wish to skip y-lines to save time.

#### FORTRAN LISTINGS

Appendix III contains FORTRAN listings for the principal subroutines of the OUTPUT code: RAD, PTRANS, STRANS, STEP, SCAT, QUE8, QUE9, QUE10, HYDRO, RTAPE, WTAPE, and the MAP. For brevity, SPUTTER blank common, which is used in all the subroutines, is listed only in RAD. The equivalence table used in RAD, PTRANS, STRANS, and STEP is also listed only in RAD.

INPUT QUANTITIES

The input quantities listed below are (1) parameters used in the Compton and Thomson scattering treatment, (2) special parameters for the "OUTPUT-output," a special edit, and (3) various quantities used by the nonstandard source and opacity routines.

THETAK(103)	Required input. Sets value of HNUP. Compton scattering does not allow large values of HNUP. There is no safety coding, such as "if zero, set to $10^6$ ." It must be entered.
SOLID(36)	Compton switch. If zero, Compton scattering is calculated; if nonzero, Thomson scattering.
SOLID(37)	Scattering coefficient. If negative, it is not added to the absorption coefficient from the DIANE tape, but the absolute value is used as the scattering coefficient. If positive, $\kappa$ (DIANE) is assumed to be $\kappa_a$ only, and $\kappa_s = \text{SOLID}(37)$ is added.  $\text{SOLID}(37) = 0$ would provide a very inefficient "no scattering" calculation.
LMDA(36)	Required input for NHNU. The subroutines WTAPE and RTAPE, which set up the drum storage for the scattering source terms and process these on a non-standard SPUTTER dump tape, must know the value of NHNU before the DIANE tape is read. Using the dimension limit of 20 would avoid this, but would lead to inefficiencies. (MUST BE LOADED IN FIRST SET OF CARDS.)
CAPAC(150)	Scattering iteration convergence coefficient, usually 0.05.
CAPAC(151)	Scattering iteration extrapolation parameter, usually 0.5.
CAPAC(152)	Scattering iteration recycle limit. Zero, no iteration. Maximum allowed value is 8. Usual value is 4.
EPSI	Radius of OUTPUT sample tube.
LMDA(26)	Index of zone boundary at which OUTPUT sample tube is affixed.

- CVB      If CVB is zero, negative scattering intensities are set to zero and the problem continues. If CVB is nonzero and a negative scattering intensity arises, the code calls UNCLE.
- HVB      If HVB is zero, negative scattering intensities go unreported. If HVB is nonzero and a negative scattering intensity arises, a brief data report is given, and a y-line print is triggered.
- HCB      The normal SPUTTER RAD meaning applies here. Negative forces diffusion; positive forces transport; zero lets the code decide.
- RDK      This array is used in standard SPUTTER to provide input for the various source routines, and the OUTPUT code does the same.
- QUE8: Energy is supplied uniformly in mass space within a defined region (i.e., all zones within the region receive a constant ergs/sec/gm) at a rate that varies stepwise with time. RDK(1) is the index of the first zone in the region; RDK(2) is the index of the last zone in the region. Up to 50 time cuts can be specified in RDK(3) - RDK(52). The 49 rates for the corresponding periods (in ergs/sec for the whole region) are given in RDK(54) - RDK(102). Energy source rate for  $t > RDK(52)$  is zero.
- QUE9: This routine is similar to QUE8, except that the fine time definition is sacrificed for some spatial definition. Up to six different contiguous regions can be defined (seven bounding zone indices provided). Each has a separate set of up to six source rates and time cuts. Detailed input specification appears in the FORTRAN listing.
- QUE10: Energy is supplied as  $g^{-x}$  in mass space, where g is a mass space coordinate and x is an exponent determined by the average energy rate and by the rate specified at the right boundary. The two rates are specified for each given period. Ten periods are allowed. Time cuts are given by RDK(3) - RDK(13), the average rates by RDK(54) - RDK(63), and the boundary rates by RDK(44) - RDK(53).
- THETAK    This array is used as input to KAP6/JP. If one wishes to run a multifrequency problem without using a DIANE tape, one must specify the number of frequency groups and the boundaries, in eV, of the desired

frequency groups. KAP6 specifies these and also a constant absorption coefficient, independent of density and temperature, for each frequency group. Up to 20 frequency groups are allowed. THETAK(61) - THETAK(80) gives the lower boundary frequencies. THETAK(81) - THETAK(100) gives the absorption coefficients. THETAK(101) gives the number of frequency groups; THETAK(102) gives the grey absorption coefficient; and THETAK(103) gives the upper boundary frequency of the top group, as indicated on p. 86 of this report.

### S1 FLAGS IN THE RADIATION ROUTINES

The SPUTTER code follows the practice of setting the variable S1 to some value and calling UNCLE in case of serious trouble. The integer portion of the S1 flag indicates the subroutine, and the four digits after the decimal point indicate the FORTRAN statement number. The S1 flags for the radiation routines in the OUTPUT code are listed and commented on below.

#### RAD

<u>S1 Flag</u>	<u>Immediate Cause</u>	<u>Probable Remote Causes and Comments</u>
13.0102	Zero or negative CAPAC or CAPAR (grey)	An opacity of interest has not been evaluated. This usually happens with untested KAP routines.
13.0112	Small $\Delta t$ causes more than 50 radiation subcycles	Can be caused on first cycle by improperly set TELM(25) or SLUG, or any time by actual pathologies in CV, THETA, SV, or opacity, or simply by failure to start the problem with a small enough time step.
13.0150	Divide check	Division by zero. See the preceding DIVCHK diagnostic print for locations of divide instruction and divisor.
13.0119	K > !0	At least 10 attempts have been made to space y-lines so that x-storage is not overtaxed, and these have failed. A code change is necessary, or the problem must be rerun with fewer zones.

<u>S1 Flag</u>	<u>Immediate Cause</u>	<u>Probable Remote Causes and Comments</u>
13.0452	Divide check	Same as 13.0150.
13.0460	Negative INM1	Absurd. Indicates a radiation zone with zero or negative index.
13.0530	Zero or negative CAPAC or CAPAR (multifrequency)	See 13.0102
13.0652	Zero or negative optical depth	Something has gone wrong with the definition of H3 in the preceding lines of code.
13.0654	GL > 0.9, HCB < 0	An attempt is being made to force diffusion and provide external source intensities, two incompatible problem specifications.
13.0692	Divide check	Same as 13.0150.
13.0720	Positive X3	Absurd. Code sets X3 either zero or negative as a diffusion flag.
13.0962	IALPHA = 2	A transport calculation in cylindrical geometry is being attempted.
13.0982	IR < IM	Absurd in present code. IR set to IM earlier.
13.1070	Divide check	Same as 13.0150.

PTRANS

14.0100	Divide check	Same as 13.0150
14.0120	IALPHA > 1	Somehow PTRANS called for nonplane geometry. Absurd.
14.0160	INM1 negative	See 13.0460.
14.0350	GL > 0, GL ≠ 0.5	The external source intensity option has been deleted. GL therefore has a meaningless value.
14.0360	IBX > IM	Absurd. A portion of the transport region seems to be outside the radiation region.

STRANS

<u>S1 Flag</u>	<u>Immediate Cause</u>	<u>Probable Remote Causes and Comments</u>
14.0001	Divide check	Same as 13.0150.
14.0005	IALPHA ≠ 3	Somehow STRANS called for nonspherical geometry. Absurd.
14.0011	IBX > IM	See PTRANS 14.0360.
14.0087	y-line index out of range	The most likely cause for this stop is scrambled radii.
14.0152	Divide check	Same as 13.0150.
14.0171	IBX > IM	See PTRANS 14.0360.
14.0207	X(K) = 0.0	This is an absurd stop. Any occurrence of it in the OUTPUT code would have to be carefully examined for the cause.

STEP

No S1 interrupts appear in the program.

SCAT

75.001	Divide check	Same as 13.0150.
75.008	FS < 0, CVB ≠ 0	Negative scattering intensity with abort flag set. If HVB also ≠ 0, a diagnostic print is given.
75.0009	Divide check in SCAT	Same as 13.0150.

KAPPA

15.0500	Bad QLM (J + 17)	Improper optical property specification resulting in an attempt to call a non-existent KAP routine.
15.0810	Divide check	Same as 13.0150.

KAP6

<u>S1 Flag</u>	<u>Immediate Cause</u>	<u>Probable Remote Causes and Comments</u>
	No S1 interrupts appear in this program.	

KAP12

No S1 interrupts appear in this program.
--

## APPENDIX II

### DESCRIPTION OF INPUT FOR A SAMPLE PROBLEM

The debug test problem described in this appendix was run on the UNIVAC-1108 about September 15, 1967, and on the CDC 6600 at the Air Force Weapons Laboratory, Kirtland AFB, New Mexico, on September 22. The two runs matched within expected roundoff differences. The calculation was a multifrequency transport calculation with Compton scattering, using dummy hydrodynamics and the KAP6 opacity routine.

The problem consisted of a sphere with a radius of 10 cm, which contained 10 zones, with a temperature of 8 keV, and was surrounded by a shell with a thickness of 10 cm, which contained 20 zones, at a temperature of 600 eV. The density was taken as one in all zones, and the Rosseland opacity as  $0.2 \text{ cm}^2/\text{g}$  in all groups.

The input deck consists of SPUTTER Common input cards only, except for an initial header card which contains arbitrary identifying information in Columns 1-72. These Common input cards have the following format:

Col. 1: 1 denotes the last card of the deck and must appear on that card. 2 specifies that the data on the card will be converted to fixed point and appear as INTEGER variables in Common.

Blank (or 1) specifies that the data on the card will appear as REAL variables (floating-decimal form) in Common.

Cols. 2-6: Location relative to start of Blank Common in which the first data word on the card is to be stored.

Col. 7: Number of data fields on the card (maximum 7).

Cols. 8-70: Up to seven data fields, of nine columns each, contain numerical data to be read under the format OP7E9. 4. Blank fields generate zeros (unless they are excluded by the field number in Col. 7).

Cols. 71-80: Card identification information.

The test problem input deck is shown in table III. Each card, identified by its initial location number, is discussed below. Data fields containing 0. can, of course, be omitted, but are included in several cases for expository reasons.

- 68: FREQ gives the number of cycles between prints. CNTMAX, in 69, gives the cycle number at which the problem is to be terminated.
- 1: Zone-boundary indices for the "lambda" regions (which are usually regions of different materials). For one lambda region of 30 zones, the two limits, 1 and 31, are given.
- 19: LMADA(19) is a flag which, if >0, will trigger a complete print of Blank Common as if there were an error exit, even if the problem terminates normally. Since the extra information is often useful, setting LMADA(19) to 1 has become standard practice.
- 36: LMADA(36) is a special quantity specifying the number of frequencies. This is discussed further on p. 86 of this report. In the present case, there were 20 frequencies.
- 40: Several numbers are given on this card: the SPUTTER variables IA, IB, ICA, ICB, and KMAX, respectively. The nonzero value of KMAX dictates multifrequency. The other numbers, index limits, are required by the special XCARDS as a substitute for normal SPUTTER problem generation.
- 53: IG is another quantity needed by XCARDS. This is the extreme upper-limit index to the problem. When one remembers that SPUTTER has indices for upper limit to vapor, upper limit to solid, upper limit to radiation regions, upper limit to non-LTE region, etc., the significance of the term "extreme upper limit" becomes clear; and IG does serve a purpose.
- 65: TMAX is the problem time at which the problem is terminated. An absurdly large value is set here, since one wishes to limit the run by cycle count instead.

TABLE III  
TEST PROBLEM INPUT DECK

COMPTON TEST --	20 FREQ	WITH CAPAC(151) = -0.2	SEPTEMBFR 18, 1967	Column
1 2-6	7	8-16	17-25	26-34
2 68	3 5.0	20.	35-43	44-52
2 12	1.	31.	53-61	62-70
2 19	1 1.0			FREQCNTMAX
2 36	1 20.			LMDA(19)
2 40	5 1.	31.		LMDA(36)
2 53	1 31.			IG
2 65	1 1.0	+05		TMAX
	77	1 0.		CVB
	78	1 .1		SLUG
	79	1 3.0		ALPHA
	81	1 0.		HVB
	83	1 0.		HCB
	90	1 0.0		GL
	93	1 .	+10	RHOR
100	1			JC
107	1	1.0	-11	TE
109	1	9.0	-09	DTMAX1-2-3
115	3	1.0	-13	TRDBG
127	1	0.		S2
137	1	3.		S3 NO COND
138	1	1.0		S4 NO BOIL
139	1	-1.0		S5
140	1	10.0		EDITMF
147	1	1.0		
151	1	0.0	1.	
160	7	0.	2.	6.
			3.	5.
			4.	

(Table continued on next page)

TABLE III (continued)

COMPTON TEST -	-20 FREQ	WITH CAPA	C(151) =	-0.2	SEPTEMBER 18, 1967	
Column						
1	2-6	7	8-16	17-25	26-34	35-43
	1677	7.	3.	9.	10.	10. 5
	1747	12.	12. 5	13.	13. 5	14.
	1817	15. 5	16.	16. 5	17.	14. 5
	1883	19.	19. 5	20.	17.	18.
2021	7	1.	1.	1.	1.	11. 5
2028	7	1.	1.	1.	1.	15.
2035	7	1.	1.	1.	1.	18. 5
2042	7	1.	1.	1.	1.	
2049	3	1.	1.	1.	1.	
2325	7	8000.	8000.	8000.	8000.	
2332	7	8000.	8000.	8000.	8000.	
2339	7	600.	600.	600.	600.	
2346	7	600.	600.	600.	600.	
2353	3	600.	600.	600.	600.	
7338	3	.01	-0.2	8.		
8405	1	1.				
8681	1	206.				
9405	1	1.				
9332	7	76000.	72000.	64000.	60000.	\$2000.
9339	7	48000.	44000.	36000.	32000.	24000.
9346	6	20000.	16000.	8000.	4000.	
9352	7	.2	.2	.2	.2	
9359	7	.2	.2	.2	.2	
9366	6	.2	.2	.2	.2	
9372	3	20.	.2	80000.		
8466	1	0.5				
8853	1	10.0				
8884	1	0.				
8885	1	- 1999				
8882	1	0.				
	1					
						TEL M(25) SOLID(10) SOLID(36) SOLID(37) SOLID(34)

- 77: CVB, described on p. 87 of this volume, is set zero in this case because one knows that negative scattering intensities will arise in the first frequency group, and one wishes to ignore them.
- 78: This quantity, SLUG, is described on p. 71 of this report. Its normal value, 0.1, appears here.
- 79: ALPHA (see p. 17 of this report). Set 3. for spherical geometry.
- 81: HVB (see p. 87 of this report). Set zero to bypass debug print.
- 83: HCB (see p. 87 of this report). Set zero to let code decide.
- 90: GL (see pp. 58 and 67 of this report). Set zero to provide exterior vacuum.
- 93: RHOR. Special flag for air equation of state. Set big by habit. Not needed for this problem.
- 100: RPIA. If zero, radiation terms are added in the equation-of-state calculation. Normally set zero.
- 107: TC is the time before which DTMAX1 applies, and after which DTMAX2 applies.
- 109: TE is the time before which DTMAX2 applies, and after which DTMAX3 applies.
- 115: The time step is usually controlled by physical considerations (e. g., hydro or radiation time-step controls). However, if the user feels that these might be too generous in some cases, he may specify an upper limit to the time step, one for each of three specified time intervals, namely, DTMAX1, DTMAX2, DTMAX3. Some number must be specified for these. In this problem, it was desired that a quasi-steady-state radiation field be established in a time short compared with cooling times. Since retardation is not included here, an arbitrarily small time step is indicated, and one was chosen.
- 127: TRDBG, otherwise known as AC03T4, described on p. 54 of this report. Set zero to avoid debug print.
- 137: S2 is the radiation flag, tested in MP2 and in the switching routine RADTN. A value of 2 means that RAD is to be called. Zero bypasses radiation altogether.
- 138: S3. Set nonzero to bypass conduction. Since the appropriate subroutine has been deleted from the OUTPUT code, this input quantity is unnecessary.

- 139: S4. Its original meaning, "don't call BOIL," has, with the deletion of BOIL, been changed to allow for "dummy hydro" or no motion if set negative. This was desired for this test problem.
- 140: S5. Used to indicate the logical unit of the SPUTTER dump tape. The value 10 is a typical one.
- 147: EDITMF, otherwise known as S12. Described on p. 66 of this report. Set to 1 to deliver the multifrequency edit.
- 151: S16. This is a flag that tells MP1 (the generator-setup-rezone section of SPUTTER) whether the problem is an "initial" or cold start or a "restart." Zero, the former; one, the latter. This problem is an "initial" start, where all parameters must be supplied by the user and a zero-cycle dump prepared. Hence the value zero.
- 160-188: R (radii) for the 31 zone boundaries.
- 2021-2049: SV (inverse densities) for the 30 zones.
- 2325-2353: THETA (temperatures in eV) for the 30 zones.
- 7338: These are the three parameters controlling iteration on the moment quantities. They are described on p. 86 of this volume and are also discussed on p. 83. It is stated on p. 83 that CAPAC(151) = 0. did not work, whereas CAPAC(151) = -0.5 worked well. CAPAC(151) = -0.2, which was used in this problem, proved even better.
- 8405: OKLM(1). This equation-of-state flag results in a call to EIONM5 for hydrogen. A suitable dummy material.
- 8681: The value 206. results in a call to KAP6, a special OUTPUT opacity subroutine described in Appendix A.
- 9405: MAXLM is the number of "lambda" regions in this configuration.
- 9332-9346: THETAK(61-80) is used by KAP6 to define the frequency groups.
- 9352-9366: THETAK(81-100) is used by KAP6 to define the opacities.
- 9372: THETAK(101-103) is used for KAP6 to define NHNU, the grey opacity, and the top frequency limit.
- 8466: TELM(25) (defined on p. 71 of this report). Set to 0.5 as a holdover from another problem. Not needed as long as the time step is held down by DTMAX1, 2, 3.
- 8858: SOLID(10) (defined on p. 71 of this report). Set nonzero for Rosseland opacities. Superfluous when KAP6 is used, as in this case.

- 8884: SOLID(36) (defined on p. 86 of this report). Set zero for Compton scattering.
- 8885: SOLID(37) (defined on p. 86 of this report). The peculiar value of -0.1999 is given so that  $\bar{\kappa} > \kappa_s$  in all cases.
- 8882: SOLID(34), the starting cycle number. The 1 in Col. 1 indicates that input is complete. Contrary to normal SPUTTER generation, this style of problem specification has only one set of input cards.

A second test calculation has been carried out to test the sensitivity of the results to the parameter S4. The problem was as follows:

Region 1       $0 < r < 15$  cm

$$\rho = 1.$$

$$\epsilon = 8-.433 r \text{ keV}$$

$$\Delta r = 1. \text{ cm}$$

Region 2       $15 < r < 150$  cm

$$\rho = 31/r^3$$

$$\theta = 600 \text{ eV}$$

$$g = 10 \text{ (i. e., total mass of each zone is } 4/3 \pi \times 10 \text{ g)}$$

The Rosseland optical depth of Region 2 was 0.7. Three values of S4 were tried: 0.33, 0.5, and 1. The first of these effectively forces use of motion equation differencing based on Eq. (64) even in the diffusion limit. The singularity at  $r = 0$ , represented by the last term in Eq. (64), causes the velocity to diverge near the center, as shown in figure 7 for time  $2.36 \times 10^{-9}$  sec. The use of an S4 value higher than 1/3 allows the use of Eq. (63) in the central region, and more reasonable results are obtained. No significant difference between the results for  $S4 = 0.5$  and 1.0 at low radii or between any of the results for  $r > 4$  cm was noted. A calculation with a very small core radius would presumably indicate that the value 0.5 is superior, but additional work is required before this point can be established.

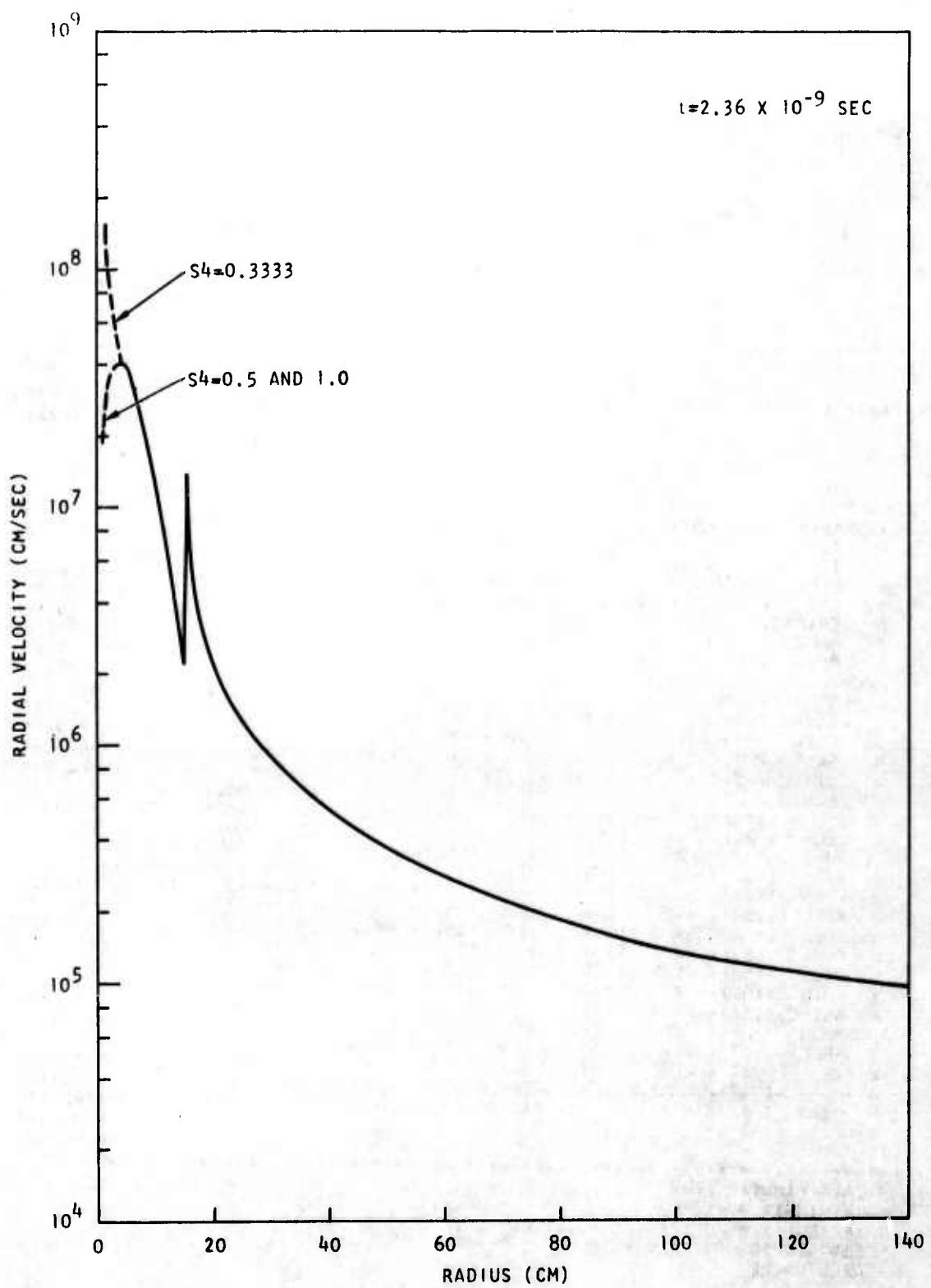


Figure 7. Dependence of Radial Velocity on Parameter S4

## APPENDIX III

LISTING OF OUTPUT CODE ROUTINES

```

WIT FOR QUES/OUT, QUER/OUT, QUE8/OUT1
SUBROUTINE QUES
C SPECIAL SOURCE FOR J. PALMER
C RDK(1)= LEFT ZONE BOUNDARY
C RDK(2)= RIGHT ZONE BOUNDARY
C RUK(3-52)= TIMES FOR GIVEN EDOT
C RUK(53-104)= EQUITS FOR GIVEN TIMES
C*****KAP20070
C*****KAP20080
C*****KAP20090
C*****KAP20100
C*****KAP20110
C S P U T I E R C O M M O N
COMMON LMDA(37), NR , NSMLR , IA , IB , ICA , ICB , KAP20120
1 KMAX , BLANK1, BLANK2, BLANK3, IAPI , IBP1 , ICAP1 , ICBP1 , KAP20130
2 II , IG , NRAD , BLANK4, IAM1 , IBM1 , ICAM1 , ICBM1 , KAP20140
3 IIP1 , IGM1 , IALPHA, BLANK5, TH , TMAX , BLANK6, DELPRT, KAP20150
4 FREQ , CNTMAX, AR , ASMLR , PUSHA , PUSHB , BOILA , BOILB , KAP20160
5 CVA , CVB , SLUG , ALPHA , HVA , HVB , HCA , HCB , KAP20170
6 EMINA , EMINB , CA , CB , GA , GB , GL , GR , KAP20180
7 RHOL , RHOR , EPIO , EPSI , RIA , RIB , RDIA , RDIB , KAP20190
8 RPIA , RPIB , RPDIA , RPDIB , TPRINT, TA , TB , TC , KAP20200
COMMON TD , TE , DTH2 , DTH2P , DTH1 , DTRMIN, DTMAX , KAP20210
1 DTMX1, DTMX2, DTMX3, DTR , SWITCH, CO , CMIN , DELTA , KAP20220
2 GAMA , WCRT , SIGMAQ, AC , AC03T4, CNVRT , SUMRA , SUMRB , KAP20230
3 ROIA , ROIAM1, ROI8 , ROI8P1, GMS , S1 , S2 , S3 , KAP20240
4 S4 , S5 , S6 , S7 , S8 , S9 , S10 , S11 , KAP20250
5 S12 , S13 , S14 , S15 , S16 , S17 , S18 , S19 , KAP20260
6 S2U , EO , FO , TAU , ZERO , R (152) , DELTAR(152) , KAP20270
7 ASQ (152) , RU (152) , VD (152) , RDD (152) , SMLR (152) , KAP20280
8 DELR ( 37) , P (152) , PI (152) , PB (152) , PB1 (152) , KAP20290
COMMON P2 (152) , SV (152) , RHO (152) , THETA (152) , KAP20300
1 W (152) , E (152) , EI (152) , EK (152) , A (152) , KAP20310
2 V (152) , G (152) , D (152) , C (152) , X2 (152) , KAP20320
3 X3 (152) , X4 (152) , X5 (152) , X6 (152) , X7 (152) , KAP20330
4 SMLA (152) , SMLB (152) , SMLC (152) , SMLD (152) , SMLE (152) , KAP20340
5 EC (152) , ER (152) , SMLQ (152) , SMLH (152) , BIGA (152) , KAP20350
6 BIGD (152) , CV (152) , BC (152) , BR (152) , CHIC (152) , KAP20360
7 CHIR (152) , CAPAC (152) , CAPAH (152) , CRTC (152) , CRTR (152) , KAP20370
8 CRTPC (152) , GOFK (152) , FEW (152) , CAR (152) , OKLM ( 37) , KAP20380
COMMON TELM ( 37) , EKLM ( 37) , ELM ( 37) , FCLM ( 37) , KAP20390
1 FRLM ( 37) , WLM ( 37) , QLM ( 37) , AMASNO( 37) , CHRNO ( 37) , KAP20400
2 ZP1 ( 37) , ZP2 ( 37) , SOLID ( 37) , ECHCK ( 37) , RK (104) , KAP20410
3 RL ( 37) , RHOK (104) , RDK (104) , THETAK(104) , TEMP ( 16) , KAP20420
4 HEAD ( 12) , MAXL , MAXLM , KAP20430
KAP20440
**KAP20450
*****KAP20460
C DATA FIRST/0./
IF(FIRST) 40,10,40
10 IF (RDK(1)* RDK(2) .GT. 0.1) GO TO 20
S1=74.0010
CALL UNCLE
20 IL=RUK(1)
IR=RUK(2)
SUMG=0.
DO 30 I=IL,IR

```

```
30 SUMG=SUMG+G(I)
40 DO 50 I=1,152
50 SMLQ(I)=0.
      T2= TH + DTH2
      IF(T2 .LT. RDK(3)) GO TO 1000
      DO 60 J=4,52
      IF(RDK(J) .GT. T2) GO TO 70
      IF(RDK(J) .EQ. 0.) GO TO 1000
60 CONTINUE
70 J= J-1
      SOLID(8)=RDK(J+50)
      IF(RDK(J) .LE. TH) GO TO 80
      TEMP(1)=(RDK(J)-TH)*RDK(J+49)+RDK(J+50)*(T2-RDK(J))
      SOLID(8)= TEMP(1)/UTH2
80 CNVRT= CNVRT + SOLID(8)*DTH2
      TEMP(1) = SOLID(8) * .25873241 / SUMG
      DO 90 I=IL,IR
90 SMLQ(I) = TEMP(1) * G(I)
1000 FIRST= 1.
      RETURN
      END
```

WI FOR QUE9/JPB, QUE9/JPB, QUE9/JP1  
SUBROUTINE QUE9

```
C*****
C
C*          S P U T T E R   C O M M O N
C*          ****
C*
COMMON LMUA(37), NR      , NSMLR , IA      , IB      , ICA      , ICB      ,
1 KMAX , BLANK1, BLANK2, BLANK3, IAP1 , IBP1 , ICAP1 , ICBP1 ,
2 II    , IG     , NRAU , BLANK4, IAM1 , IBM1 , ICAM1 , ICBM1 ,
3 IIP1 , IGM1 , IALPHA, BLANK5, TH    , TMAX , BLANK6, DELPRT,
4 FREQ , CNTMAX, AR     , ASMLR , PUSHA , PUSHB , BOILA , BOILB ,
5 CVA   , CVB   , SLUG  , ALPHA  , HVA   , HVB   , HCA   , HCB   ,
6 EMINA , EMINB , CA     , CB     , GA     , GB     , GL     , GR     ,
7 RHOL , RHOR , EPIO   , EPSI   , RIA   , RIB   , RDIA , RDIB ,
8 RPIA , RPIB , RPDIA , RPDIB , TPRINT, TA    , TB     , TC     ,
COMMON TD    , TE    , DTH2 , DTH2P , DTH1 , DTRMIN, DTMAX ,
1 DTMAX1, DTMAX2, DTMAX3, DTR , SWITCH, CO    , CMIN   , DELTA  ,
2 GAMA  , WCRIT , SIGMAQ, AC    , AC03T4, CNVRT , SUMRA , SUMRB ,
3 ROI A , HOIAM1, ROIH , ROIUP1, GMS   , S1     , S2     , S3     ,
4 S4    , S5    , S6    , S7    , S8    , S9    , S10   , S11   ,
5 S12   , S13   , S14   , S15   , S16   , S17   , S18   , S19   ,
6 S20   , EO    , FO    , TAU   , ZERO  , R    (152), DELTAR(152),
7 ASQ  (152), RD    (152), VD    (152), RDD   (152), SMLR (152),
8 DELR ( 37), P     (152), P1    (152), PB    (152), PB1   (152)
COMMON P2    (152), SV    (152), RHO   (152), THETA (152),
1 W    (132), E     (152), EI    (152), EK    (152), A     (152),
2 V    (152), G     (152), D     (152), C     (152), X2    (152),
3 X3   (152), X4    (152), X5    (152), X6    (152), X7    (152),
4 SMLA (152), SMLB (152), SMLC (152), SMLD (152), SMLE (152),
5 EC    (152), ER    (152), SMLQ (152), SMLH (152), BIGA (152),
6 BIGB (152), CV    (152), BC    (152), BR    (152), CHIC (152),
7 CHIR (152), CAPAC (152), CAPAR (152), CRT C (152), CRT R (152),
8 CRTPC (152), GOFR (152), FEW   (152), CAR   (152), OKLM ( 37)
COMMON TELM ( 37), EKLM ( 37), ELM   ( 37), FCLM ( 37),
1 FRLM ( 37), WLM  ( 37), QLM   ( 37), AMASNO( 37), CHRNO ( 37),
2 ZP1  ( 37), ZP2  ( 37), SOLID ( 37), ECHCK ( 37), RK    (104),
3 KL    ( 37), RHOK (104), RDK   (104), THETAK(104), TEMP  ( 16),
4 HEAD ( 12), MAXL      , MAXLM
C
C*
C*****
```

DATA FIRST/0./  
NR=RDK(101)  
IF(FIRHST.GT..5) GO TO 25  
DO 20 K=1,N  
IL=RDK(K)  
IR=RDK(K+1)-1.  
RDK(94+K)=0.  
DO 10 I=IL,IR  
10 RDK(94+K)=RDK(94+K)+G(I)  
20 CONTINUE  
25 DO 30 I=1,152  
30 SMLQ(I)=0.

```

DO 99 K=1,N
T2=TH+DTH2
IF(T2.LT. 1.E-20) GO TO 100
M=10+(K-1)*12
IF(T2.LT.RDK(M)) GO TO 99
M1 = M + 1
M5 = M + 5
DO 40 J = M1, M5
L = J
IF(RDK(J).GT. T2) GO TO 50
40 CONTINUE
GO TO 99
50 SOLID(8)=RDK(L+5)
IF (RDK(L-1) .LT. TH) GO TO 60
Q = (RDK(L-1) - TH) * RDK(L+4)
IF (L .EQ. M1) Q = 0.
TEMP(2) = (T2 - RDK(L-1)) * RDK(L+5) + Q
SOLID(8)=TEMP(2)/DTH2
60 CNVRT=CNVRT+SOLID(8)*DTH2
TEMP(1)=.23873241*SOLID(8)/RDK(94+K)
IL=RDK(K)
IR=RDK(K+1)-1.
DO 70 I=IL,IR
70 SMLQ(I)=TEMP(1)*G(I)
99 CONTINUE
100 FIRST=1.
RPIB=CNVRT
RETURN
C
C
C      RDK(1)-RDK(7) CONTAIN HORY OF SOURCE REGIONS
C      RDK(10)-RDK(15) CONTAIN TIME CUTS FOR FIRST GROUP
C      RDK(16)-RDK(21) CONTAIN RATES FOR FIRST GROUP
C
C      RDK(22)-RDK(27) - TIME    2ND GROUP
C      RDK(28)-RDK(33)  RATES
C      RDK(34)-RDK(39)  TIME    3RD GROUP
C      RDK(40-45)   RATES
C
C      RDK(46-51)  TIMES  4TH GROUP
C      RDK(52-57)  RATES
C
C      RDK(58-63)  TIMES  5TH GROUP
C      RDK(64-69)  RATES
C
C      RDK(70-75)  TIMES  6TH GROUP
C      RDK(76-81)  RATES
C
C      RDK(95-100) CONTAIN MASS OF EACH GROUP
C      RDK(101) NO. OF MATERIAL GROUPS
C
C      K IS SOURCE REGION INDEX
C      J, L ARE TIME CUT INDICES
C      M(K) IS THE RDK TABLE INDICATOR

```

C.V.EAT. IF TIME EXCEEDS LAST CUT, SOURCE SET ZERO.  
END

Q,T FOR QUE10/OUT, QUE10/OUT, QUE10/OUT1

SUBROUTINE QUE10

C~ MPILED JULY 18, 1967 JEZ

C~ SPECIAL SOURCE FOR J. PALMER

C~ SOURCE VARIES AS R\*\*-N

C~ RDK(1)= LEFT ZONE BOUNDARY

C~ RDK(2)= RIGHT ZONE BOUNDARY

C~ RDK(3-52)= TIMES FOR GIVEN EDOT

C~ RDK(53-104)= EDOTS FOR GIVEN TIMES

C\*\*\*\*\*KAP20070

C\*\*\*\*\*KAP20080

C\*\*\*\*\*KAP20090

C\*\*\*\*\*KAP20100

C\*\*\*\*\*KAP20110

SPUTTER COMMON

COMMON LMDA(37), NR , NSMLR , IA , IB , ICA , ICB , KAP20120  
 1 KMAX , BLANK1, BLANK2, BLANK3, IAP1 , IBP1 , ICAP1 , ICBP1 , KAP20130  
 2 II , IS , NRAD , BLANK4, IAM1 , IBM1 , ICAM1 , ICBM1 , KAP20140  
 3 IIP1 , IGM1 , IALPHA, BLANK5, TH , TMAX , BLANK6, DELPRT , KAP20150  
 4 FREQ , CNTMAX, AR , ASMLR , PUSH , PUSHB , ROILA , BOILB , KAP20160  
 5 CVA , CVB , SLUG , ALPHA , HVA , HVB , HCA , HCB , KAP20170  
 6 EMINA , EMINB , CA , CB , GA , GB , GL , GR , KAP20180  
 7 RHOL , RHOR , EPI0 , EPSI , RIA , RIB , RDIA , RDIB , KAP20190  
 8 RPIA , RPIB , RPUIA , RPDIB , TPRINT , TA , TB , TC , KAP20200  
 COMMON TU , TE , DTH2 , UTH2P , DTH1 , DTRMIN , DTMAX , KAP20210  
 1 DTMX1, DTMX2, DTMX3, DTR , SWITCH, CO , CMIN , DELTA , KAP20220  
 2 GAMA , WCRIT , SIGMAG, AC , AC03T4, CNVRT , SUMRA , SUMRB , KAP20230  
 3 RO1A , RO1AM1, ROIB , ROIWP1, GMS , S1 , S2 , S3 , KAP20240  
 4 S4 , S5 , S6 , S7 , S8 , S9 , S10 , S11 , KAP20250  
 5 S12 , S13 , S14 , S15 , S16 , S17 , S18 , S19 , KAP20260  
 6 S20 , EO , FO , TAU , ZERO , R (152) , DELTAR(152) , KAP20270  
 7 ASU (152) , RD (152) , VD (152) , RDD (152) , SMLR (152) , KAP20280  
 8 DELR ( 37) , P (152) , P1 (152) , PB (152) , PB1 (152) , KAP20290  
 COMMON P2 (152) , SV (152) , RHO (152) , THETA (152) , KAP20300  
 1 W (152) , E (152) , EI (152) , EK (152) , A (152) , KAP20310  
 2 V (152) , G (152) , D (152) , C (152) , X2 (152) , KAP20320  
 3 X3 (152) , X4 (152) , X5 (152) , X6 (152) , X7 (152) , KAP20330  
 4 SMLA (152) , SMLB (152) , SMLC (152) , SMLD (152) , SMLE (152) , KAP20340  
 5 EC (152) , ER (152) , SMLQ (152) , SMLH (152) , BIGA (152) , KAP20350  
 6 BIGB (152) , CV (152) , BC (152) , BR (152) , CHIC (152) , KAP20360  
 7 CHIR (152) , CAPAC (152) , CAPAR (152) , CRTG (152) , CRTR (152) , KAP20370  
 8 CRTPC (152) , GOFR (152) , FEW (152) , CAR (152) , OKLM ( 37) , KAP20380  
 COMMON TELM ( 37) , EKLM ( 37) , ELM ( 37) , FCLM ( 37) , KAP20390  
 1 FRLM ( 37) , WLH ( 37) , QLM ( 37) , AMASNO( 37) , CHRNO ( 37) , KAP20400  
 2 ZP1 ( 37) , ZP2 ( 37) , SOL10 ( 37) , ECHCK ( 37) , RK (104) , KAP20410  
 3 RL ( 37) , KHOK (104) , RDK (104) , THETAK(104) , TEMP ( 16) , KAP20420  
 4 HEAD ( 12) , MAXL , MAXLM , KAP20430  
 KAP20440  
 \*\*\*\*\*KAP20450  
 \*\*\*\*\*KAP20460

DATA FIRST/0./

IF(FIRST) 40,10,40

10 IF (RDK(1)\* RDK(2) .GT. 0.1) GO TO 20

S1=74,0010

CALL UNCLE

```

20 IL=RDK(1)
IR=RDK(2)
SUMG=0.
DO 30 I=IL,IR
30 SUMG=SUMG+G(I)
SUMG = SUMG-.5*G(IR)
40 DO 50 I=I,152
50 SMLQ(I)=0.
T2= TH + DTH2
IF(T2 .LT. RDK(3)) GO TO 1000
DO 60 J=4,52
IF(RDK(J) .GT. T2) GO TO 70
IF(RDK(J) .EQ. 0.) GO TO 1000
60 CONTINUE
70 J= J-I
TEMP(2) = RDK(J+40)
SOLID(8)=RDK(J+50)
IF(RDK(J) .LE. TH) GO TO 80
TEMP(1)=(RDK(J)-TH)*RDK(J+49)+RDK(J+50)*(T2-RDK(J))
TEMP(2) = (RDK(J)-TH)*RDK(J+39)+RDK(J+40)*(T2-RDK(J))
TEMP(2) = TEMP(2)/DTH2
SOLID(8)= TEMP(1)/DTH2
80 BETA = 4.18879*SUMG*TEMP(2)/SOLID(8)-I.
SUMQ = 0.
SUMS = -.5*G(IL-1)
DO 90 I=IL,IR
SUMS = SUMS+.5*(G(I-1)+G(I))
TEMP(I) = TEMP(2)*(SUMS/SUMG)**BETA
SMLQ(I) = TEMP(I)*G(I)
90 SUMQ = SUMQ+SMLQ(I)*4.18879
RPIB = SUMQ-SOLID(8)
DO 95 I=IL,IR
95 SMLQ(I) = SMLQ(I)*(SOLID(8)/SUMQ)
1000 FIRST= 1.
RETURN
END

```

QIT FOR KAPPA/ST, KAPPA/ST, KAPPA/STD  
 SUBROUTINE KAPPA(I,IM)  
 C OUTPUT VERSION THAT INCLUDES DIVIDE CHECK TEST  
 COMPILED MAY 22, 1967 KBL  
 C  
 C COMPILED ON DECEMBER 8, 1966, BY LAURA NORRIS

```

C
C
C+***** KAPP0020
C+***** KAPP0040
C+***** KAPP0050
C+***** KAPP0060
C+***** KAPP0070
C+***** KAPP0080
C+***** KAPP0090
C+***** KAPP0100
C+
C S P U T T E R   C O M M O N
C+
COMMON LMDA(37), NR , NSMLR , IA , IB , ICA , ICB , KAPP0110
1 KMAX , BLANK1, BLANK2, BLANK3, IAP1 , IBP1 , ICAP1 , ICBP1 , KAPP0120
2 II , IG , NRAD , BLANK4, IAM1 , IBM1 , ICAM1 , ICBM1 , KAPP0130
3 IIP1 , IGM1 , IALPHA, BLANK5, TH , TMAX , BLANK6, DELPRT, KAPP0140
4 FREQ , CNTMAX, AR , ASMLR , PUSHA , PUSHB , BOILA , BOILB , KAPP0150
5 CVA , CVB , SLUG , ALPHA , HVA , HVB , HCA , HCB , KAPP0160
6 EMINA , EMINB , CA , CB , GA , GB , GL , GR , KAPP0170
7 RHOL , RHOR , EPIO , EPSI , RIA , RIB , RDIA , RDIB , KAPP0180
8 RPIA , RPIB , RPDIAB , RPDIIB , TPRINT, TA , TB , TC , KAPP0190
COMMON TO , TE , DTH2 , DTH2P , DTH1 , DTRMIN, DTMAX , KAPP0200
1 DTMX1, DTMX2, DTMX3, DTR , SWITCH, CO , CMIN , DELTA , KAPP0210
2 GAMMA , WCRT , SIGMAQ, AC , AC03T4, CNVRT , SUMRA , SUMRB , KAPP0220
3 RO1A , ROIAM1, KOIB , KOIBP1, GMS , S1 , S2 , S3 , KAPP0230
4 S4 , S5 , S6 , S7 , S8 , S9 , S10 , S11 , KAPP0240
5 S12 , S13 , S14 , S15 , S16 , S17 , S18 , S19 , KAPP0250
6 S20 , EO , FO , TAU , ZERO , R (152), DELTAR(152), KAPP0260
7 ASQ (152), RD (152), VD (152), ROD (152), SMLR (152), KAPP0270
8 DELR (37), P (152), PI (152), PB (152), PB1 (152) KAPP0280
COMMON P2 (152), SV (152), RHO (152), THETA (152), KAPP0290
1 W (152), E (152), EI (152), EK (152), A (152), KAPP0300
2 V (152), G (152), D (152), C (152), X2 (152), KAPP0310
3 X3 (152), X4 (152), X5 (152), X6 (152), X7 (152), KAPP0320
4 SMLA (152), SMLB (152), SMLC (152), SMLD (152), SMLE (152), KAPP0330
5 EC (152), ER (152), SMLQ (152), SMLH (152), BIGA (152), KAPP0340
6 BIGB (152), CV (152), BC (152), BR (152), CHIC (152), KAPP0350
7 CHIK (152), CAPAC (152), CAPAR (152), CRTA (152), CRTR (152), KAPP0360
8 CRTPC (152), GOFR (152), FE# (152), CAR (152), OKLM (37) KAPP0370
COMMON TELM (37), EKLM (37), ELM (37), FCLM (37), KAPP0380
1 FRLM (37), #LM (37), QLM (37), AMASNO(37), CHRNO (37), KAPP0390
2 ZP1 (37), ZP2 (37), SOLID (37), ECHCK (37), RK (104), KAPP0400
3 RL (37), RHOK (104), ROK (104), THETAK(104), TEMP (16), KAPP0410
4 HEAU (12), MAXL , MAXLM
C
C+***** KAPP0420
C+***** KAPP0430
C+***** KAPP0440
C+***** KAPP0450
C+***** KAPP0460
C+***** KAPP0470
C+***** KAPP0480
C+
C COMMON /LINLTY/ HNU,SGNL,IHNU,IHNU,HNUP,NT,IM,IN,DHNU,THICK,NY KAPP0460
C+
C AMASNO(18-34) RESERVED FOR DESIGNATION OF DIANE TAPE UNIT.
C

```

```

C      NT=0
C      SGNL=0.
C
C      M = IM
C      IL = N
100  DO 200 J=1, MAXL.
     IF(IL.LT.LMDA(J)) GO TO 300
200  CONTINUE
     S1 = 15.0200
     CALL UNCLE
300  IR=MINU(LMDA(J)-1,M)
     JE J - 1
     IF(QLM(J+17))400,350,400
350  QLM(J+17)=ABS(OKLM(J))
400  CONTINUE
     SOLIO(20)= IL
     SOLID(21)=IR
     SOLIO(22)=J
     IF(ANASH0(J+17) .GT. 0.) GO TO 600
     L = 2
     MATERL=QLM(J+17) + .5
     IF(MATERL .GT.200) L=MATERL - 200
     IF(MATERL .EQ. 1) L=3
     IF(MATERL .EQ. 6) L=1
     IF(MATERL .EQ. 13) L=5
        IF (MATERL .EQ. 101 .OR. MATERL .EQ. 103) L=I
     IF(MATERL .EQ. 102)L=8
     IF(L .LT. 13) GO TO 500
     S1= 15.0500
     CALL UNCLE
500  GO TO (1,2,3,4,5,6,7,8,9,10,11,12),L
     1 CALL KAP1
        GO TO 700
     2 CALL KAP2
        GO TO 700
     3 CALL KAP3
        GO TO 700
     4 CALL KAP4
        GO TO 700
     5 CALL KAP5
        GO TO 700
     6 CALL KAP6
        GO TO 700
     7 CALL KAP7
        GO TO 700
     8 CALL KAP8
        GO TO 700
     9 CALL KAP9
        GO TO 700
    10 CALL KAP10
        GO TO 700
    11 CALL KAP11
        GO TO 700
     12 CALL KAP12
        GO TO 700
600  CALL UIAHA(J)
700  DO 800 I=IL,IR
     CAPAR(I) = AMINI(CAPAK(I), 2.620E6/THTA(I))
800  CONTINUE
     CALL UVCHK(KX)
     GO TO (810, 820), KX
810  S1 = 15.0810
     CALL UNCLE
820  IL=IR + 1
     IF (IL .GT. M) RETURN
     GO TO 100
END

```

KAPP1170  
KAPP1190  
KAPP1210  
KAPP1230  
KAPP1250  
KAPP1270  
KAPP1290  
KAPP1310  
KAPP1330  
KAPP1350  
KAPP1370  
KAPP1390

```

C, FOR KAP6/JP, KAP6/JP, KAP6/JPI
SUBROUTINE KAP6
COMPILED NOVEMBER 11, 1966 ABL
C      SPECIAL KAP ROUTINE FOR CONSTANT INPUT MULTIFREQUENCY OPACITIES
C      THETAK(61-80) SPECIFIES FREQUENCY
C      THETAK(61-100) SPECIFIES OPACITY
C,*****          *****          *****          *****          *****
C
C          S P U T T E R   C O M M O N
C
C
COMMON LMDA(37), NH    , NSMLR , IA    , IB    , ICA   , ICB   ,
1  KMAX , BLANK1, BLANK2, BLANK3, IAP1 , IBP1 , ICAP1 , ICBP1 ,
2  II   , IG   , NRAD , BLANK4, IAM1 , IBM1 , ICAMI , ICBM1 ,
3  IIP1 , IGM1 , IALPHA, BLANK5, TM   , TMAX , BLANK6, DELPRT,
4  FREQ  , CNTMAX, AR   , ASMLR , PUSHA , PUSHB , HOILA , BOILH ,
5  CVA   , CVB  , SLUG  , ALPHA  , HVA   , HVR   , HCA   , HCB   ,
6  EMINA , EMINB , CA   , CB   , GA   , GB   , GL    , GR    ,
7  RHOL , RHOR , EPIO , EPSI , RIA  , RID  , RDIA , RDIB ,
8  RPIA , RPIB , RPDIA , RPDIB , TPRINT, TA   , TB    , TC    ,
COMMON TU   , TE   , UTH2 , UTH2P , UTH1 , DTRMIN, DTMAX ,
1  UTMAX1, UTMAX2, DTMAX3, UTR  , SWITCH, CO   , CMIN  , DELTA ,
2  GAMA  , WCRI1 , SIGMA1, AC   , AC03T4, CNVRT , SUMRA , SUMRB ,
3  ROI1 , ROIAM1, HOIB , HOIHP1, GMS  , S1   , S2   , S3   ,
4  S4   , S5   , S6   , S7   , S8   , S9   , S10  , S11  ,
5  S12  , S13  , S14  , S15  , S16  , S17  , S10  , S19  ,
6  S20  , EO   , FO   , TAU  , ZERO , R    , (152), DELTAR(152),
7  ASQ  , (152), NU   , (152), VD  , (152), RDD  , (152), SMLR (152),
8  UELK , (37), P   , (152), P1  , (152), PB   , (152), PB1  , (152)
COMMON P2   , SV   , (152), RHO , (152), THETA (152),
1  W   , (152), E   , (152), EI   , (152), EK   , (152), A   , (152),
2  V   , (152), G   , (152), D   , (152), C   , (152), X2  , (152),
3  X3   , X4   , (152), XS   , (152), X6   , (152), X7  , (152),
4  SHLA , SMLB , (152), SMLC , (152), SMLD , (152), SMLE (152),
5  EC   , ER   , (152), SMLQ , (152), SMLH , (152), BIGA (152),
6  BIGB , CV   , (152), BC   , (152), BR   , (152), CHIC (152),
7  CHIR , CAPAC, (152), CAPAH, (152), CRTG , (152), CRTR (152),
8  CRTPC , GOFR , (152), FEW  , (152), CAR  , (152), OKLM ( 37),
COMMON TELM , (37), EKLM , (37), ELM  , (37), FCLM , (37),
1  FRLM , (37), ULM , (37), QLM , (37), AMASNO( 37), CHRNO ( 37),
2  ZP1  , (37), ZP2 , (37), SOLID , (37), ECHCK , (37), RK  , (104),
3  NL   , (37), NMOK , (104), RDK  , (104), THETAK(104), TEMP  , (16),
4  HEAD , (12), MAXL , , MAXLM
C
C
COMMON /LINOLY/ HNU,SGNL,INNU,NHNU,HNUP,NT,IM,IN,DHNU,THICK,NY
IF (INNU .EQ. 1) NHNU = THETAK(101)
NHNU = THETAK(INNU+60)
IK = SOLID(20)
IL = SOLID(21)
DO 10 I = IK, IL
IF (INNU .GT. 0) GO TO 6
C      GREY KAPPA
C
CAPAH(1) = THETAK(102)
GO TO 10
C      MULTIFREQUENCY KAPPA
6  CAPAH(I) = THETAK((INNU+80))
10 CAPAC(I) = CAPAH(I)
RETURN
END

```

W,T FOR KAP12/JP, KAP12/JP, KAP12/JP1  
SUBROUTINE KAP12

KAP10010  
KAP10020  
ES3 0040

C  
C  
C  
C  
C

## S P U T T E R C O M M O N

```

COMMON LMDA(37), NR      , NSMLR , IA      , IB      , ICA      , ICB      ,
1  KMAX   , BLANK1, BLANK2, BLANK3, IAP1   , IBP1   , ICAP1   , ICBP1   ,
2  II     , IG     , NRAD   , BLANK4, IAM1   , IBM1   , ICAM1   , ICBM1   ,
3  IIP1   , IGN1   , IALPHA , BLANK5, TH     , TMAX   , BLANK6, DELPRT,
4  FREQ   , CNTMAX, AR     , ASMLR  , PUSHA  , PUSHB  , BOILA  , BOILB  ,
5  CVA    , CVB    , SLUG   , ALPHA   , HVA    , HVB    , HCA    , HC8     ,
6  EMINA  , EMINB  , CA     , CB     , GA     , GH     , GL     , GR     ,
7  RHOL   , RHOK   , EPI0   , EPSI   , RIA    , RIB    , ROIA   , ROIB   ,
8  RPIA   , RPIB   , RPDIA  , RPDIB  , TPRINT, TA    , TB     , TC     ,
COMMON TD     , TE     , UTH2   , DTH2P  , DTH1   , DTRMIN, DTMAX  ,
1  DTMAX1, DTMAX2, DTMAX3, UTR    , SWITCH, CO    , CMIN   , DELTA   ,
2  GAMMA  , ACINIT , SIGMAQ , AC    , ACO3T4, CNVRT , SUMRA  , SUMRB  ,
3  ROIA   , KOIAM1 , KOIB   , KOIBP1, GMS    , S1     , S2     , S3     ,
4  S4     , S5     , S6     , S7     , S8     , S9     , S10   , S11   ,
5  S12   , S13   , S14   , S15   , S16   , S17   , S18   , S19   ,
6  S20   , EO     , FU     , TAU   , ZERO   , R     (152), DELTAR(152),
7  ASQ   (152), RD     (152), VD   (152), RDD   (152), SMLR  (152),
8  DELK   (37), P     (152), P1   (152), PB    (152), PB1   (152)
COMMON P2   (152), SV     (152), RHO   (152), THETA (152),
1  W     (152), E     (152), EI    (152), EK    (152), A     (152),
2  V     (152), G     (152), D     (152), C     (152), X2    (152),
3  X3   (152), X4   (152), X5   (152), X6   (152), X7    (152),
4  SMLA  (152), SMLB  (152), SMLC (152), SMLD (152), SMLE (152),
5  EC    (152), ER    (152), SMLQ (152), SMLH (152), BIGA (152),
6  HIGH  (152), CV    (152), BC    (152), BH    (152), CHIC (152),
7  CHIR  (152), CAPAC (152), CAPAR (152), CRTA (152), CRTR (152),
8  CRTPC (152), GUFR  (152), FEW   (152), CAR   (152), OKLM (37)
COMMON TELM (37), EKLM (37), ELM   (37), FCLM (37),
1  FRLM (37), WLM  (37), QLM  (37), AMASNO (37), CHR10 (37),
2  ZP1  (37), ZP2  (37), SOLID (37), ECHCK (37), RK    (104),
3  HL   (37), RHOK (104), ROK   (104), THETAK(104), TEMP  (16),
4  HEAD (12), MAXL   , MAXLM
IL=SOLID(21)
```

KAP10050

C  
C  
C

```

IK=SOLID(20)
J=SOLID(22)
CHRN=92.0
AMASH=4.017E+9
ZPA=2.408E-03
ZPB=3.821E+04
DO 1000 I=IK,IL
IF(THTETA(I).NE.0.) GO TO 1
```

KAP10030  
KAP10040  
KAP10060  
KAP10070  
KAP10080  
KAP10090  
KAP10100  
KAP10110

```

CAPAR(1)=.2
CAPAC(1)=.2
GO TO 1000
1 IAT=1
TEMP(3)=THETA(1)*2
TEMP(4)=TEMP(3)**2
TEMP(5)=THETA(1)**1.5
IF(SV(I).GE.10. ) GO TO 200
IF(SV(I).LE.1. ) GO TO 10
IXT=2
SVI=SV(I)
SV(I)=10.
DYDY=ALOG( 10./SVI)/( -2.3025851)
GO TO 200
5 SV(I)=1.
TEMP(6)=CAPAR(1)
10 IF(THETA(1)-1000.0)12,12,14
12 CAPAR(1)=5.49E+6/(TEMP(5)*(1.0+1.E-6*SV(I)*TEMP(3)))+13.1*
1 SQRT(SV(I))
CAPAR(1) = AMAX1(OKLM(J+17) * CAPAR(1), .2)
GO TO 20
14 CAPAR(1)=(1.0+1.75E12/TL*P(4))*(.0789+2.784E21/(TEMP(3)*(7.6E13*P(4)*SV(I))))
1 SQRT(SQRT(SV(I)))+TEMP(4)*SV(I)))
CAPAR(1) = AMAX1(OKLM(J+17) * CAPAR(1), .2)
20 GO TO (350,340)+IXT
200 CONTINUE
IF (THETA(I)-2.0) 201,201,202
201 EFF=1.E2
GO TO 300
202 IF (THETA(I)-20.0) 203,203,204
203 EFF=9.E2/(THETA(I)+1.0)**2+2.5E-5*THETA(I)**5
GO TO 300
204 IF (THETA(I)-100.0) 205,205,206
205 EFF=9.56E5/(THETA(I)**3+5.E-5*THETA(I)**5)
GO TO 300
206 EFF=0.0+4.4E5/THETA(I)**2
301 CONTINUE
WS=SQRT(THETA(I))
WSA = FEW(I)*2/SV(I)
IF(FLA(I).LT.1.) GO TO 330
WSA=WSA*FEW(I)
GO TO 335
330 WSA = AMIN1(WSA, 1.E10/AMASN*THETA(I)+WS*EXP(-8.0/THETA(I)))
335 CAPAR(1) = AMASN**2*9.06E-14/THETA(I)**3*WSA/WS
CAPAR(1)=AMAX1(OKLM(J+17)*EFF*CAPAR(1),.2)
GO TO (350,5)+IXT
340 CAPAR(1)=EXP(ALOG(TEMP(6))+ALOG(CAPAR(1))-ALOG(TEMP(6))+DYDY)
SV(I)=SVI
350 NSB=13.6/THETA(I)+(CHRN**2)
IF(FEW(I).EQ.0.) GO TO 370
360 NSC=1.E+10/AMASN*TEMP(5)/FEW(I)*SV(I)
GO TO 380
370 NSC= 1.
380 CONTINUE
CAPAC(1)=.8/*CAPAR(1)+AMIN1(1.0+NSB,1.0+ALOG(NSC))
CAPAC(1)=AMAX1(CAPAC(1),CAPAR(1))
.000 CONTINUE
RETURN
END

```

KAP10200  
KAP10210  
KAP10220  
KAP10240  
KAP10260  
KAP10270  
KAP10280  
KAP10300  
KAP10310  
KAP10320  
KAP10410  
KAP10420  
KAP10430  
KAP10440  
KAP10450  
KAP10460  
KAP10470  
KAP10480  
KAP10490  
KAP10500  
KAP10510  
KAP10600  
KAP10610  
KAP10630  
KAP10640  
KAP10650  
KAP10660  
KAP10670

Q FOR RTAPE/ORIG, RTAPE/ORIG, RTAPE/A1  
SUBROUTINE RTAPE(INTAPE,CYCLE)

RTAP0010

C

C COMPILED APRIL 4, 1967 WBL

C SPECIAL VERSION FOR COMPTON SCATTERING CODES

C\*\*\*\*\*RTAP0020

C\*\*\*\*\*RTAP0030

C\*\*\*\*\*RTAP0040

C\*\*\*\*\*RTAP0050

C\*\*\*\*\*RTAP0060

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COMMON  LMDA(37), NR    , NSMLR , IA    , IB    , ICA    , ICB    , RTAP0070
1   KMAX  , BLANK1, BLANK2, BLANK3, IAP1  , IBP1  , ICAP1  , ICBP1 , RTAP0080
2   II    , IG    , NRAD  , BLANK4, IAM1  , IBM1  , ICAM1  , ICAM1 , RTAP0090
3   IIP1  , IGM1  , IALPHA, BLANK5, TH    , TMAX  , BLANK6, DELPR, RTAP0100
4   FREQ  , CNTMAX, AR    , ASMLR , PUSHA , PUSHB , HOILA , BOILB , RTAP0110
5   CVA   , CVB   , SLUG  , ALPHA  , HVA   , HVB   , HCA   , HCB   , RTAP0120
6   EMINA , EMINB , CA    , CB    , GA    , GB    , GL    , GR    , RTAP0130
7   RHOL  , RHOK  , EPIO  , EPSI  , RIA   , RIB   , RDIA  , RDIB  , RTAP0140
8   RPIA  , RPIB  , RPDIA , RPDIB , TPRINT, TA    , TB    , TC    , RTAP0150
COMMON  TU    , TE    , UTH2  , DTH2P , DTH1  , DTRMIN, DTMAX , RTAP0160
1   DTMAX1, DTMAX2, DTMAX3, DTR   , SWITCH, CO    , CMIN  , DELTA  , RTAP0170
2   GAMA  , WCRIT , SIGMAJ, AC    , AC03T4, CNVPT , SUMRA , SUMRB , RTAP0180
3   ROIA  , ROIAMI, HOIH  , HOIBP1, GMS   , S1    , S2    , SJ    , RTAP0190
4   S4    , S5    , S6    , S7    , S8    , S9    , S10   , S11   , RTAP0200
5   S12   , S13   , S14   , S15   , S16   , S17   , S18   , S19   , RTAP0210
6   S20   , EU    , FU    , TAU   , ZERO  , R    (152), DELTAR(152), RTAP0220
7   ASQ   (152), RU   (152), VD   (152), ROD  (152), SMLR (152), RTAP0230
8   DELR  (37), P    (152), P1   (152), PB   (152), PB1  (152) RTAP0240
COMMON  P2    (152), SV    (152), RHO  (152), THETA (152), RTAP0250
1   N    (152), E    (152), EI   (152), EK   (152), A    (152), RTAP0260
2   V    (152), G    (152), D    (152), C    (152), X2   (152), RTAP0270
3   X3   (152), X4   (152), X5   (152), X6   (152), X7   (152), RTAP0280
4   SMLA (152), SMLU (152), SMLC (152), SMLD (152), SMLE (152), RTAP0290
5   EC   (152), ER   (152), SMLQ (152), SMLH (152), BIGA (152), RTAP0300
6   BIGB (152), CV   (152), BC   (152), BH   (152), CHIC (152), RTAP0310
7   CHIR  (152), CAPAC (152), CAPAR (152), CRTC (152), CRTR (152), RTAP0320
8   CRTPC (152), GUFK (152), FEW  (152), CAR  (152), OKLM (37) RTAP0330
COMMON  TELM (37), EKLM (37), ELM  (37), FCLM (37), RTAP0340
1   FRLM (37), WLM (37), QLM (37), AMASNO(37), CHRNO (37), RTAP0350
2   ZP1  (37), ZP2 (37), SOLID (37), ECHCK (37), RK (104), RTAP0360
3   RL   (37), RHOK (104), RUK (104), THETAK(104), TEMP (16), RTAP0370
4   HEAD (12), MAXL , MAXLM , RTAP0380
RTAP0390
RTAP0400
C*****RTAP0410
```

C

C

```

COMMON /LINDLY/ MNNU,SGNL,INNU,NMNU,MNUP,NT,IM,IN,DMNU,THICK,NY
COMMON /PALMER/ FI0(152), FI1(152), FI2(152), FI3(152), F00(152),
2 F01(152), F02(152), F03(152), JDUM
```

C RTAPE DEFINES JDUM ON RESTARTS. WTape DEFINES IT ON INITIAL STARTS.

NMNU = LHOA(36)

REWIND NTAPE

RTAP0420

10 READ (INTAPE) COUNT

RTAP0430

IF (COUNT.LT.0..OR.COUNT.GT.CYCLE) GO TO 30

RTAP0440

IF (ABS(F0) .GT. 1.E-20) HEAD (INTAPE)

```
IF (ABS(COUNT - CYCLE) .LT. 1.E-20) GO TO 20
WRITE (6,50) COUNT
50 FORMAT (1HO 5X, 6HCYCLE F0.0, 8H SKIPPED)
READ (INTAPE)
DO 15 I = 1, NHNU
15 READ (INTAPE)
GO TO 10
10 READ (INTAPE) (LMUA(I),I=1,9405)
JURUM = 25
LMUA(36) = NHNU
DO 25 I = 1, NHNU
READ (INTAPE) F10, F11, F12, F13
25 WRITE (JURUM) F10, F11, F12, F13
RETURN
30 SI = 40.
WHITE (6, 40) SOLID(34)
40 FORMAT (1HO 5X 6HCYCLE F6.0, 17H NOT ON DUMP TAPE )
CALL EXIT
END
```

RTAP0460

RTAP0470

RTAP0480

RTAP0490

RTAP0500

RTAP0520

## AFWL TR 67-131, Vol III

Q FOR \*TAPE/SP, \*TAPE/SP, \*TAPE/SP1  
 SUBROUTINE \*TAPE  
 COMPILED APRIL 27, 1967 NOL  
 C SPECIAL VERSION FOR COMPTON SCATTERING CODES

WTAP0010

C \*\*\*\*=  
 C SPUTTER COMMON  
 C \*\*\*\*=  
 C  
 COMMON LMDA(37), NR, NSMLR, IA, IB, ICA, ICB, WTAP0030  
 1 KMAX, BLANK1, BLANK2, BLANK3, IAP1, IBP1, ICAPI, ICBP1, WTAP0100  
 2 II, IG, INRAD, BLANK4, IAM1, IBM1, ICAM1, ICBM1, WTAP0110  
 3 IIP1, IGM1, IALPHA, GLANK5, TH, TMAX, BLANK6, DELPRT, WTAP0120  
 4 FREQ, CNTMAX, AR, ASMLR, PUSHA, PUSHB, BOILA, BOILB, WTAP0130  
 5 CVA, CVB, SLUG, ALPHA, HVA, HVB, HCA, HCB, WTAP0140  
 6 EMINA, EMINB, CA, CB, GA, GB, GL, GR, WTAP0150  
 7 RHOL, RHOR, EPIO, EPSI, RIA, RIR, RDIA, RDIB, WTAP0160  
 8 RPRA, RPIB, RPDIA, RPDIB, TPRINT, TA, TB, TC, WTAP0170  
 COMMON TD, TE, DTH2, UTH2P, DTH1, DTRMIN, DTMAX, WTAP0180  
 1 DTMAX1, DTMAX2, DTMAX3, DTR, SWITCH, CO, CMIN, DELTA, WTAP0190  
 2 GAMA, WCRIT, SIGMAU, AC, AC03T4, CNVRT, SUMRA, SUMRB, WTAP0200  
 3 ROIA, ROIAMI, ROIB, ROIBP1, GMS, S1, S2, S3, WTAP0210  
 4 S4, S5, S6, S7, S8, S9, S10, S11, WTAP0220  
 5 S12, S13, S14, S15, S16, S17, S18, S19, WTAP0230  
 6 S20, EO, FO, TAU, ZERO, R, (152), DELTAR(152), WTAP0240  
 7 ASG (152), RD, (152), VD, (152), RUD, (152), SMLR, (152), WTAP0250  
 8 DELR (37), P, (152), PI, (152), PB, (152), PB1, (152), WTAP0270  
 COMMON P2, (152), SV, (152), RHO, (152), THETA, (152), WTAP0280  
 1 X, (152), E, (152), EI, (152), EK, (152), A, (152), WTAP0290  
 2 V, (152), G, (152), D, (152), C, (152), X2, (152), WTAP0300  
 3 X3, (152), X4, (152), X5, (152), X6, (152), X7, (152), WTAP0310  
 4 SMLA, (152), SMLH, (152), SMLC, (152), SMLD, (152), SMLE, (152), WTAP0320  
 5 EC, (152), ER, (152), SMLQ, (152), SMLH, (152), BIGA, (152), WTAP0330  
 6 BIGB, (152), CV, (152), UC, (152), BR, (152), CHIC, (152), WTAP0340  
 7 CHIR, (152), CAPAC, (152), CAPAR, (152), CRTC, (152), CRTR, (152), WTAP0350  
 8 CRTPC, (152), GOFH, (152), FEW, (152), CAR, (152), OKLM, (37), WTAP0360  
 COMMON TELM, (37), EKLM, (37), ELM, (37), FCLM, (37), WTAP0370  
 1 FRLM, (37), WLM, (37), QLM, (37), AMASHO, (37), CHRNO, (37), WTAP0380  
 2 ZP1, (37), ZP2, (37), SOLID, (37), ECHCK, (37), RK, (104), WTAP0390  
 3 RL, (37), RMOK, (104), RDK, (104), THETAK(104), TEMP, (16), WTAP0400  
 4 HEAU, (12), MAXL, MAXLM, WTAP0410  
 \*\*\*\*=  
 C  
 COMMON /LIMOLY/, MNJ, SGIL, INHU, NHNU, HNUP, NT, IM, IN, DMNU, THICK, NY  
 COMMON /PALMEK/, FIU(152), FI1(152), FI2(152), FI3(152), FOO(152),  
 2 FQ1(152), FQ2(152), FQ3(152), JORUM  
 COMMON /CNTML/, SCYCLE, JMULT  
 C  
 N=55  
 WS=3.0  
 WRITE (N) SOLID(18)  
 1F (AHS(F0) .GT. 1.E-20) WRITE (N) TH  
 \*\*\*\*=  
 WRITE (N)(LMDA(1), I=1, 9405)  
 NNU = LMDA(36)  
 1F (SOLID(18) .GT. SCYCLE) REIND JORUM  
 DO 10 I = 1, NNU  
 1F (SOLID(18) .GT. SCYCLE) HEAD (JORUM) FI0, FI1, FI2, FI3  
 10 WRITE (N) FI0, FI1, FI2, FI3  
 C 10 WRITE (N) FI0, FI1, FI2, FI3  
 \*TAPE DEFINES JNUH ON INITIAL STARTS. \*TAPE DOES IT ON RESTARTS.  
 IF (SOLID(34) .GT. 0. .OR. SOLID(18) .GT. 0.) GO TO 20  
 JORUM = 25  
 REWIN JORUM  
 DO 15 I = 1, NNU  
 15 WRITE (JNUH) FI0, FI1, FI2, FI3  
 20 WRITE (N) WS  
 BACKSPACE N  
 RETURN  
 END  
 \*\*\*\*=  
 WTAP0440  
 WTAP0450  
 WTAP0460  
 WTAP0470  
 WTAP0490  
 WTAP0500  
 WTAP0510  
 WTAP0560  
 WTAP0570

```

D FOR MP2/XX,MP2/XX,MP2/XXX
C SUBROUTINE MP2
C OUTPUT VERSION -- BOIL, CNDCTN DELETED
C
C*****S P U T T E R C O M M O N *****C
C
COMMON LMDA(37), NR , NSMLR , IA , IB , ICA , ICB , 3MAI0010
1 KMAX , BLANK1, BLANK2, BLANK3, IAP1 , IBP1 , ICAP1 , ICBP1 , 3MAI0030
2 II , IG , NRAD , BLANK4, IAM1 , IBM1 , ICAM1 , ICBM1 , 3MAI0040
3 IIP1 , IGM1 , IALPHA, BLANK5, TH , TMAX , BLANK6, DELPRT, 3MAI0120
4 FRLW , CNTMAX, AR , ASMLR , PUSHA , PUSHB , HOILA , BOILB , 3MAI0130
5 CVA , CVB , SLUG , ALPHA , HVA , HVB , HCA , HCB , 3MAI0140
6 EMINA , EMINU , CA , CH , GA , GB , GL , GR , 3MAI0150
7 RHOL , RHOK , EPIO , EPSI , RIA , RIB , RDIA , RDIB , 3MAI0160
8 RPIA , RPIB , RPDIA , RPDIB , TPRINT, TA , TB , TC , 3MAI0170
COMMON1 TU , TE , UTH2 , UTH2P , DTH1 , UTRMIN, DTMAX , 3MAI0180
1 DTMAX1, DTMAX2, DTMAX3, UTR , SWITCH, CO , CMIN , DELTA , 3MAI0190
2 GAMA , WCHIT , SIGMAQ, AC , AC03T4, CNVRT , SUMRA , SUMRB , 3MAI0200
3 ROIA , ROIAM1, ROIUB , ROIIP , GMS , S1 , S2 , S3 , 3MAI0210
4 S4 , S5 , S6 , S7 , S8 , S9 , S10 , S11 , 3MAI0220
5 S12 , S13 , S14 , S15 , S16 , S17 , S18 , S19 , 3MAI0230
6 S20 , EO , FU , FAU , ZERO , R (152), DELTAR(152), 3MAI0240
7 ASQ (152), RD (152), VD (152), RDD (152), SMLR (152), 3MAI0250
8 DELR (37), P (152), PI (152), PB (152), PB1 (152) 3MAI0260
COMMON P2 (152), SV (152), RHO (152), THETA (152), 3MAI0270
1 W (152), E (152), EI (152), EK (152), A (152), 3MAI0280
2 V (152), G (152), D (152), C (152), X2 (152), 3MAI0290
3 X3 (152), X4 (152), X5 (152), X6 (152), X7 (152), 3MAI0300
4 SMLA (152), SMLB (152), SMLC (152), SMLD (152), SMLE (152), 3MAI0310
5 EC (152), ER (152), SMLQ (152), SMLH (152), BIGA (152), 3MAI0320
6 BIGB (152), CV (152), BC (152), BR (152), CHIC (152), 3MAI0330
7 CHIR (152), CAPAC (152), CAPAR (152), CRTC (152), CRTR (152), 3MAI0340
8 CRTPC (152), GOFK (152), FEW (152), CAR (152), OKLM (37) 3MAI0350
COMMON TELM (37), EKLM (37), ELM (37), FCLM (37), 3MAI0360
1 FRLH (37), WLM (37), OLM (37), AMASNO(37), CHRNO (37), 3MAI0370
2 ZP1 (37), ZP2 (37), SOLID (37), ECHCK (37), RK (104), 3MAI0380
3 RL (37), RHOK (104), HDK (104), THETAK(104), TEMP (16), 3MAI0390
4 HEAD (12), MAXL , MAXLM , 3MAI0400
3MAI0410
3MAI0420
COMMON /CNTRL/ SCYCLE, JMULT 3MAI0430
EQUIVALENCE (S12,ED1TMF) 3MAI0440
HYDRO SUB-CYCLE VERSION 4/25/63 3MAI0450
3MAI0460
TO50, 3MAI0470
IF(BLANK2 .GT.1.E-15)TPRINT=TH + BLANK2 3MAI0480
3 IF (WARN(IHR)) 10, 600, 600 3MAI0490
3MAI0570

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```

10 IF (EKLM(19)) 501, 520, 501          3MAI0580
501 CONTINUE                            3MAI0590
    CALL AUTORZ                           3MAI0600
520 IF(GMS)530,201,530                  3MAI0610
530 CALL DEZONE                           3MAI0620
201 CONTINUE                            3MAI0630
    40 CALL TDELT
        TD=TU+1.
        IF(BLANK2 .LT. 1.E-15) GO TO 170
        IF(TU .GT. FREQ) GO TO 171
        IF(TPRINT .GT. TH+ DTH2) GO TO 100
        TPRINT=TPRINT + BLANK2
        IF(TH .GT. TPRINT) BLANK2 = 0.
        GO TO 171
170 IF(AMOD(SOLID(18)+1.,FREQ).GT. 0.5)GO TO 100
171 TU=0.
100 CONTINUE
C ***** INTEGRAL SUR                   3MAI0650
C *****                                     *3MAI0660
C *****                                     *3MAI0670
    IF(PUSHA) 2000,2010,2001
2000 IBD=IBM1
    GO TO 2010
2001 IBD = IA
2010 CONTINUE                            3MAI0720
    204 CALL HYDRO                           3MAI0730
    202 CALL PTWO                            3MAI0740
    3MAI0750
    3MAI0760
C
C      SELECT SOURCE
    IF (ABS(S6) .LT. 1.E-20) GO TO 80
    CALL QUE
C
80 CONTINUE
    DTH=DTH2
    ZP1(18)=DTH2
    ZP1(19)=0.0
84 CONTINUE
    IF (ABS(S2) .LT. 1.E-20) GO TO 88
C     CHECK TO SEE IF RADIATION NEEDS TO BE DONE ON CURRENT SUBCYCLE
85 IF(TD.LT.0.5 .OR. SOLID(18)+1.1.GT.CNTMAX)GO TO 86
    IF (BLANK3-TH-DTH2-DTR+ZP1(18)) 86,88,88
    3MAI1050
    86 CALL RADTN                           3MAI1060
    88 CONTINUE
        CALL ENCALC                           3MAI1110
        CALL ECALC                            3MAI1120
        NRAD=NRAD-1                          3MAI1130
        IF(NRAD.LE.0.0K.(ZP1(18)/DTH).LT.1.5) GO TO 110
105 CALL SSWTCH (6, K000FX)
    GO TO (106, 107), K000FX
106 S1 = 1.0106
    GO TO 610
107 CONTINUE
    IF (WARN(IHR)) 108, 605, 605
108 ZP1(18)=ZP1(18)-DTH
    ZP1(19)=ZP1(19)+1.0
    3MAI1150
    3MAI1160
    3MAI1170
    3MAI1180
    3MAI1190
    3MAI1200
    3MAI1210
    3MAI1220

```

```

      GO TO 84
110 CONTINUE
C   FORCE SV TO BE CONSISTENT WITH R AND G
    DO 111 I=IA,IB,I1
      SV(I)=DELTAR(I)/G(I)
      IF(IALPHA.EQ.1) GO TO 111
      TEMP(1)=R(I+I)+R(I)
      IF(IALPHA.EQ.3) TEMP(1)=TEMP(1)*R(I+1)+R(I)**2
      SV(I)=SV(I)*TEMP(1)

111 CONTINUE
      SOLID(28) = F(IND)
      BLANK4 = BLANK4 + SOLID(28)*DTH2
      209 SOLID(18)=SOLID(18)+1.0
      COUNT=SOLID(18)
      TH=TH+DTH2
      180 IF (COUNT - CNTMAX) 190, 301, 301
      190 IF(TH .GT. THMAX) GO TO 301
      IF (ABS(BLANK1) .LT. 0.1) GO TO 195
      IF (AMOD(SOLID(18), BLANK1) .LT. 0.5) CALL WTAPE
      195 IF (TU .GT. 0.5) GO TO 210
      CALL PRINT
      IF(BLANK1 .LT. 0.5) CALL WTAPE
      210 CALL SSWTCH(6,KUUUFX)
          GO TO(250,5),K000FX
      250 S1 = 1.0250
          GO TO 610
      301 S1 = 1.0301
          GO TO 610
      600 S1 = 1.0600
          GO TO 610
      605 S1 = 1.0605
      610 CALL UNCLE
      END

```

3MAI1230  
3MAI1240

3MAI1300  
3MAI1310  
3MAI1320

3MAI1390  
3MAI1400  
3MAI1410  
3MAI1420  
3MAI1430  
3MAI1440  
3MAI1450  
3MAI1460  
3MAI1470

3MAI1510

BIT FOR HYDRO/OUT, HYDRO/UUT, HYDRO/OUTI  
SUBROUTINE HYDRO  
C OUTPUT VERSION OF HYDRO  
C DEVELOPED BY JIM PALMER  
COMPILED SEPTEMBER 1, 1967 WBL  
C  
C USES RADIATION PRESSURE TENSOR

HYDRO0010

```

C*****S P U T T E R C O M M O N *****C
C*                                         HYDR0050
C*                                         HYDR0060
C*                                         *HYDR0070
C*                                         **HYDR0080
C*                                         **HYDR0090
C*                                         **HYDR0100
C*
COMMON LMDA(37), NR , NSMLR , IA , IB , ICA , ICB , 'HYDR0110
1 KMAX , BLANK1, BLANK2, BLANK3, IAP1 , IBP1 , ICAP1 , ICBP1 , HYDR0120
2 II , IG , NRAD , BLANK4, IAM1 , IBM1 , ICAM1 , ICBM1 , HYDR0130
3 IIP1 , IGM1 , IALPHA, BLANK5, TH , TMAX , HLANK6, DELPRT, HYDR0140
4 FREQ , CNTMAX, AR , ASMLR , PUSHA , PUSHB , HOILA , BOILB , HYDR0150
5 CVA , CVB , SLUG , ALPHA , HVA , HVB , HCA , HCB , HYDR0160
6 EMINA , EMINB , CA , CB , GA , GB , GL , GR , HYDR0170
7 RHUL , RHOR , EPIO , EPSI , RIA , RIR , RDIA , RDIB , HYDR0180
8 RPIA , RPIU , RPDIA , RPDIB , TPRINT, TA , DTRMIN, DTMAX , HYDR0200
COMMON TD , TE , UTH2 , DTH2P , UTH1 , SWITCH, CO , CMIN , DELTA , HYDR0210
1 DTMAX1, DTMAX2, DTMAX3, UTR , AC03T4, CNVRT , SUMRA , SUMRB , HYDR0220
2 GAMA , WCRIT , SIGMAJ, AC , ACO3T4, CNVRT , SUMRA , SUMRB , HYDR0230
3 ROI1A , ROIAMI1, ROIU , ROIBP1, GMS , S1 , S2 , S3 , HYDR0240
4 S4 , S5 , S6 , S7 , S8 , S9 , S10 , S11 , HYDR0250
5 S12 , S13 , S14 , S15 , S16 , S17 , S18 , S19 , HYDR0260
6 S20 , EO , FO , TAU , ZERO , R (152) , DELTAR(152) , HYDR0270
7 ASQ (152) , RU (152) , VD (152) , RDD (152) , SMLR (152) , HYDR0280
8 DELR ( 37) , P (152) , P1 (152) , PH (152) , PBI (152) , HYDR0290
COMMON P2 (152) , SV (152) , RHO (152) , THETA (152) , HYDR0300
1 W (152) , E (152) , EI (152) , EK (152) , A (152) , HYDR0310
2 V (152) , G (152) , D (152) , C (152) , X2 (152) , HYDR0320
3 X3 (152) , X4 (152) , XS (152) , X6 (152) , X7 (152) , HYDR0330
4 SMLA (152) , SMLB (152) , SMLC (152) , SMLD (152) , SMLE (152) , HYDR0340
5 EC (152) , ER (152) , SMLQ (152) , SMLH (152) , BIGA (152) , HYDR0350
6 HIGH (152) , CV (152) , BC (152) , BR (152) , CHIC (152) , HYDR0360
7 CHIR (152) , CAPAC (152) , CAPAR (152) , CRTC (152) , CRTR (152) , HYDR0370
8 CRTPC (152) , GOFK (152) , FEW (152) , CAR (152) , OKLM ( 37) , HYDR0380
COMMON TELM ( 37) , EKLM ( 37) , ELM ( 37) , FCLM ( 37) , HYDR0390
1 FRLM ( 37) , WLM ( 37) , QLM ( 37) , AMASNO( 37) , CHRNO ( 37) , HYDR0400
2 ZP1 ( 37) , ZP2 ( 37) , SOLID ( 37) , ECHCK ( 37) , RK (104) , HYDR0410
3 RL ( 37) , RHOK (104) , RDK (104) , THETAK(104) , TEMP ( 16) , HYDR0420
4 HEAU ( 12) , MAXL , MAXLM , HYDR0430
C*                                         **HYDR0440
C*                                         **HYDR0450
C*                                         HYDR0460
C*                                         HYDR0470
C*                                         HYDR0480
C*                                         HYDR0490

```

**IL = IA**

```

IR = IB
CALL UVCHK(KDHY)
IAM1 = IA-1
C S4 NEGATIVE GIVES DUMMY HYDRO
IF(S4)10,20,20
10 IBM1 = IB-1
DO 15 I=IA,IBM1
C(I) = R(I)
15 DELTAR(I) = R(I+1)-R(I)
C(IB) = R(IB)
GO TO 70
20 P(IB) = 0.
G(IB) = 0.
RHO(IB) = 0.
SV(IB) = 0.
IALPHA = ALPHA
DO 30 I=IL,IR
C(I) = R(I)
IF(IALPHA.EQ.3)GO TO 25
IF(I.EQ.IB)P(IU) = 0.
IF (I .NE. 1) GO TO 23
RD(I) = RD(I) - P(I) / G(I) + A(I) * DTH1
GO TO 30
23 RD(I) = RU(I) + (2. * (P(I-1) - P(I)) / (G(I) + G(I-1)) + A(I)) +
2 DTH1
GO TO 30
25 IF (I .NE. 1) GO TO 27
RU(I) = 0.
GO TO 30
27 CONTINUE
AA1=2.*A(I)*(P(I-1)-P(I)+(SMLR(I+1)-SMLR(I-1))/2.)/(G(I)+G(I-1))
IF(AMIN1(RHO(I),RHO(I+1)).LT. 1.E-10) GO TO 28
IF (AMIN1(SMLR(I) / RHO(I), SMLR(I+1) / RHO(I+1)) .LT. S4)
2 GO TO 28
AA2=(A(I-1)*SMLR(I-1)-A(I+1)*SMLR(I+1))/(G(I-1)+G(I))
AA3 = (RHO(I) - SMLR(I)) + (SV(I) + SV(I-1)) / (C(I) + C(I))
RD(I) = RD(I) + DTH1 * (AA1 + AA2 + AA3)
GO TO 30
28 CONTINUE
IF(I.EQ.IB)P(IU) = 0.
BB2=A(I)*(SMLR(I-1)-SMLR(I+1))/(G(I)+G(I-1))
BB3=(RHO(I)-3.*SMLR(I))*(SV(I)+SV(I-1))/(C(I)+C(I))
RD(I)=RD(I)+DTH1*(AA1+BB2+BB3)
CALL UVCHK(KP)
IF(KP.LT. 2) X=SQRT(-KP)
30 CONTINUE
R(IA) = R(IA)+RU(IA)+DTH2
DO 50 I=IA,IBM1
TEMP(I) = DELTAR(I)+(RD(I+1)-RD(I))+DTH2
R(I+1) = R(I)+TEMP(I)
IF(S1U)41,42,43
41 A(I) = 1.
VD(I) = RD(I+1)-RD(I)
GO TO 50
42 S1 = 7.1

```

```
CALL UNCLE
43 A(I) = 3.*R(I)*R(I)
    CALL UVCHK(KP)
    IF(KP.LT. 2) X=SQRT(-KP)
    IF((ABS(R(I+1)-C(I+1))+ABS(R(I)-C(I)))/(R(I+1)+R(I)).LT.1.E-6)GO T
10 45
    VD(I) = (TEMP(I)*(R(I+I)+R(I+I)+R(I))+R(I)+R(I))-DELTAR(I)*(C(I+I
    I)*(C(I+1)+C(I))+C(I)*C(I)))/DTH2
    GO TO 50
45 VD(I) = (R(I+1)*R(I+1)+R(I)*(R(I+1)+R(I)))**(RD(I+1)-RD(I))
50 DELTAR(I) = TEMP(I)
    A(1B) = ALPHA*R(1B)**(ALPHA-1.)
    CALL UVCHK(KP)
    IF(KP .GT. 1) GO TO 70
    SI=7.129
    X=-I.
    X=SQR(X)
    CALL UNCLE
70 CONTINUE
    RETURN
    ENO
```

## EQUIVALENCE (LMDA(24), IBB)

DIMENSION ERLM(1)  
 EQUIVALENCE (FCLM, ERLM)

```

C
C
1 IF (PUSHA) 1, 5, 5
1 IC = 16
  FRLM(16) = FRLM(16) + UT*X2(IA)
  IF (EMINA - EI(ICUH1)) 2, 2, 10
2 IF (ICU - IC) 3, 3, 10
3 ICHM1 = ICU
  ICU = ICUP1
  ICUP1 = ICU + 1
  CALL EOS (ICCHM1)
  GO TO 10
5 IC = IA
  FRLM(18) = FRLM(18) + UT*X2(IU)
  IF (IFIX(PUSHA) .EQ. 0) GO TO 10
  IF (S3) 10, 6, 10
6 IF (EMINA .GT. EI(ICA)) GO TO 10
  IF (ICA .EQ. 1) GO TO 10
  ICAP1 = ICA
  ICA = ICAM1
  ICAM1 = ICA - 1
  CALL EOS (ICA)
10 CONTINUE
  DO 160 J = 1, MAXL
    ELM(J) = 0.
    IL = LMDA(J)
    NVAP = 0
    IF (IL .GE. IA .AND. IL .LE. IB) NVAP = 1
    IF (J .EQ. MAXL) GO TO 115
    IR = LMDA(J+1) - 1
    DO 110 I = IL, IR
      ELM(I) = ELM(I) + G(I)*E(I)
      IF (IFIX(S2) .EQ. 0) GO TO 105
      TEMP(1) = DELTAK(1)
      IF (S10) 104, 105, 102
102 TEMP(1) = TEMP(1) + (3. * R(I) + (K(I) + TEMP(1)) * TEMP(1)**2)
      GO TO 104
103 TEMP(1) = TEMP(1) + (2.*R(I) + TEMP(1))
104 ERLM(J) = ERLM(J) + RHO(I) * TEMP(1)
      IF (IFIX(S2) .NE. 3) GO TO 105
      ELM(J) = ELM(J) + RHO(I) * TEMP(1)
105 IF (IFIX(S6) .EQ. 0) GO TO 110
      QLM(J) = QLM(J) + SMLQ(I)*DTR
110 CONTINUE
115 IM = IL
  IMM1 = IM - 1
  IF (NVAP .NE. 1) GO TO 170
  FRLM(J) = FRLM(J) + DTR*X2(IM)
  IF (IM .EQ. 1) GO TO 180
  GX = G(IM)
  IF (IFIX(BOILA) .EQ. 0 .OR. IM .NE. IC) GO TO 160
  GX = SOLID(24)*G(IBB)
152 IF (PUSHA) 160, 180, 155

```

ECAL0490  
 ECAL0500

```
155 WLM(J) = WLM(J) + (P(IA-1)*G(IA) + P(IA)*GX)*RD(IA)*DTR/GX +
1 G(IA)) * A(IA)
GO TO 180
160 WLM(J) = WLM(J) + (P(IMM1)*GX + P(IM)*G(IMM1))*RD(IM)*DTR/GX +
1 G(IMM1))*A(IM)
GO TO 180
170 IF (SS) 180, 175, 160
175 IF (IM .NE. IA .AND. IM .NE. IB) FCLM(J) = FCLM(J) + DTR*(X7(IM) -
1 DELTA*(W(IM)*X5(IM) - K(IMM1)*X6(IM)))
180 CONTINUE
IF (IFIX(SUMRA) .EQ. 0) GO TO 200
DO 190 I = IA, IMM1
190 RDD(I) = RDU(I) + EK(I)*(SV(I)*P(I) + E(I) + (RD(I)*RD(I) +
1 RU(I+1)*RD(I+1))/4.)*DTR
200 FRLM(19) = FRLM(19) + DTR*X2(NR)
P(152) = AMAX1(P(152), SOLID(28))
RETURN
END
```

ECAL1110  
ECAL1120

U,T FOR EOS/UT, EOS/OUT, EOS/JUT1  
SUBROUTINE EOS(UJ1)

C OUTPUT VERSION OF EOS -- CHANGES TO USE THE RADIATION PRESSURE  
CALCULATED IN THE SCATTERING TRANSPORT CODE

COMPILED SEPTEMBER 1, 1967 BY

C COMPILED FEB. 15, 1967 H. SCHLAUG

C INCLUDES SPECIAL SECTION FOR X-RAY SOURCE ROUTINE  
THAT HANDLES DEPOSITION IN MULTI ELEMENT MATERIALS

C COMPILED ON MAY 31, 1966 ADDING TEST ON NONEO OPTION BEFORE  
TAKING PATH I=3

C COMPILED MAY 18, 1966 BY G. A. LANE

C COMPILED MAY 10, 1965. RSE TAKES ADVANTAGE OF NEW PATH (-3.) IN  
EIONX.

EOS 0030

C COMPILED ON MARCH 31, 1966 ADDING COMMON FOR NRAR (GE)

EOS 0040

C.....

C.....

C.....

SPUTTER COMMON

COMMON	LMDA(37)	, NLR	, NSHLR	, IA	, IB	, ICA	, ICB	,	
1	KMAX	, BLANK1	, BLANK2	, BLANK3	, IAP1	, IUP1	, ICAP1	, ICBP1	
2	II	, IG	, NRAD	, BLANK4	, IAM1	, IBM1	, ICAM1	, ICBM1	
3	IIP1	, IGH1	, IALPHA	, BLANK5	, TH	, TMAX	, BLANK6	, DELPRT	
4	FREQ	, CNTMAX	, AR	, ASMLR	, PUSHU	, PUSHU	, ROILA	, BOILB	
5	CVA	, CVB	, SLUG	, ALPHA	, NVA	, HVB	, HCA	, HCB	
6	EMINA	, EMINB	, CA	, CB	, GA	, GB	, GL	, GR	
7	RHOL	, RHOK	, EPIU	, EPSI	, RIA	, RIB	, RODA	, RDIB	
8	RHIA	, RPIB	, RPUIA	, RPOIB	, TPRINT	, TA	, TB	, TC	
	COMMON	TU	, TE	, UTH2	, UTH2P	, UTH1	, DTRMIN	, DTMAX	
1	DTMAX1	, DTMAX2	, DTMAX3	, DTR	, SWITCH	, CO	, CMIN	, DELTA	
2	GAMA	, WCRIT	, SIGMAU	, AC	, AC03T4	, CNVRT	, SUNRA	, SUMRU	
3	ROIA	, RUIAM1	, RUIB	, ROIBP1	, GMS	, S1	, S2	, S3	
4	S4	, S5	, S6	, S7	, S8	, S9	, S10	, S11	
5	S12	, S13	, S14	, S15	, S16	, S17	, S18	, S19	
6	S2U	, EO	, FO	, TAU	, ZERO	, H	(152),	DELTAR(152),	
7	ASU	(152),	RU	(152),	VD	(152),	ROD	(152),	SMLR (152),
8	DELK	(37),	P	(152),	P1	(152),	PB	(152),	PB1 (152)
	COMMON	P2	(152),	SV	(152),	RHO	(152),	THETA	(152),
1	D	(152),	E	(152),	EI	(152),	EK	(152),	A (152),
2	V	(152),	G	(152),	D	(152),	C	(152),	X2 (152),
3	X3	(152),	X4	(152),	X5	(152),	X6	(152),	X7 (152),
4	SMLA	(152),	SMLB	(152),	SMLC	(152),	SMLU	(152),	SMLE (152),
5	EC	(152),	ER	(152),	SMLQ	(152),	SMLH	(152),	BIGA (152),
6	B1GU	(152),	CV	(152),	DC	(152),	BR	(152),	CHIC (152),
7	CHIK	(152),	CAPAC	(152),	CAPAK	(152),	CRTC	(152),	CRTR (152),
8	CRTPC	(152),	GOFR	(152),	FEW	(152),	CAH	(152),	OKLM (37)
	COMMON	TELM	(37),	EKLM	(37),	ELM	(37),	FCLM	(37),
1	FRLM	(37),	WLM	(37),	QLM	(37),	AMASNO	(37),	CHRNO (37),
2	ZP1	(37),	ZP2	(37),	SOLID	(37),	ECHK	(37),	RK (104),
3	RL	(37),	RHOK	(104),	ROK	(104),	THETAK	(104),	TEMP (16),
4	HEAD	(12),	MAXL	,	MAXLM	,			

C



```

C MATERIAL NUMBER OF ZERO LEADS TO AN ERROR EXIT
C S1 = 4.0004
C CALL UNCLE
C OKLM(J) NOT PROPERLY ENTERED IN INPUT
C
57 L= .5*OKLM(J)
RHOR= 0.
IF ((I-IA)*NA.GT.0 .AND. ZP1(26).NE.0.) EROR = -3.0
58 CONTINUE
IF (L.LE.200 OR.L.GE.301.AND.L.LE.400) GO TO 7
IF (L.EU. 200 .AND. SMIA(I) .GT. RHOR) GO TO 10
C TO AVOID 'EIONX', SET RHOR .EQ. 1.E38.
C
GO TO 8
C
10 CONTINUE
C THIS PATH IF OKLM(J) .EQ. 208, FORCES 'EIONX' IF AND ONLY IF
C TEMPERATURE OF ZONE AT START OF CYCLE .GT. RHOR.
L = 102
GO TO 7
C
23 CONTINUE
C THIS PATH IF OKLM(J) .EU. 102, FORCES 'ESB' IF AND ONLY IF
C TEMPERATURE OF ZONE AT START OF CYCLE .LE. RHOR.
L = 208
GO TO 8
C
7 CONTINUE
IF (L.EQ. 102 .AND. SMIA(I) .LE. RHOR) GO TO 23
C TO AVOID 'ESB' SET RHOR .EQ. -1.E38.
IF(L.EQ.101.OR.L.EQ.102.OR.L.EQ.6.OR.L.EQ.306) EION(14)=RIA
IF(ERUR.EQ.(-3.0)) ZBAR = FLW(I)
IF (OKLM(J) .GT. 100.) GO TO 60
CALL EIONMS(THETA(I),SV(I),L,EROR)
GO TO 61
60 CALL EIONX (THETA(I),SV(I),L,EROR)
61 CONTINUE
IF (L.GE.201 .AND. L .LE. 300 ) GO TO 9
NBAR= EION(17)
IF (S11) 54.52*51
52 IF (EROR) 53.54*55
55 S11= EHUR + 100. + FLOAT(I)
GO TO 57
51 IF (ERUR) 53.54*54
53 S11= EHUR-FLOAT(I)+100.
C
54 CONTINUE
C
IF (EION(14) .EQ. 0. ) GO TO 901
S1= EION(14)
CALL UNCLE
C EION(14) IS SET EQUAL TO ZERO BY ANY VALID EXIT
C
901 CONTINUE
C

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P1(I) = EION(7)
E(I) = EION(8)
IF ((I-IX)*NX.GT.0 .AND. ZP1(26).NE.0.) GO TO 203
CV(I) = EION(9)
ASQ(I) = EION(11)
FEW(I) = EION(3)
PB1(I) = EION(10)
IF (R1B.E0.0.) GO TO 199
IF (SV(I).GE.SOLIU(17)) GO TO 199
P1(I) = P1(I) + 1.E12*(SOLIU(17)/SV(I)-1.)**2
ASU(I) = SQRT(ASQ(I)**2+2.E12*SOLIU(17)*(SOLID(17)/SV(I)-1.))
GO TO 194
8 CONTINUE
C
9 CONTINUE
L= L-200
GO TO (11,12,13,14,15,16,17,18,19,20,21,22),L
11 CALL ES1      EOS 1240
12 CALL ES2      EOS 1250
13 CALL ES3      EOS 1260
14 CALL ES4      EOS 1270
15 CALL ES5      EOS 1280
16 CALL ES6      EOS 1290
17 CALL ES7      EOS 1300
18 CALL ES8      EOS 1310
19 CALL ES9      EOS 1320
20 CALL ES10     EOS 1330
21 CALL ES11     EOS 1340
22 CALL ES12     EOS 1350
6 CONTINUE
100 CONTINUE
C
ASQ IS SPEED OF SOUND
200 ASQ(I) = SQRT(.25*GAMA*P1(I)*SV(I))
199 IF (IFIX(S2).EQ.3) GO TO 404
GO TO 150
404 P1(I) = RHO(I) + 0.3333333 + P1(I)
GO TO 400
150 IF(RPIA)400,201,400
201 TEMP(1)=THETA(I)**4      EOS 1520
P1(I)=P1(I)+.5*(SMLR(I)+SMLR(I+1))
PB1(I)=137.0*TEMP(1)+PB1(I)
E(I)=E(I)+137.0*TEMP(1)*SV(I)
CV(I)=CV(I)+548.0*TEMP(1)/THETA(I)*SV(I)
EOS 1540
EOS 1550
EOS 1560

400 IF (S4.NE.0.) GO TO 401
IF (ECHCK(J+17).EQ.0.) ECHCK(J+17) = HVA + HCA - 1.5 * AMASNO(J)
1           +(HCA/CVA+.025)*(1./NBAR + FEW(I))
401 E(I) = E(I) + ECHCK(J+17)
500 RETURN
ENO
EOS 1580

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B, FOR RAD/BL, RAD/BL, RAD/BLI  
SUBROUTINE RAD

COMPILED OCTOBER 9, 1967 ABL

COMPTON SCATTERING WITH OPTION FOR THOMSON SCATTERING

C UNIFIED RADIATION CODE -- PLANE OR SPHERICAL TRANSPORT

C DIFFUSION IN PLANE, CYLINDRICAL, OR SPHERICAL GEOMETRY

CAVEAT. MIXTURE OF TRANSPORT AND DIFFUSION IN CYL. GEOMETRY PROHIBITED.

C\*\*\*\*\*PRAD 60  
C S P U T T E R C O M M O N \*PRAD 70  
C \*PRAD 80

COMMON	LMUA(37)	, NR	, NSMLR	, IA	, IB	, ICA	, ICB	, PRAD	90	
1	KMAX	, BLANK1	, BLANK2	, BLANK3	, IAPI	, IBP1	, ICAP1	, ICAP1	, PRAD 100	
2	II	, IG	, NRAD	, BLANK4	, IAM1	, IBM1	, ICAMI	, ICBM1	, PRAD 110	
3	IIP1	, IGM1	, IALPHA	, BLANK5	, III	, TMAX	, BLANK6	, DELPRT	, PRAD 120	
4	FREQ	, CNTMAX	, AK	, ASMLR	, PUSHA	, PUSHB	, BOILA	, BOILB	, PRAD 130	
5	CVA	, CVB	, SLUG	, ALPHA	, HVA	, HVH	, HCA	, HCB	, PRAD 140	
6	EMINA	, EMINU	, CA	, CU	, GA	, GB	, GL	, GR	, PRAD 150	
7	RHOL	, RHOR	, EPIO	, EPSI	, KIA	, KIB	, ROIA	, ROIB	, PRAD 160	
8	RPIA	, RPIB	, RPOIA	, RPOIB	, TPRINT	, TA	, TB	, TC	, PRAD 170	
	COMMON	TU	, TE	, DTH2	, DTH2P	, DTH1	, DTRMIN	, DTMAX	, PRAD 180	
1	DTMAX1	, DTMAX2	, DTMAX3	, UTR	, SWITCH	, CO	, CMIN	, DELTA	, PRAD 190	
2	GAMA	, WCRIT	, SIGMAJ	, AC	, AC03T4	, CNVRT	, SUMRA	, SUMRB	, PRAD 200	
3	ROIA	, ROIAM1	, ROIU	, ROIUP1	, GMS	, S1	, S2	, S3	, PRAD 210	
4	S4	, S5	, S6	, S7	, S8	, S9	, S10	, S11	, PRAD 220	
5	S12	, S13	, S14	, S15	, S16	, S17	, S18	, S19	, PRAD 230	
6	S20	, EO	, FO	, TAU	, ZERO	, R	(152)	, DELTAR(152)	, PRAD 240	
7	ASQ	(152)	, RD	(152)	, VD	(152)	, RUD	(152)	, SMLH (152)	, PRAD 250
8	DELR	(37)	, P	(152)	, P1	(152)	, PB	(152)	, PB1 (152)	, PRAD 260
	COMMON	P2	(152)	, SV	(152)	, RHO	(152)	, THETA	(152)	, PRAD 270
1	W	(152)	, E	(152)	, EI	(152)	, EK	(152)	, A (152)	, PRAD 280
2	V	(152)	, G	(152)	, D	(152)	, C	(152)	, X2 (152)	, PRAD 290
3	X3	(152)	, X4	(152)	, X5	(152)	, X6	(152)	, X7 (152)	, PRAD 300
4	SMLA	(152)	, SMLB	(152)	, SMLC	(152)	, SMLD	(152)	, SMLE (152)	, PRAD 310
5	EC	(152)	, ER	(152)	, SMLQ	(152)	, SMLH	(152)	, RIGA (152)	, PRAD 320
6	BIGU	(152)	, CV	(152)	, BC	(152)	, BR	(152)	, CHIC (152)	, PRAD 330
7	CHIK	(152)	, CAPAC	(152)	, CAPAR	(152)	, CRTG	(152)	, CRTR (152)	, PRAD 340
8	CRTPC	(152)	, GOFR	(152)	, FEW	(152)	, CAR	(152)	, OKLM (37)	, PRAD 350
	COMMON	TELM	(37)	, EKLM	(37)	, ELM	(37)	, FCLM	(37)	, PRAD 360
1	FRLM	(37)	, WLM	(37)	, QLM	(37)	, AMASNO(37)	, CHRNO (37)	, PRAD 370	
2	ZP1	(37)	, ZP2	(37)	, SOLID	(37)	, ECHCK (37)	, RK (104)	, PRAD 380	
3	RL	(37)	, RHOK	(104)	, ROK	(104)	, THETAK(104)	, TEMP (16)	, PRAD 390	
4	HEAD	(12)	, MAXL	,	MAXLM				, PRAD 400	

\*PRAD 410

C	COMMON /LINULY/	MNU,SGIL,IHNU,NHNU,HNUP,NT,IM,IN,DHNU,THICK,NY	PRAD 460
C	COMMON /CNTRL/	SCYCLE, JMULT	PRAD 470
C	COMMON /DAVIS/	X(4000), ICX, ICY	
C	COMMON /PALMER/	FI0(152), FI1(152), FI2(152), FI3(152), F00(152),	
2	FQ1(152), FQ2(152), FQ3(152), JDRLM		
C	COMMON /JIM/	NN, FMU, R1, R2, HO, EST, IL, I2, GMP, A1, A3, FMUS,	
2	FS, LUF, LRI, IZN, TG1, TG2, F2		

PRAD 500

DIMENSION	C500 (1), PR (1), FM (1), H (1), H2 (1),
1	H3 (1), H4 (1), FMS (1), Q1 (1), TG (1), Q3 (1),

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C  Q37  (I), Q38  (I), SUMX2 (I), SUMX3 (I), SUMX4 (I), X8  (I), P  550
C  Y  (I), T2  (I), UX  (I), FL  (I), TR  (I), FSM  (I), P
C  4 FSP  (I)  P  570
C
C      EQUIVALENCE          (AC03T4,THDBG), (BC, +SUMX4), (BIGA, Y ), P  580
C      1  (BIGB, H ), (LI, H3 ), (CAR, Q37 ), (CHIC, +SUMX3), P  590
C      2  (CHIR, Q38 ), (CTR, +SUMX2), (X7, PR ), (GOFK, Q3 ), P  600
C      3  (PB, Q1 ), (S12, EDITMF), (EC, H2 ), ( W, QX ), P
C      4  (SMLA, FMS ), (SALB, FL ), (SMLC, TR ), (SMLH, H4 ), P
C      5  (ER, FM ), (V, TG ), (SMLD, FSM ), (CRTA, CS30 ), P
C      6  (AS, T2 ), (X4, X8 ), (SMLE, FSP )  P  660
C
C ***** 660  P
C
C      OX CONTAINS X FROM THE PREVIOUS Y LINE  P  670
C
C      CSOD  SAME AS  CRTA  P  680
C      EDITMF SAME AS  S12  P  690
C      LI  SAME AS  PB  P
C      FM  SAME AS  ER  P
C      H  SAME AS  BIGB  P  700
C      H2  SAME AS  EC  P  710
C      H3  SAME AS  BR  P
C      H4  SAME AS  SMLH  P  720
C      PR  SAME AS  X7  P
C      FMS  SAME AS  SMLA  P
C      FL  SAME AS  SMLB  P
C      TR  SAME AS  SMLC  P
C      FSM  SAME AS  SMLD  P
C      FSP  SAME AS  SMLE  P
C      Y2  SAME AS  X5  P
C      OX  SAME AS  W  P
C      TG  SAME AS  V  P
C      Q3  SAME AS  GOFK  P  830
C      Q37  SAME AS  CAR  P  840
C      Q38  SAME AS  CHIR  P  850
C      SUMX2  SAME AS  CTR  P  860
C      SUMX3  SAME AS  CHIC  P  870
C      SUMX4  SAME AS  BC  P  880
C      THDBG  SAME AS  AC03T4  P  890
C      Y  SAME AS  BIGA  P  910
C      X8  SAME AS  X4  P  940
C
C ***** 940  PHAD  800
C
C      SOLID(37) IS THE SCATTERING COEFFICIENT  P
C      SOLID(36) IS THE COMPTON SWITCH. ZERO, COMPTON; NONZERO, THOMSON. P
C      F10, F11, F12, F13, FREQUENCY-DEPENDENT SCATTERING SOURCES. P
C      READ IN FROM DRUM, ALTERNATELY FROM LOGICAL UNITS 25 AND 26. P
C      RECIPROCAL ELECTRON REST ENERGY IN EV--1 IS 1.95692E-6  P
C      IDMX = 4000  P
C      T4 = 1.  P
C      TAX = ABS(SOLID(37)) * 1.95692E-6  P
C      IF (ABS(SOLID(36)) .GT. 1.E-20) TAX = 0.  P
C      JORUMI = 51 - JORUM  P
C      REWIND JORUM  P

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      IF (ABS(ER(I)) .LT. 1.E-20 .OR. WSHU .LT. .001 * WSH) GO TO 105      P
C      ACCURACY CRITERION
      TEMP(1)=SLUG*WSHU/ABS(ER(I))
C      STABILITY CRITERION
105  TEMP(2) = TELM(25) + (AHINI(TEMP(1), (.5 + 1.5 * H(I)**2) * CV(I)/P
     1 (4.1132E12*TE-M(3)*THETA(I)**3)))
     IF (TEMP(2) .LT. 1.E-20) GO TO 107
     IF (TEMP(2).GT.UTR1) GO TO 106
     UTR2=UTR1
     IMN2=IMN1
     UTR1=TEMP(2)
     IMN1=I
     GO TO 107
106  IF (TLMP(2).GT.UTR2) GO TO 107
     UTR2=TEMP(2)
     IMN2=I
107  CONTINUE
     UTRMIN=UTR1
     EO=IMN1
     IF (UTR1.GT.TELM(26)) GO TO 108
     TELM(26)=UTR1
     TELM(27)=IMN1
     TELM(28)=UTR2
     TELM(29)=IMN2
     TELM(30)=SGL10(10)+1.
108  IF (UTRMIN-UTR) 111,112,109
109  UBLANK3=TH+AMINI(UTRMIN,GR*DTH2)
     GO TO 112
111  NRAD=2P1(18)/UTRMIN + 1.
     UTR=2P1(18)/FLOAT(NRAD)
     IF (NRAD .LE. 50) GO TO 112
     S1=13.0112
     CALL UNCLE
112  CALL UVCHK(KX)
     GO TO (150,160), KX
150  S1 = 13.0150
     CALL UNCLE
160  IF (GL .GT. 0.) IR = IM
     GO TO (128, 113, 113), IALPHA
C      CHARACTERISTICS IN NONPLANE GEOMETRY -- SET UP X AND Y OUTSIDE FREQ LOOP P
C
113  K = 0
C      DETERMINE WHETHER TO SKIP ZONES WITH Y-LINES
114  K1=K
     Y(1)=0.0
     X4(1) = 0.0
     JK=1
*****SET UP Y LINES*****
C      SET UP Y LINES
C
1530
*****DRAW Y LINES---ONE FOR EACH LARGE TEMPERATURE CHANGE*****
C      DRAW Y LINES--ONE FOR EACH LARGE TEMPERATURE CHANGE
C      FORCE Y LINE ON SURFACE OF SOLID IF ANY
1540
1550
1560
1570
1580
1600
1610

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C      NO Y GREATER THAN C(IR+1)          P 1630
C
C      IF (IN .LE. 1) GO TO 115          P
C      JK = 2                            P
C      Y(2) = C(IN)                      P
C      X4(2) = CSQD(IN)                  P
C 115 DO 116 I = IN, IR                P
C      IF (ABS(G37(I+1) - G37(I)) = 0.5) 117, 117, 116    P
C 116 JK=JK+1                          P 1670
C      Y(JK)=C(I+1)                      P
C      X4(JK) = CSQD(I+1)                P 1690
C      K1=K                            P
C      GO TO 118                        P
C 117 IF (K1 .LE. 0 .OR. I .EQ. IR) GO TO 116    P
C      K1=K1-1                          P
C 118 CONTINUE                         P 1740
C      NY=JK                          P
C      GO TO 120                        P
C 119 K=K+1                            P
C      IF (K .LE. 10) GO TO 114          P
C      S1 = 13.0119                      P
C      CALL UNCLE                       P
C***** F I N D   C O M P L E T E   S E T   O F   X   V A L U E S *P 1780
C      (INTERSECTIONS OF RAVII WITH Y LINES)          *P 1790
C***** *P 1800
C***** *P 1810
C***** *P 1820
C***** *P 1830
C***** *P 1840
C      F O R M A T = -X4, -N U M B E R   O F   I N T E R S E C T I O N S , X ' S . *P 1860
C
C 120 K2=1                            P
C      DO 125 J=2,JK                    P 1890
C      I=IMP1                          P
C      X(K2) = -X4(J)                  P 1910
C      K2=K2+2                          P 1920
C      KK=1                            P
C 121 TS1 = CSQD(I) - X4(J)          P
C      IF (TS1) 124, 124, 122          P
C 122 X(K2)=SQRT(TS1)                P 1960
C      K2=K2+1                          P
C      IF (K2-IDMX) 123, 123, 119    P
C 123 I=I-1                            P 1990
C      KK=KK+1                          P
C      GO TO 121                        P
C 124 KKK=K2-KK                      P 2020
C      X(KKK)=-1(KK-1)                P
C 125 CONTINUE                         P
C      FINISH X-BLOCK WITH A NEGATIVE NUMBER          P 2070
C      X(K2)=-1.0                      *PRAD1420
C***** *PRAD1430
C***** *PRAD1440
C***** *PRAD1460
C***** PRAD1480
C
C      B E G I N   F R E Q U _ N C Y   L O O P       *PRAD1460
C***** *PRAD1480

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C           SET UP MAX FREQ BOUNDARY          PRAD1490
C           COMPTON SCATTERING FORBIDS VERY HIGH FREQUENCIES.    PRAD1500
C           THETAK(103) CHOSEN TO BE COMPATIBLE WITH KAP6/JP    P
C           128 HNU = THETAK(103)                                P
C           IHNUP=1                                              PRAD1520
C           DO 129 I=IN,IMP1
C           X4(I) = 0.
C           FQ0(I) = 0.
C           FQ1(I) = 0.
C           FQ2(I) = 0.
C           FQ3(I) = 0.
C           RDD(I) = EK(I)
C           EK(I) = 0.
C           SMLR(I) = 0.
C           129 SUMX2(I)=0.0
C           IF (KMAX .NE. 0) GO TO 310
C
C           MONOFREQUENCY CALCULATION          PRAD2070
C
C           210 NHNU=1
C           DO 220 I=IN,IM
C           X0(I)=Q1(I)
C           240 DFB = 1.0
C           HNU = .001
C           DHNU = THETAK(103)
C           ICX = IM
C           ICY=IN
C           GO TO 460
C
C           TYPICAL GROUP CALCULATION OF SOURCES      PRAD2220
C
C           310 CALL KAPPAIN,IM
C           ZZ = 0.
C           DHNUP = DHNU
C           UHNU=HNUP-HNU
C           350 ICX = IM
C           ICY = IN
C           IF (GL .LT. 1.E-20) GO TO 370
C           DO 360 I=IN,IM
C           DFB = PLNKUT(HNU / THETA(I), HNUP / THETA(I))
C           360 X6(I)=DFB*Q1(I)
C           GO TO 460
C           370 DO 450 I=IN,IR
C           BETA=HNU/THETA(I)
C
C           AVOID CALCULATION OF DFB LESS THAN 1E-5      PRAD2490
C
C           IF (BETA .GT. 19.) GO TO 430
C           BETAP = HNUP / THETA(I)
C           IF (BETAP .GT. 0.01) GO TO 440
C           430 X6(I)=0.0
C           GO TO 450
C
C           FORM SOURCE X6

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C
440 DFB=PLNKUT(BETA,BETAP)          PRAD2910
X6(I)=DFB*Q1(I)                   PRAD2920
450 CONTINUE                         PRAD2930
CALL UVCHK(KX)
GO TO (452, 460), KX               PRAD2980
452 S1 = 13.0452
CALL UNCLE
460 IF (INM1) 470, 490, 480

C
C           SET BLACKBODY CONDITION FOR IA GREATER THAN 1      PRAD3000
C
C           470 S1 = 13.0470                                     PRAD3010
C           CALL UNCLE                                      PRAD3020
480 DFB = PLNKUT (HNU/THETA(INM1),HNUP/THETA(INM1))    PRAD3040
X6(INM1) = DFB * THETA(INM1)**4                         PRAD3050
C           SET BLACKBODY CONDITION IF DESIRED FOR IMP1       PRAD3060
490 IF (ABS(GL - 0.5) .GT. 1.E-5) GO TO 510             PRAD3070
IF (ABS(THETA(IMP1)) .LT. 1.E-20) GO TO 500
FMS(INM1) = 0.
DFB = PLNKUT(HNU / THETA(IMP1), HNUP / THETA(IMP1))   PRAD3100
X6(IMP1) = DFB * THETA(IMP1)**4                         PRAD3110
FMS(IMP1) = 0.
GO TO 510
500 X6(IMP1)=0.                                         PRAD3120
510 IF (ABS(SOLID(37)) .LT. 1.E-20) GO TO 515          PRAD3130
C           SCATTERING. SET RADIATION REGION ACTIVE THROUGHOUT, TAKING IN DATA
C           FROM DRUM, AND SET UP COMPTON SCATTERING FREQUENCY PARAMETERS
ICY = IN
ICX = IM
READ (JDRUM) FI0, FI1, FI2, FI3
HNUX = AMIN1(HNUP, 1.E5)
GAMMA = AMIN1(0.2, 0.9764E-6 * (HNU + HNUX))
A1P = HNU**2 / DHNU * 1.95692E-6
IF (IHNU .EQ. 1) GO TO 512
A3 = HNUX**2 / DHNUP * 1.95692E-6
512 A1 = A1P + 3. * GAMMA
GMP = 1. - 2. * GAMMA
515 Q31=0.0

C           FORM ROSSELAND AND PLANCK OPTICAL DEPTHS          PRAD3150
C
DO 570 I=IN,IM
IF (CAPAR(I)) 530,530,520
520 IF (CAPAC(I)) 530,530,540
530 S1=13.0530
CALL UNCLE
540 TAUX = AMAX1(SOLID(37), 0.)
C           SPECIAL CODING TO FORCE KAPPAS TO BE AT LEAST 0.2     PRAD3160
CPC = AMAX1(CAPAR(I) + TAUX, 0.2)
C           HHTAX MUST BE LIMITED TO 2 * GAMMA * KAPPA(S) = .08   PRAD3170
HHTAX = AMIN1(TAX * (HNU + HNUX), 0.08)
GO TO (542, 544, 544), IALPHA
542 QQ1 = G(I)
GO TO 546

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544 QQ1 = (C(I+1) - C(I)) / SV(I)
546 CPA = CPC - MHTAX
      H3(I) = CPA + QQ1
C CHOOSE ALL ROSELAND IF SOLID(10) IS NONZERO
      IF (ABS(SOLID(10)) .LT. 1.E-20) GO TO 550
      H2(I) = H3(I)
      H(I) = CPA / SV(I)
      GO TO 560
PRAD3280
550 CPB = AMAX1(CAPAC(I) + TRUX, 0.2)
      H2(I) = CPB + QQ1
      H(I) = CPB / SV(I)
PRAD3290
560 Q31=Q31+H3(I)
      Q3(I+1)=Q31
      H(I) = 0.5 * H(I)
      H2(I)=0.5*H2(I)
      H3(I)=0.5*H3(I)
      X2(I)=0.0
      X3(I)=0.0
      X4(I)=0.0
      Y2(I) = 0.
      TG(I) = 0.
      RHO(I) = 0.
      PR(I) = 0.
      FL(I) = 0.
PRAD3310
PRAD3320
570 TR(I) = 0.
      X2(IMP1)=0.0
      X3(IMP1)=0.0
      X4(IMP1)=0.0
      TG(IMP1) = 0.
      RHO(IMP1) = 0.
      PR(IMP1) = 0.
      FL(IMP1) = 0.
      TR(IMP1) = 0.
      Y2(IN)=X6(IN)
      TG(IN)=0.0
PRAD3420
PRAD3430
PRAD3440
C
C FORM Y2 AND TG SET X3=-1 IF A DIFFUSION CRITERION MET USING HCB
C
600 ICXM1=ICX-1
      IF (ICY.GT.ICXM1) GO TO 650
      DO 640 I=ICY,ICXM1
      TEMP(1)=H3(I+1)+H3(I)
      IF (AMAX1(X6(I),X6(I+1)).LT.1.E-30) GO TO 610
      IF (ABS((H3(I) - H3(I+1)) / TEMP(1)) .GT. 0.333) GO TO 610
      IF (ABS((X6(I)-X6(I+1)) / (X6(I)+X6(I+1))) .LT. 0.333) GO TO 620
PRAD3520
PRAD3570
PRAD3580
PRAD3590
PRAD3600
PRAD3610
PRAD3620
PRAD3630
610 TG(I+1)=0.0
      GO TO 640
PRAD3680
PRAD3690
620 TG(I+1)=(X6(I+1)-X6(I))/TEMP(1)
      Y2(I+1)=(X6(I+1)*H3(I)+X6(I)*H3(I+1))/TEMP(1)
PRAD3700
PRAD3740
C
C FORCE TRANSPORT FOR RAPIDLY VARYING SOURCE OR POSITIVE HCB
      IF (ABS(TG(I+1)) .GT. 0.1 * Y2(I+1) .OR. HCB .GT. 0.) GO TO 640
      630 X3(I+1)=-1.0
      640 CONTINUE
PRAD3760
PRAD3770
C
C FORCE DIFFUSION FOR NEGATIVE HCB
650 IF (HCB) 651, 655, 655

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```

651 IF (GL .GT. 0.9) GO TO 654
INP1 = IN + 1
X3(1) = -1.
DO 653 I = INP1, IM
Q = H3(I-1) + H3(I)
IF (Q .GT. 0.) GO TO 652
S1 = 13.0652
CALL UNCLE
652 TG(I) = (X6(I) - X6(I-1)) / Q
653 X3(I) = -1.
GO TO 655
654 S1 = 13.0654
CALL UNCLE
C                                     PRAD3780
C LAST ZONE MUST BE TRANSPORT IF EXTERNAL INPUT INTENSITIES PROVIDED
655 IF (GL .GT. 0.9) X3(IM) = 0.0
IF (ABS(GL - 0.5) .LT. 1.E-5) GO TO 690
Y2(IMP1)=X6(ICX)
TG(IMP1)=0.0
PRAD3910
PRAD3920
PRAD3930
C EXTEND TRANSPORT REGION BOUNDARIES TO PROVIDE 5 MEAN FREE PATHS
C                                     PRAD3960
C                                     PRAD3970
690 I=IN+1
CALL DVCHK(KX)
GO TO (692, 700), KX
692 S1 = 13.0692
CALL UNCLE
700 IF (X3(I)) 710,730,720
710 I=I+1
IF (I-ICX-1) 700,730,820
720 S1=13.0720
CALL UNCLE
730 IF (I .EQ. IMP1) GO TO 820
J = I - 1
740 IF (Q3(I)-Q3(J)- 5.) 750,750,760
750 X3(J) = 0.0
J=J-1
IF (J-IN) 760,740,740
760 I=I+1
IF (I-ICX-1) 770,770,820
770 IF (X3(I)) 780,760,720
780 J=I
790 IF (Q3(J)-Q3(I-1)- 5.) 800,800,710
800 X3(J) = 0.0
J=J+1
IF (J-ICX-1) 790,810,810
810 I = J
GO TO 710
820 I=IN+1
C                                     PRAD4050
C                                     PRAD4060
C                                     PRAD4070
C                                     PRAD4080
C                                     PRAD4090
C                                     PRAD4100
C                                     PRAD4110
C                                     PRAD4130
C                                     PRAD4140
C                                     PRAD4150
C                                     PRAD4160
C                                     PRAD4170
C                                     PRAD4180
C                                     PRAD4190
C                                     PRAD4200
C                                     PRAD4210
C TEST TO FORM TRANSPORT REGIONS
C                                     PRAD4230
C                                     PRAD4240
IF (X3(IN)) 890,830,720
830 IAX=IN
840 IF (X3(I)) 860,850,720

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C           REMOVE ONE ZONE DIFFUSION REGION          PRAD4250
C
C   850 I=I+1                                     PRAD4260
    IF (I-ICX-1) 840,950,950
  860 I=I+1                                     PRAD4270
    IF (I-ICX-1) 870,950,950
  870 IF (X3(I)) 880, 875, 720                 PRAD4280
  875 X3(I-1) = 0.                               PRAD4290
    GO TO 840
  880 IBX=I-3                                     PRAD4300
    IF (IBX.LT.IAX) GO TO 970
    GO TO 960
  890 IF (IN.GT.1) GO TO 910                   PRAD4310
C   ASSUME C(1) = 0.
  900 X2(1) = 0.0
    FL(1) = 0.
    GO TO 920
  910 X2(IN) = 1.0285E12 + A(I,I) * (X6(IN-1) - X6(IN))
    FL(IN) = 0.5 * (X6(I,I-1) - X6(IN))          PRAD4330
  920 PR(IN) = Y2(IN) + 0.66666667                PRAD4340
    RHO(IN) = Y2(IN) + Y2(IN)
    FL(IN) = -.66666667 + TG(IN)
    TR(IN) = 0.6 * FL(IN)
  925 IF (X3(I)) 930, 940, 720
C   FORM X2 FOR DIFFUSION ZONES IN ORDER          PRAD4350
  930 X2(I) = -1.37E12 + TG(I) + A(I)          PRAD4360
    PR(I) = Y2(I) + .66666667
    RHO(I) = Y2(I) + Y2(I)
    FL(I) = -.66666667 + TG(I)
    TR(I) = .6 * FL(I)
    I=I+1
    IF (I = ICX - 1) 925, 980, 980             PRAD4410
C           DO TRANSPORT TO IM IN REGION OF NO SOURCE PRAD4420
C
C   940 IAX=I
    GO TO 850
  950 IBX=IR
  960 IF (IAX .GT. ICX) GO TO 965
    GO TO (961, 962, 963), 1ALPHA
  961 CALL PTRANS(IAX, IBX)
    GO TO 965
  962 S1 = 13.0962
    CALL UNCLE
  963 CALL STRANS(IAX, IBX)
  965 IF (IBX-IM) 970,1030,1030               PRAD4430
  970 I=IBX+2
    GO TO 930
  980 IF (I.GT.IR) GO TO 981
    IAX=I
    GO TO 950
  981 IF (IR .EQ. IM) GO TO 990               PRAD4440
    S1 = 13.0982
    CALL UNCLE

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C      RIGHT-HAND BOUNDARY CONDITION FOR DIFFUSION ZONES          PRAD4580
C      VEAT. MAY NEED RHO, FL, TR HERE.
990 IF (GL) 1000,1010,1020
,000 X2(IMP1) = 0.                                              PRAD4590
      GO TO 1030                                              PRAD4600
      GO TO 1030                                              PRAD4610
,010 X2(IMP1) = 1.0283E12 * X6(IM) * A(IMP1)
      GO TO 1025
1020 X2(IMP1) = 1.0283E12 * (X6(IM) - X6(IMP1)) * A(IMP1)
1025 RHO(IMP1) = Y2(IMP1) + Y2(IMP1)
      PR(IMP1) = 0.3333333 * RHO(IMP1)
      FL(IMP1) = X2(IMP1) / A(IMP1) * 4.8624E-13
      TR(IMP1) = FL(IMP1) * 0.6                                PRAD4650
C      OPTIONAL EDIT OF X2 ETC.                                    PRAD4660
C
1030 IF (ABS(EDITMF) .LT. 1.E-20 .OR. ZZ .GT. 1.E-20) GO TO 1040    PRAD4670
      CNT1=SOLID(18)+1.0
      IF(TD.GT.0.5.AND.CNT1.LT.CNTMAX)GO TO 1040
      WRITE (6,3) CNT1, TH, HNU, HNUP, IR, ICX, ICY
      WRITE (6,4)
      DO 1035 I = IN,IMP1
      WRITE (6,6) I, C(I), X6(I), H2(I), H3(I), TG(I), Y2(I), X3(I),
      1X2(I), RHO(I), PR(I), TH(I)
1035 CONTINUE
      3 FORMAT (9H1CYCLE = F7.0, 9H TIME = E13.6, 12H HNU FROM F8.2,
      2 4H TO F8.2, 10X2HRI4, 10X3HICX14, 10X3HICY14/)
      4 FORMAT (3X1HI, 1UX1HR, 10X2HX6, 10X2HH2, 10X2HH3, 10X2HTG,
      2 10X2HY2, 3X2HX5, 10X2HX2, 9X3HRHO, 10X2HPR, 10X2HTR)
      6 FORMAT (I4, 1PE11.4, 5E12.5, 0PF5.1, 1P4E12.5)
1040 IF (T4 .GT. CAPAC(152)) GO TO 1055
      ZZ=0.
      QEM = 0.
      DO 1053 I = IN, IMP1
      FIOSV=3.*RHO(I)-PR(I)
      FI2SV=3.*PR(I)-RHO(I)
      FI1SV = 3. * FL(I) - 5. * TR(I)
      FI3SV = 3. * TR(I) - 5. * FL(I)
      QN = FIOSV - FIU(I)
      QU = FIOSV + FIU(I)
      IF (ABS(QD) .GT. 0.) GO TO 1051
      GO TO 1053
1051 QE = ABS(QN / QD)
      IF (QE .LT. CAPAC(150)) GO TO 1052
      ZZ = 1.
      IF (QE .LT. QEM) GO TO 1052
      IQEM = I
      QEM = QE
1052 FI0(I) = FIOSV + CAPAC(151) + QN
      FI2(I) = FI2SV + CAPAC(151) + (FI2SV - FI2(I))
      FI1(I) = FI1SV + CAPAC(151) + (FI1SV - FI1(I))
      FI3(I) = FI3SV + CAPAC(151) + (FI3SV - FI3(I))
1053 CONTINUE
      T4 = T4 + 1.
      IF (T4 .GT. 9. .OR. ZZ .LT. 0.1) GO TO 1055
      CNT1 = SOLID(18) + 1.0

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      WRITE (6,8) CNT1, IHNNU, T4, OEM, IREM
8 FORMAT (29H SCATTERING ITERATION, CYCLE F5.0, 9H IHNNU = I2, 7H,
2T4 = F3.0, 32H LARGEST RELATIVE DIFFERENCE IS F5.3, 9H IN ZONE I3)
GO TO 515

1055 DO 1057 I = IN, IMP1
FQ0(I) = FI0(I)
FQ1(I) = FI1(I)
FQ2(I) = FI2(I)
FQ3(I) = FI3(I)
SUMX2(I)=SUMX2(I)+X2(I)                                     PRAD4810
EK(I) = RHO(I) * 68.5 + EK(I)
SMLR(I) = SMLR(I)+PR(I)*68.5
FI0(I) = 3. * RHO(I) - PR(I)
FI2(I) = 3. * PR(I) - RHO(I)
FI1(I) = 3. * FL(I) - 5. * TR(I)
1057 FI3(I) = 3. * TR(I) - 5. * FL(I)
Q = THICK * 1.026E12
CAPAR(IHNNU+135) = CAPAR(IHNNU+135) + 0.5 * DTR * (CAPAR(IHNNU+120) +
2 Q)
CAPAR(IHNNU+120) = Q
IF (ABS(SOLID(37)) .GT. 1.E-20) WRITE(JDRUMI) FI0, FI1, FI2, FI3
C ADVANCE FREQ, STORE EMERGENT FLUX, TEST FOR COMPLETION OF GROUPS PRAD4830
C                                                 PRAD4840
C                                                 PRAD4850
C                                                 PRAD4860
C HNUP=HNNU
C DHNUP = DHNU
C IHNNU = IHNNU + 1                                         PRAD4880
C T4 = 1.
C IF (IHNNU-NHNU) 1060,1060,1080                           PRAD4890
1060 CALL JVCHK (K000FX)                                     PRAD4900
GO TO (1070,310), KU00FX                                     PRAD4910
C*****END FREQUENCY LOOP*****                                PRAD4920
C                                                 *PRAD4930
C                                                 *PRAD4940
C                                                 *PRAD4950
C*****END*****                                              PRAD4960
1070 S1 = 13.1070                                           PRAD4970
CALL UNCLE
1080 SUMX2(INM1) = 0.0                                       PRAD4980
JURUM = JORUMI
REWIND JDRUM
DO 1090 I = IN, IM
EC(I) = 0.
X2(I) = SUMX2(I)                                             PRAD5020
RHO(I) = EK(I)
ER(I) = SUMX2(I) - SUMX2(I+1) + 0.5 * (A(I+1) * (3. * SMLR(I+1) -
2 EK(I+1)) - A(I) * (3. * SMLR(I) - EK(I)))
1090 CONTINUE
X2(IMP1) = SUMX2(IMP1)
RHO(IMP1) = EK(IMP1)
C EDIT OF OUTPUT OUTPUT
IF (TD .GT. 0.5 .AND. CNT1 .LT. CNTMAX) GO TO 1320
WRITE (6,5) CNT1, TH
WRITE (6,7) (CAPAR(I+120), I=1,NHNU)
WRITE (6,7) (CAPAR(I+135), I=1,NHNU)

5 FORMAT (25H OUTPUT OUTPUT FOR CYCLE F7.0, 10X7HTIME = 1PE13.6)
7 FORMAT (1P10E12.5)
1320 RETURN
END                                         PRAD6030

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WIT FOR PTRANS/P, PTRANS/P, PTRANS/P1
      SUBROUTINE PTRANS(N,M)
C COMPILED OCTOBER 9, 1967 WBL
C      PLANE CHARACTERISTIC TRACE WITH DOUBLE GAUSSIAN INTEGRATION      PTRAO030
C      USES THE RADIATION TRANSPORT 'STEP' SUBROUTINE
C***** PTRAO050
C* S P U T T E R   C O M M O N   **PTRAO060
C*                                     *PTRAO070
COMMON LMDA(37), NR      , NSMLR , IA      , IB      , ICA      , ICB      , PTRAO080
1 KMAX   , BLANK1, BLANK2, BLANK3, IAP1   , IOP1   , ICAP1   , ICBP1   , PTRAO090
2 II     , IG      , NRAD   , BLANK4, IAM1   , IBM1   , ICAM1   , ICBM1   , PTRAO100
3 IIP1   , IGM1   , IALPHA , BLANK5, TH     , TMAX   , BLANK6, DELPRT, PTRAO110
4 FREQ   , CNTMAX, AR      , ASMLR , PUSHB , BOILA , BOILB , PTRAO120
5 CVA    , CVB    , SLUG    , ALPHA  , HVA    , HVB    , HCA    , HCB    , PTRAO130
6 EMINA , EMINB  , CA      , CB     , GA     , GB     , GL     , GR     , PTRAO140
7 RHOL   , RHOR   , EP10   , EPS1   , RIA    , RIB    , RDIA   , RDIB   , PTRAO150
8 RPIA   , RPIB   , RPDIA  , RPDIB  , TPRINT, TA     , TB     , TC     , PTRAO160
COMMON TD      , TE      , DTH2   , DTH2P  , DTH1   , DTRMIN, DTMAX , PTRAO170
1 DTMAX1, DTMAX2, DTMAX3, UTR    , SWITCH , CO      , CMIN   , DELTA   , PTRAO180
2 GAMA   , WCRIT  , SIGMAQ , AC     , AC03T4, CNVRT , SUMRA  , SUMRB  , PTRAO190
3 ROI1A , ROIAM1 , ROIB   , ROIBP1 , GMS    , S1     , S2     , S3     , PTRAO200
4 S4     , S5     , S6     , S7     , S8     , S9     , S10   , S11   , PTRAO210
5 S12   , S13   , S14   , S15   , S16   , S17   , S18   , S19   , PTRAO220
6 S20   , E0     , FO     , TAU    , ZERO   , R     (152), DELTAR(152), PTRAO230
7 ASQ   (152), RD     (152), VD     (152), RDD   (152), SMLR (152), PTRAO240
8 DELR  ( 37), P      (152), P1    (152), PB    (152), PB1   (152) PTRAO250
COMMON P2     (152), SV     (152), RHO   (152), THETA (152), PTRAO260
1 W      (152), E      (152), EI     (152), EK     (152), A     (152), PTRAO270
2 V      (152), G      (152), D      (152), C      (152), X2     (152), PTRAO280
3 X3    (152), X4    (152), X5    (152), X6    (152), X7    (152), PTRAO290
4 SMLA  (152), SMLB  (152), SMLC  (152), SMLD  (152), SMLE (152), PTRAO300
5 EC     (152), ER     (152), SMLQ  (152), SMLH  (152), BIGA (152), PTRAO310
6 BIGB  (152), CV     (152), BC     (152), BR     (152), CHIC (152), PTRAO320
7 CHIR  (152), CAPAC (152), CAPAR (152), CRTC  (152), CRTR (152), PTRAO330
8 CRTPC (152), GOFR  (152), FEW   (152), CAR   (152), OKLM ( 37) PTRAO340
COMMON TELM ( 37), EKLM ( 37), ELM   ( 37), FCLM ( 37), PTRAO350
1 FRLM  ( 37), WLM   ( 37), QLM   ( 37), AMASNO( 37), CHRNO ( 37), PTRAO360
2 ZP1   ( 37), ZP2   ( 37), SOLID ( 37), ECICK ( 37), RK    (104), PTRAO370
3 RL    ( 37), RHOK (104), RDK   (104), THETAK(104), TEMP ( 16), PTRAO380
4 HEAD ( 12), MAXL   , MAXLM , PTRAO390
**PTRAO400
C***** PTRAO410
COMMON /LINDLY/ HNU,SGML,IHNU,NHNU,HNUP,NT,IM,IN,DHNU,THICK,NY PTRAO450
COMMON /CNTRL/ SCYCLE, JMULT PTRAO460
COMMON /DAVIS/ X(4000), ICX, ICY
COMMON /JIM/ NN, FMU, R1, R2, HD, EST, I1, I2, GMP, A1, A3, FMUS,
2 FS, LDF, LRI, IZN, TG1, TG2, F2
C***** PTRAO540
DIMENSION RR(40) PTRAO550
DATA RR/2.113248E-01,7.886752E-01,1.056624E-01,3.943376E-01, PTRAO560
1 1.127017E-01,5.000000E-01,8.872983E-01,3.130600E-02, PTRAO570
2 2.22222E-01,2.464718E-01,6.943180E-02,3.300095E-01, PTRAO580
3 6.699905E-01,9.305682E-01,1.207610E-02,1.076071E-01, PTRAO590
4 2.184655E-01,1.618513E-01,4.691010E-02,2.307653E-01, PTRAO600

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5      5.00000E-01, 7.092347E-01, 9.530899E-01, 5.557100E-03, PTRAO610
6      5.522540E-02, 1.422222E-01, 1.840889E-01, 1.129063E-01, PTRAO620
7      3.376520E-02, 1.093953E-01, 3.806903E-01, 6.193096E-01, PTRAO630
8      8.306047E-01, 9.662348E-01, 2.892400E-03, 3.055570E-02, PTRAO640
9      8.906520E-02, 1.448918E-01, 1.498251E-01, 8.276980E-02, PTRAO650
C      *PTRAO660
C *****CSUD (1), PR (1), FM (1), H (1), H2 (1),
1 H3 (1), H4 (1), FMS (1), Q1 (1), TG (1), Q3 (1),
2 Q37 (1), Q38 (1), SUMX2 (1), SUMX3 (1), SUMX4 (1), X8 (1), P 550
3 Y (1), Y2 (1), OX (1), FL (1), TR (1), FSM (1), P
4 FSP (1) P
C *****EQUIVALENCE (AC03T4,TRDBG), (BC ,SUMX4), (BIGA ,Y ),P 570
1 (BIGB , H ), (BR ,H3 ), (CAR ,Q37 ), (CHIC ,SUMX3),P 580
2 (CHIR ,Q38 ), (CRTR ,SUMX2), (X7 ,PR ), (GOFR ,Q3 ),P 590
3 ( PB ,Q1 ), (S12 ,EDITMF), (EC ,H2 ), ( W ,OX ),P 600
4 (SMLA ,FMS ), (SMLB ,FL ), (SMLC ,TR ), (SMLH ,H4 ),P
5 (ER ,FM ), (V ,TG ), (SMLD ,FSM ), (CRTC ,CSQD ),P
6 (X5 ,Y2 ), (X4 ,X8 ), (SMLE ,FSP ) P
C *****OX CONTAINS X FROM THE PREVIOUS Y LINE P 660
C *****CSQD SAME AS CRTC *P 670
C *****EDITMF SAME AS S12 *P 680
C *****Q1 SAME AS PB *P 690
C *****FM SAME AS ER *P 700
C *****H SAME AS BIGB *P 710
C *****H2 SAME AS EC *P 740
C *****H3 SAME AS BR *P 760
C *****H4 SAME AS SMLH *P 770
C *****PR SAME AS X7 *P
C *****FMS SAME AS SMLA *P
C *****FL SAME AS SMLB *P
C *****TR SAME AS SMLC *P
C *****FSM SAME AS SMLD *P
C *****FSP SAME AS SMLE *P
C *****Y2 SAME AS X5 *P
C *****OX SAME AS W *P
C *****TG SAME AS V *P
C *****Q3 SAME AS GOFR *P
C *****Q37 SAME AS CAR *P 840
C *****Q38 SAME AS CHIR *P 850
C *****SUMX2 SAME AS CRTR *P 860
C *****SUMX3 SAME AS CHIC *P 870
C *****SUMX4 SAME AS BC *P 880
C *****TRDBG SAME AS AC03T4 *P 890
C *****Y SAME AS BIGA *P 910
C *****X8 SAME AS X4 *P
C *****PLANE S ONLY *P 940
C *****TRAN 930 *PTRAO910
C *****PTRAO920 *PTRAO930

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C,*****  

IAX=IN  
PTRA0940  

IBXA  
PTRA0950  

INMI=IN-1  
PTRA0960  

IMPI = IM + 1  
PTRA0980  

CALL UVCHK(K000FX)  
PTRA0990  

GO TO (100,110), K000FX  
PTRA1000  

100 S1=14.0100  
PTRA1010  

CALL UNCLE  
PTRA1020  

110 IBXP1=IBX+1  
PTRA1030  

IALPHA=ALPHA  
PTRA1040  

PTRA1050  

C  
C           ERROR IF NOT PLANE  
PTRA1060  

C  
GO TO (130,120,120), IALPHA  
PTRA1070  

120 S1=14.0120  
PTRA1080  

CALL UNCLE  
PTRA1090  

130 NY = LMDA(37) - 1  
PTRA1100  

NMU = (NY - 1) * (NY + 2) + 1  
PTRA1110  

NGS = NMU + NY + 1  
PTRA1120  

JJ = 0  
PTRA1130  

C  
C           DO POSITIVE ANGLES FIRST  
PTRA1140  

C  
140 I=IAX  
PTRA1150  

F2=0.0  
PTRA1160  

FS = 0.0  
PTRA1170  

FMU = RR(NMU)  
PTRA1180  

LRI = 1  
PTRA1190  

PTRA1200  

C  
C           IF IAX=IN TRANSFER TO 150 TO SET SPECIAL BOUNDARY CONDITIONS  
PTRA1210  

C  
150 IF (IAX-IN) 360,150,180  
PTRA1220  

C  
C           CALCULATE BOUNDARY SOURCE INTENSITY  
PTRA1230  

C  
150 IF (INMI) 160,310,170  
PTRA1240  

C  
C           SET BLACKBODY CONDITION FOR PUSHER  
PTRA1250  

C  
160 S1=14.0160  
PTRA1260  

CALL UNCLE  
PTRA1270  

170 F2=X6(INMI)  
PTRA1280  

GO TO 310  
PTRA1290  

C  
C           DIFFUSION BOUNDARY CONDITION AT IAX  
PTRA1300  

C  
180 LDF = 1  
PTRA1310  

GO TO 220  
PTRA1320  

210 IF (ABS(TG(I-1)) .LT. 1.E-20) Y2(I-1) = X6(I-1)  
PTRA1330  

X8(I-1)=TG(I-1)*RH(NMU)  
PTRA1340  

C  
C           REGULAR INTEGRATION STEP FOR F2, POSITIVE MU  
PTRA1350  

C  
220 IF (ABS(TG(I)) .LT. 1.E-20) Y2(I) = X6(I-1)  
PTRA1360  

PTRA1370  

PTRA1380  

PTRA1500  

PTRA1510  

PTRA1520  

PTRA1530

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X8(I)=TG(I)*RR(NMU)          PTR A1550
I1 = 1 - 1
I2 = 1
R1 = C(I-1)
R2 = C(1)
12N = 1 - 1
TG1 = X8(I-1)
TG2 = X8(I)
HU = H2(I-1)/RR(NMU)
CALL STEP

260 SUMX3(I)=F2                PTR A1640
    IF (ABS(TG(I)) .LT. 1.E-20) Y2(I) = X6(I)
    I=I+1                         PTR A1660
    IF (I .GT. IBXP1) GO TO 320
    IF (I .LE. ICX + 1) GO TO 220

C
C           NO SOURCE IN ZONE GREATER THAN ICX
C,VEAT. SCATTERING IN SOURCELESS REGION NOT HANDLED PROPERLY HERE
1F (F2.EQ.0.0) GO TO 260          PTR A1780
TEMP(1)=H2(I-1)/RR(NMU)
H4(I-1)=FREXP(-TEMP(1)-TEMP(1))
F2=F2*H4(I-1)
GO TO 260                         PTR A1790
PTR A1800
PTR A1810
PTR A1820
PTR A1830
PTR A1840
PTR A1850
PTR A1860

300 IF (F2.EU.0.0) GO TO 310
TEMP(1)=H2(I-1)/RR(NMU)
H4(I-1)=FREXP(-TEMP(1)-TEMP(1))
F2=F2*H4(I-1)
310 SUMX3(I) = F2
FSM(1) = FS
LDF = 2
I=I+1
IF (I-1CY) 300,300,210          PTR A1870
PTR A1880
PTR A1890
PTR A1900
PTR A1910
PTR A1920

C
C           DO NEGATIVE ANGLES SECOND
C

320 I=IBXP1                      PTR A1930
FS = 0.0
LXI = 2
IF (IBX-1M) 370,330,360
330 IF (GL) 480,520,340          PTR A1940
C GL = 1/2 MEANS BLACKBODY CONDITION SET AT IMP1
C GL = POSITIVE INTEGER INPUT OPTION DELETED
C GL = 0 MEANS VACUUM AT IMP1
C GL NEGATIVE MEANS REFLECTIVE CONDITION AT IMP1
340 IF (GL.NE.0.5) GO TO 350
F2 = X6(IMP1)
GO TO 480                         PTR A1950
PTR A1970
PTR A1980
PTR A1990
PTR A2000
PTR A2010

350 S1 = 14.0350
CALL UNCLE                          PTR A2050
PTR A2060
PTR A2070
PTR A2080
PTR A2090
PTR A2100

C
C           ERROR IF INDEX EXCEEDS NORMAL RANGE
C

360 S1=14.0360
CALL UNCLE

```

```

C           DIFFUSION BOUNDARY CONDITION AT IBXP1          PTRAB2110
C
 370 LDF = 1
      GO TO 400
 399 IF (ABS(TG(I+1)) .LT. 1.E-20) Y2(I+1) = X6(I)          PTRAB2120

C           REGULAR INTEGRATION STEP FOR F2, NEGATIVE MU          PTRAB2230
C
 400 IF (ABS(TG(I)) .LT. 1.E-20) Y2(I) = X6(I)
      I1 = I + 1
      I2 = I
      R1 = -C(I+1)
      R2 = -C(I)
      IZN = I
      TG1 = -X8(I+1)
      TG2 = -X8(I)
      HU = H2(I)/RR(NMU)
      CALL STEP
 440 SUMX4(I)=F2          PTRAB2240
C
C           FORM CONTRIBUTION TO X2          PTRAB2250
C
      X2(I)=X2(I)-(F2-SUMX3(I))*RR(NGS)
      RH0(I) = RH0(I) + (F2 + SUMX3(I)) * RR(NGS) / RR(NMU)
      PR(I) = PR(I) + (F2 + SUMX3(I)) * RR(NGS) * RR(NMU)
      FL(I) = X2(I)
      TR(I) = TR(I) - (F2 - SUMX3(I)) * RR(NGS) * RR(NMU) * RR(NMU)
      IF (ABS(TG(I)) .LT. 1.E-20) Y2(I) = X6(I-1)
      I=I-1
      IF (I-IAZ) 530,450,450          PTRAB2330
 450 IF (I-ICY) 500,400,400          PTRAB2340
C
C           NO SOURCE IN ZONE LESS THAN ICY          PTRAB2350
C
 470 IF (F2.EQ.0.0) GO TO 480          PTRAB2360
      TEMP(1)=H2(I)/RR(NMU)
      H4(I)=FREXP(-TEMP(1)-TEMP(1))
      F2=F2*H4(I)
 480 SUMX4(I)=F2          PTRAB2370
 490 X2(I)=X2(I)-(F2-SUMX3(I))*RR(NGS)
      RH0(I) = RH0(I) + (F2 + SUMX3(I)) * RR(NGS) / RR(NMU)
      PR(I) = PR(I) + (F2 + SUMX3(I)) * RR(NGS) * RR(NMU)
      FL(I) = X2(I)
      TR(I) = TR(I) - (F2 - SUMX3(I)) * RR(NGS) * RR(NMU) * RR(NMU)
      LDF = 2
      FSP(I) = FS
      I=I-1
      IF (I-1-ICX) 399,470,470          PTRAB2420
C
C           NO SOURCE IN ZONE LESS THAN ICY          PTRAB2430
C
 500 IF (F2.EQ.0.0) GO TO 510          PTRAB2440
      TEMP(1)=H2(I)/RR(NMU)
      H4(I)=FREXP(-TEMP(1)-TEMP(1))
      F2=F2*H4(I)          PTRAB2450
C
C           NO SOURCE IN ZONE LESS THAN ICY          PTRAB2460
C
      PTRAB2470
      PTRAB2480
      PTRAB2490
      PTRAB2500
      PTRAB2510
      PTRAB2520
      PTRAB2530
      PTRAB2540
      PTRAB2550
      PTRAB2560
      PTRAB2570
      PTRAB2580
      PTRAB2590
      PTRAB2600
      PTRAB2610
      PTRAB2620
      PTRAB2630
      PTRAB2640

```

```

510 SUMX4(I)=F2
      X2(I)=X2(I)-(F2-SUMX3(I))*RR(NGS)          PTR A2650
      RHO(I) = RHO(I) + (F2 + SUMX3(I)) * RR(NGS) / RR(NMU)
      PR(I) = PR(I) + (F2 + SUMX3(I)) * RR(NGS) * RR(NMU)
      FL(I) = X2(I)
      TR(I) = TR(I) - (F2 - SUMX3(I)) * RR(NGS) * RR(NMU) * RR(NMU)
      I=I-1
      IF (I-IAX) 530,500,500
520 F2=0.0
      GO TO 480
530 CONTINUE
      IF (ABS(TRDBG) .LT. 1.E-20) GO TO 539
      C DEBUG PRINT OF INTENSITIES
      JJJ = JJ + 1
      IF (JJJ .EQ. 1) WRITE (6,8)
      WRITE (6,10) JJJ, RR(NMU)
      WRITE (6,9)
      DO 65 I = IAX, IBXP1                         REDI0920
65 WRITE (6,5) I, SUMX3(I), SUMX4(I)
      5 FORMAT (I4, 1PE11.4, E14.7 )
      8 FORMAT (28H1PLANE TRANSPORT DEBUG PRINT/)
      9 FORMAT (7X1HI, 7X7HI RIGHT, 8X6HI LEFT)        REDI0640
      10 FORMAT (27H GAUSSIAN QUADRATURE ANGLE I2, 17H WHOSE COSINE IS F11REDI0660
2.8)
      IF (TRDBG .GT. 68.) TRDBG = TRDBG - 69.          REDI0670
539 UHNU=HNUP-HNU
540 JJ = JJ + 1
      NMU = NMU + 1
      NGS = NGS + 1
      IF (JJ-NY) 140,140,550
550 DO 560 I=IAX,IBXP1
560 X2(I) = X2(I)+ 2.052E12
      RETURN
      END

```

Q,T FOR STRANS/S, STRANS/S, STRANS/S1  
 SUBROUTINE STRANS(N,M)  
 COMPILED OCTOBER 9, 1967 ABL  
 C MODIFIED FOR COMPTON SCATTERING AND IMPROVED LOGIC  
 C ANGULAR INTEGRATIONS ON LINEAR FORM  
 C CONNECTED LINEAR-QUADRATIC INTERPOLATION AT X=0.  
 C EPSI IS Y LIMIT OF EDIT TUBE  
 C LMUA(26) IS INTERFACE INDEX OF EDIT TUBE APERTURE

TRAN 20  
 TRAN 30

```
C*****  

C SPUTTER COMMON **  

C**  

COMMON LMUA(37), NR , NSMLR , IA , I8 , ICA , ICB ,  

1 KMAX , BLANK1, BLANK2, BLANK3, IAP1 , IBP1 , ICAP1 , ICBP1 ,  

2 II , IG , NRAD , BLANK4, IAM1 , IBM1 , ICAM1 , ICBM1 ,  

3 IIP1 , IGN1 , IALPHA, BLANK5, TH , TMAX , BLANK6, DELPRT,  

4 FREQ , CNTMAX, AR , ASMLR , PUSHA , PUSHB , BOILA , BOILB ,  

5 CVA , CVB , SLUG , ALPHA , HVA , HVB , HCA , HCB ,  

6 EMINA , EMINB , CA , CB , GA , GB , GL , GR ,  

7 RHOL , RHOR , EPIO , EPSI , RIA , RIB , ROIA , RUIB ,  

8 RP1A , RP1B , RP0IA , RP0IB , TPRINT, TA , TB , TC  

COMMON TU , TE , UTH2 , UTH2P , UTH1 , DTRMIN, DTMAX ,  

1 DTMAX1, DTMAX2, DTMAX3, UTR , SWITCH, CO , CMIN , DELTA ,  

2 GAMA , WCRIT , SIGMAU, AC , AC03T4, CNVRT , SUMRA , SUMRB ,  

3 RO1A , ROIAM1, ROIB , ROIBP1, GMG , S1 , S2 , S3 ,  

4 S4 , S5 , S6 , S7 , S8 , S9 , S10 , S11 ,  

5 S12 , S13 , S14 , S15 , S16 , S17 , S18 , S19 ,  

6 S20 , EO , FO , TAU , ZERO , R (152), DELTAR(152),  

7 ASG (152), RD (152), VD (152), RDU (152), SMLR (152),  

8 DELR (37), P (152), P1 (152), PB (152), PB1 (152)  

COMMON P2 (152), SV (152), RHO (152), THETA (152),  

1 W (152), E (152), EI (152), EK (152), A (152),  

2 V (152), G (152), D (152), C (152), X2 (152),  

3 X3 (152), X4 (152), X5 (152), X6 (152), X7 (152),  

4 SMLA (152), SMLB (152), SMLC (152), SMLD (152), SMLE (152),  

5 EC (152), EK (152), SMLQ (152), SMLH (152), BIGA (152),  

6 BIGB (152), CV (152), BC (152), BR (152), CHIC (152),  

7 CHIR (152), CAPAC (152), CAPAR (152), CRTC (152), CRTR (152),  

8 CRTPC (152), GOFK (152), FEW (152), CAR (152), OKLM (37)  

COMMON TELM (37), EKLM (37), ELM (37), FCLM (37),  

1 FRLM (37), WLM (37), QLM (37), AMASNO(37), CHRNO (37),  

2 ZP1 (37), ZP2 (37), SOLID (37), ECHCK (37), RK (104),  

3 RL (37), RHOK (104), RUK (104), THETAK(104), TEMP (16),  

4 HEAU (12), MAXL , MAXLM  

C**  

C***** TRAN 430  

C***** *TRAN 440  

C***** TRAN 450  

C***** TRAN 480  

C COMMON /LINDLY/ HIU,SGNL,IHNU,NHNU,HNUP,NT,IM,IN,DHNU,THICK,NY  

C COMMON /DAVIS/ X(4000), ICX, ICY
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COMMON /JIM/ NN, FMU, R1, H2, HD, EST, I1, I2, GMP, A1, A3, FMUS,
2 FS, LUF, LRI, LZN, TG1, TG2, F2
  DIMENSION CSQU (1), PR (1), FM (1), H (1), H2 (1),
1 H3 (1), H4 (1), FMS (1), Q1 (1), TG (1), Q3 (1),
2 Q37 (1), Q38 (1), SUMX2 (1), SUMX3 (1), SUMX4 (1), X8 (1),P 550
3 Y (1), Y2 (1), OX (1), FL (1), TR (1), FSM (1),P
4 FSP (1) P 570
C
EQUIVALENCE (AC03T4,TRDBG), (BC ,SUMX4), (BIGA ,Y ),P 580
1 (BIGH , H ), (BR ,H3 ), (CAR ,Q37 ), (CHIC ,SUMX3),P 590
2 (CHIR ,Q38 ), (CRTH ,SUMX2), (X7 ,PR ), (GOFR ,Q3 ),P 600
3 ( PB ,Q1 ), (S12 ,EDITMF), (EC ,H2 ), ( W ,OX ),P
4 (SMLA ,FMS ), (SMLB ,FL ), (SMLC ,TR ), (SMLH ,H4 ),P
5 (ER ,FM ), (V ,TG ), (SMLD ,FSM ), (CRTC ,CSQD ),P
6 (XS ,Y2 ), (X4 ,X8 ), (SMLE ,FSP ) P 660
C ***** OX CONTAINS X FROM THE PREVIOUS Y LINE *P 680
C
C   CSQU SAME AS CRTC *P 700
C   EDITMF SAME AS S12 *P 710
C   Q1 SAME AS PB *P
C   FM SAME AS ER *P
C   H SAME AS BIGH *P 740
C   H2 SAME AS EC *P
C   H3 SAME AS BR *P 760
C   H4 SAME AS SMLH *P 770
C   PR SAME AS X7 *P
C   FMS SAME AS SMLA *P
C   FL SAME AS SMLB *P
C   TR SAME AS SMLC *P
C   FSM SAME AS SMLD *P
C   FSP SAME AS SMLE *P
C   Y2 SAME AS XS *P
C   OX SAME AS W *P
C   TG SAME AS V *P
C   Q3 SAME AS GOFH *P 830
C   Q37 SAME AS CAR *P 840
C   Q38 SAME AS CHIR *P 850
C   SUMX2 SAME AS CRTR *P 860
C   SUMX3 SAME AS CHIC *P 870
C   SUMX4 SAME AS BC *P 880
C   TRDBG SAME AS AC03T4 *P 890
C   Y SAME AS BIGA *P 910
C   X8 SAME AS X4 *P 940
C ***** TRAN 930
C   TRAN 940
C   TRAN1040
C   TRAN1050
C   TRAN1060
C
IAX=N
IBX=M
THICK = 0.0
ITUBE = LMDS(26)
IF (C(ITUBE) .LT. EPSI) GO TO 1
XTUBE = SQRT(CSQU(ITUBE) - EPSI**2)
CALL DVCHK(KX)

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```

GO TO (1,3), KX
1 S1=14.0001          TRAN1090
CALL UNCLE           TRAN1100
3 IBXP1=IBX+1         TRAN1110
IALPHA=ALPHA          TRAN1120
C
C           ERROR IF NOT SPHERE
C
GO TO (5,5,7), IALPHA
5 S1=14.0005          TRAN1130
CALL UNCLE           TRAN1140
C
C           SPHERE ONLY
C
7 I=IBXP1             TRAN1150
FS = 0.0               TRAN1160
FMU = 0.0              TRAN1170
LRI = 1                TRAN1180
IF (IBX - IM) .GT. 37, 8, 11
8 F2 = 0.0              TRAN1190
IF (GL .GT. 0.) F2 = X6(IM+1)
GO TO 23              TRAN1200
TRAN1210
TRAN1220

C
C           ERROR IF INDEX EXCEEDS NORMAL RANGE
C
11 S1=14.0011          TRAN1230
CALL UNCLE           TRAN1240
C*****CALCULATE Y = 0 RAY*****
C
C           LHS OF RAY FIRST, STORE F2 IN SUMX3.
C
13 I1 = I + 1          TRAN1250
I2 = I                TRAN1260
R1 = -C(I+1)           TRAN1270
R2 = -C(I)              TRAN1280
HD = H2(I)              TRAN1290
IZN = I
TG1 = -TG(I+1)
TG2 = -TG(I)
CALL STEP
C
C           SAVE LHS INTENSITIES IN SUMX3
C
23 SUMX3(I)=F2          TRAN1300
FSM(I) = FS            TRAN1310
IF (ABS(TG(I)) .LT. 1.E-20) Y2(I) = X6(I-1)
I=I-1                  TRAN1320
LDF = 2
IF (ABS(TG(I)) .LT. 1.E-20) Y2(I) = X6(I)
IF (I-IAX) 47,15,13
C
TRAN1330
TRAN1340
TRAN1350
TRAN1360
TRAN1370

TRAN1700
TRAN1710
TRAN1720
TRAN1730
TRAN1750
TRAN1770
TRAN1990

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C           DIFFUSION BOUNDARY CONDITION AT IBXP1          TRAN2000
C
37 IF (ABS(TG(I)) .LT. 1.E-20) Y2(I) = X6(I)          TRAN2010
LDF = 1
GO TO 13

C           RHS OF RAY                                     TRAN2380
C
C           47 I=IAX                                     TRAN2390
IAXP=IAX
LRI = 2
IF (IAX = IN) 11, 48, 51
48 IF (IN - 1) 11, 49, 50
49 IAXP=IN+1
RHO(1)=SUMX3(1)*2.
PR(1)=SUMX3(1)*.6667
GO TO 69
50 IAXP = IN
F2 = X6(IN-1)
GO TO 69

C           DIFFUSION BOUNDARY CONDITION AT IAX          TRAN2500
C
C           51 LDF = 1                                     TRAN2510
H4(I-1) = FREXP(-H2(I-1))
GO TO 59

C           REGULAR INTEGRATION STEP (Y=0,RHS)          TRAN2830
C
C           59 I1 = I - 1                               TRAN2840
I2 = I
R1 = C(I-1)
R2 = C(I)
HD = H2(I-1)
IZN = I - 1
TG1 = TG(I-1)
TG2 = TG(I)
CALL STEP

C           SAVE RHS INTENSITIES IN SUMX4             TRAN3120
C
C           69 SUMX4(I) = F2                         TRAN3130
FSF(I) = FS
OX(I)=C(I)
IF (ABS(TG(I)) .LT. 1.E-20) Y2(I) = X6(I)          TRAN3140
LDF = 2
I=I+1
IF (ABS(TG(I)) .LT. 1.E-20) Y2(I) = X6(I-1)
IF (I-IBXP1) 59, 59, 85

C*****TRAN3470
C           C O M P L E T E   Y = 0   R A Y          *TRAN3480
C
C*****TRAN3510
C*****TRAN3490
C*****TRAN3500
85 Y5QDP=0.0

```

```

JJ=1          TRAN3530
JJJ=1          TRAN3540
KK=1          TRAN3550
C             TRAN3560
C             TRAN3570
C             TRAN3580
C             TRAN3590
C             *****
C             SEARCH FOR Y-LINE NEAR HALF OF C(IAX)  TRAN3620
C             *TRAN3630
C             *TRAN3640
C             *TRAN3650
C             *TRAN3660
C             *TRAN3670
C             *TRAN3680
C             *TRAN3690
C             *****
C             SET UP Y-LINES
C             *****
C             FIRST, TEST IF Y LIES OUTSIDE OF ACTIVE MESH.
C             *****
C             87 IF (JJ .LE. NY) GO TO 88
C             S1 = 14.0087
C             CALL UNCLE
C             88 IF (C(IBXP1) .LT. Y(JJ) + 1.000001) GO TO 127
C             Y-LINE IS INSIDE ACTIVE MESH.
C             IF (C(IAX)-Y(JJ)) 89,107,97
C             89 IF (C(IT) .GT. Y(JJ) + .999999) GO TO 107
C             IT = IT + 1
C             GO TO 89
C             97 IF (C(IAX-1) - Y(JJ)) 98, 107, 98
C             98 IF (C1 .LT. 0.5 * (Y(JJ) + Y(JJ+1))) GO TO 103
C             JJ=JJ+1
C             KK=KK-IFIX(X(KK+1))+2
C             GO TO 87
C             *****
C             SEARCH FOR Y-LINE NEAR THREE QUARTERS OF C(IAX)
C             *****
C             103 C1=C1+0.25*C(IAX)
C             LCULABLE Y-LINE FOUND -- PROCEED
C             107 YSQDP=X(KK)
C             FMU = YSQDP
C             K=KK+IM-I3X+2
C             TEMP(5)=YSQDP-YSQD1
C             GO TO 167
C             *****
C             COMPLETE X2 INTEGRATION WHEN LAST Y-LINE USED
C             *****
C             127 DO 139 I=IAXP,IBXP1
C             *****
C             ANGULAR INTEGRATION OF X2 USING LINEAR INTERPOLATION
C             *****
C             FNL = SUMX4(I) - SUMX3(I)
C             XSQ = OX(I)*#2
C             X2(I) = X2(I) + FNL * (XSQ + XSQ)
C             129 IF (IBX-IM) 131,135,171
C             131 IF (I-IAXP) 171,139,137
C             133 XS=SGRT(CSQD(IBXP1+1)-CSQD(IBXP1))  TRAN4290
C                                         TRAN4300
C                                         TRAN4310

```

```

FM(IBXP1) = Y2(IBXP1) + XS / C(IBXP1+1) * TG(IBXP1+1) +
2 FHEAP(-XS * H(IBXP1))
GO TO 137
135 FM(IBXP1)=0.0
137 TEMP(5)=CSQD(IBXP1)-YSQD1
TEMP(11)=CSQD(I)-YSQD1
FU=(TEMP(11)*(FM(IBXP1)+F4(IBXP1))+(CSQD(IBXP1)-
1 CSQD(I))*(SUMX4(IAXP)+SUMX3(IAXP)))/TEMP(5)
FLX = SUMX3(I) + SUMX4(I)
FP = FLX + FU
FPL = FP + FLX
RHO(I)=RHO(I)+OX(I)*FP
PR(I) = PR(I) + XSQ * (FPL + FLX) * OX(I)
TR(I) = TR(I) + XSQ**2 * FNL * 4.
139 CONTINUE
C*****TRAN4500
C      COMPLETION OF X2 INTEGRATION AT *TRAN4510
C      END OF TRANS REGION *TRAN4520
C      *TRAN4530
C      *TRAN4540
C*****TRAN4550
DO 151 I=IAX,IBXP1
IF (I.EQ.1) GO TO 151
PR(I)=PR(I)+.08334/(CSQD(I)+C(I))
RHO(I)=RHO(I)+.50/C(I)
FL(I) = .1666667 * X2(I) / CSQD(I)
IF (I.GT. 1) GO TO 147
TR(I) = 0.
GO TO 151
147 TR(I) = TR(I) + .05 / CSQD(I)**2
151 X2(I)=X2(I)*1.026E12
CALL DVCHK(KX)
GO TO (152, 153), KX
152 S1 = 14.0152
CALL UNCLE
153 RETURN
C*****TRAN4740
C      TYPICAL Y-LINE INTEGRATION *TRAN5230
C      *TRAN5240
C      *TRAN5250
C      *TRAN5260
C*****TRAN5270
C      LHS CALCULATION FIRST-STOKE F2 IN FM *TRAN5280
C      *TRAN5290
C      *TRAN5300
C      *TRAN5310
167 I=IBXP1
FS = 0.0
LRI = 1
IF (IBX - IM) 173, 168, 171
168 F2 = 0.0
IF (GL.GT. 0.) F2 = X6(IM+1)
GO TO 205
C      ERROR IF INDEX EXCEEDS NORMAL RANGE *TRAN5340
C      *TRAN5350
C      *TRAN5360
171 S1=14.0171 *TRAN5370

```

```

C CALL UNCLE
C
C DIFFUSION BOUNDARY CONDITION AT IBXP1
C
173 IF (ABS(TG(1)) .LT. 1.E-20) Y2(I) = X6(I)
    X8(I+1) = X(K-1) / C(I+1) * TG(I+1)
    LDF = 1
C
C SAVL X8 FOR RMS=(DIFF INTENSITY)
C
181 X8(I)=X(K)/C(I)+TG(I)
    HD = (X(K-1) - X(K)) * H(I)
C
C REGULAR INTEGRATION STEP(LHS)
C
I1 = I + 1
I2 = I
R1 = -X(K-1)
R2 = -X(K)
IZN = I
TG1 = -X8(I+1)
TG2 = -X8(I)
CALL STEP
C
C SAVE F3 OF LHS IN FM FOR INTEGRATION.
C
191 FM(I)=F2
    FSM(I) = FS
    IF (ABS(TG(I)) .LT. 1.E-20) Y2(I) = X6(I-1)
    I=I-1
    LDF = 2
    K=K+1
    IF (ABS(TG(I)) .LT. 1.E-20) Y2(I) = X6(I)
    IF (I-IAX) 239,193,193
193 IF (X(K)) 195,207,181
C Y-LINE HAS MADE CLOSEST APPROACH
195 TEMP(2)=SQRT(DELTAR(I)*(4.0+C(I)+DELTAR(I)))*H(I)
    IF (ABS(TEMP(2)) .LT. 1.E-20) GO TO 223
    HD = X(K-1) * H(I)
    TEMP(1) = HD + HD
    I1 = I + 1
    I2 = I
    R1 = -X(K-1)
    R2 = 0.
    IF (TEMP(1) = .02) 197, 197, 201
    THIN LHS XK=0.
197 TEMP(16)=(Y2(I+1)+X6(I))*0.5*TEMP(1)+(0.667*Y2(I)-0.5*Y2(I+1))
    1 -0.167*X6(I))*TEMP(2)
    NN = 0
    CALL SCAT
    Q = AMAX1(0., 1. - FMUS / H(I))
    F2 = FS + F2 * (1. - TEMP(1)) + Q * TEMP(16)
    GO TO 223
201 H4(I)=FREXP(-TEMP(1))
    NORMAL LHS XK=0.

```

```

NN = 1
TEMP(7)=FREXP(-TEMP(2))
TEMP(13)=(X6(I)-Y2(I))/TEMP(2)**2*2.0
TEMP(15)=(X6(I)-Y2(I+1))/(TEMP(1)-TEMP(2))
C AVEAT. H4 IS THE SQUARE OF ITS NORMAL VALUE.
EST = 1. - H4(I)
CALL SCAT
Q = AMAX1(0., 1. - FMUS / H(I))
F2 = FS + F2 * H4(I) + G * (Y2(I) + TEMP(13) + TEMP(7) * (-TEMP
2 (15) - TEMP(13) * (TEMP(2) + 1.)) + H4(I) * (TEMP(15) - Y2(I+1)))
GO TO 223
C
C          FIRST TRANSPORT ZONE ON LHS
C
205 X8(I)=X(K)/C(I)*TG(I)
GO TO 191
C
C          X = ZERO ERROR
C
207 S1=14.0207
CALL UNCLE
219 S1=14.0219
CALL UNCLE
221 F2=0.0
GO TO 225
C
C          UP Y-INTEGRATION TO X=0
C
223 IF (F2.LT.0.) GO TO 221
225 FM(I)=F2
FSM(I) = FS
C
C          FORM TOP SLICE CONTRIBS TO X2
C
DO 231 J=IAXP,I
TEMP(11)=CSQD(J)-YSQD1
FNL = SUMX4(J) - SUMX3(J)
XSQ = OX(J)**2
X2(J) = X2(J) + FNL * (XSQ + XSQ)
FU=(TEMP(11)*(FM(I)+F2)+(YSQDP-CSQU(J))*(SUMX4(IAXP)
1 +SUMX3(IAXP)))/TEMP(5)
FLX = SUMX3(J) + SUMX4(J)
FP = FLX + FU
FPL = FP + FLX
RHO(J)=KHO(J)+OX(J)*FP
PR(J) = PR(J) + XSQ * (FPL + FLX) * OX(J)
TR(J) = TR(J) + XSQ**2 * FNL * 4.
231 CONTINUE
LRI = 2
LDF = 2
IAXP=I+1
SUMX3(I)=F2
SUMX4(I)=F2
I=IAXP
K=K-1

```

TRAN6400  
TRAN6410  
TRAN6420  
TRAN6490  
TRAN6550  
TRAN6560  
TRAN6570  
TRAN6580  
TRAN6590  
TRAN6600  
TRAN6610  
TRAN6620  
TRAN6630  
TRAN6640  
TRAN6840  
TRAN6850  
TRAN6860  
TRAN6870  
TRAN6880  
TRAN6890  
TRAN6900  
TRAN6910  
TRAN6920  
TRAN6930  
TRAN6940  
TRAN6950  
TRAN6960  
TRAN7080  
TRAN7130  
TRAN7140  
TRAN7150  
TRAN7160  
TRAN7170  
TRAN7180

```

C      IF(ABS(TEMP(7)).LT.1.E-20)GO TO 257  DELETED 3/16/67, FCT
I1 = I - 1
I2 = I
R1 = 0.
R2 = X(K)
IF (TEMP(1) = .02) 235,233,237
C
C          SMALL OPTICAL DEPTH EXPANSION (X=0,RHS)           TRAN7240
C
C
233 NN = 0
CALL SCAT
Q = AMAX1(0., 1. - FMUS / H(I-1))
F2 = FS + F2 * (1. - TEMP(1)) + Q * TEMP(16)
GO TO 257
C          NORMAL RHS   (X=0)           TRAN7250
C
234 NN = 1
C        VEAT.  H4 IS THE SQUARE OF ITS NORMAL VALUE.
EST = 1. - H4(I-1)
CALL SCAT
Q = AMAX1(0., 1. - FMUS / H(I-1))
F2 = FS + F2 + H4(I-1) + (Y2(I) + TEMP(15) + FREXP(-TEMP(1) +
2 TEMP(2)) * (-TEMP(15) + TEMP(13) * (1. - TEMP(2))) - H4(I-1) *
3 (Y2(I-1) + TEMP(13)))
257 IF (F2 .LT. 0.) F2 = 0.
GO TO 259
C
C          CALCULATE RHS Y-LINE INTEGRATION           TRAN7270
C          ARRIVE HERE IF Y-LINE INTERSECTS INNER RADIUS
239 IAXP=IAX
I=IAX
K=K-1
LRI = 2
LDF = 1
IF (IAX .NE. IH .OR. IN .EQ. 1) GO TO 238
C          BLACKBODY
F2 = X6(IH-1)
GO TO 259
238 IF (X(K+1)) 240, 207, 241
C          DIFFUSION BOUNDARY CONDITION WHEN X=0.
240 X8(I-1) = 0.
HD = X(K) * H(I-1)
R1 = 0.
GO TO 242
241 X8(I-1)=X(K+1)/C(I-1)*T0(I-1)
HD = (X(K) - X(K+1)) * H(I-1)
R1 = X(K+1)
242 H4(I-1) = FREXP(-HD)
C
C          DIFFUSION BOUNDARY CONDITION AT IAX           TRAN7430
C
C          GO TO 250
249 HD = (X(K)-X(K+1))*H(I-1)
R1 = X(K+1)
250 I1 = I - 1
I2 = 1

```

```

R2 = X(K)
IZN = I - 1
TG1 = XB(I-1)
TG2 = XB(I)
CALL STEP
C
C      ANGULAR INTEGRATION OF X2 USING LINEAR INTERPOLATION
C
259 FNU = F2 - FM(I)
      FNLL = SUMX4(I) - SUMX3(I)
      FXM=OX(I)-X(K)
      X2(I) = X2(I) + ((X(K) + X(K) + OX(I)) * FNU + (OX(I) + OX(I) +
      2 X(K)) * FNLL) * FXM
      IF (I .NE. ITUHE) GO TO 263
      IF (EPS1 .GT. Y(JJ)) GO TO 261
      IF (EPS1 .LT. TEMP(9)) GO TO 263
      THICK = THICK + ((XTUHE + XTUHE + OX(I)) * F2 + (OX(I) + OX(I) +
      2 XTUHE) * SUMX4(I)) * (OX(I) - XTUHE)
      GO TO 263
261 THICK = THICK + ((X(K) + X(K) + OX(I)) * F2 + (OX(I) + OX(I) +
      2 X(K)) * SUMX4(I)) * (OX(I) - X(K))
263 FU=FM(I)+F2
      FLX = SUMX3(I) + SUMX4(I)
      TEMP(6)=OX(I)*OX(I)
      TEMP(7)=X(K)*X(K)
      FXP=OX(I)+X(K)
      FXP3 = FXP**3
      FM1 = FU - FLX
      HHO(I)=H(I)(I)+FXM*(FU+FLX)
      PR(I) = PR(I) + FXP + (TEMP(6) + TEMP(7)) * FM1 + 4. *
      2 (FLX + OX(I) + TEMP(6) - FU * X(K) + TEMP(7))
      TR(I) = THICK + FXM * (FNLL + (FXP3 + OX(I) * (3. + TEMP(6) -
      2 TEMP(7))) + FNU * (FXP3 + X(K) * (3. + TEMP(7) - TEMP(6))))
C
C      SAVE F2 AND FM FOR NEXT Y-LINE
C
      SUMX4(I)=T2
      SUMX3(I)=FM(I)
      FSF(I) = FS
      OX(I)=X(K)
      IF (ABS(TG(I)) .LT. 1.E-20) Y2(I) = X6(I)
      IZI+1
      KEK-1
      LDF = 2
      IF (ABS(TG(I)) .LT. 1.E-20) Y2(I) = X6(I-1)
      IF (I-IBXP1) 249,249,283
C
C      DEBUG PRINT
C
183 CONTINUE
      DHNU=HNUP-HNU
      IF (ABS(TRDBG) .LT. 1.E-20) GO TO 301
      CNTI = SOLIU(18) + 1.0
      WRITE (6,307) HNU, HNUP, IHNU, CNTI
      WRITE (6,309) Y(JJ)
      TRAN8120
      TRAN8130
      TRAN8140
      TRAN8340
      TRAN8350
      TRAN8360
      TRAN8380
      TRAN8390
      TRAN8410
      TRAN8420
      TRAN8740
      TRAN8750
      TRAN8760
      TRAN8780

```

```

      WRITE (6,311) JJ,IAX,IBX
      WRITE (6,315)
      IAXP1=MAX0(IAX,IAXP-1)
      DO 299 I=IAXP1,IAXP1
      IF (JJ.GT.1) GO TO 297
C
C           PRINT Y=0 INTEGRATION
C
      295 WRITE(6,313) I,C(I),X6(I),X8(I),H4(I),FSM(I),
     2FSP(I),SUMX3(I),SUMX4(I),X2(I)
      GO TO 299
C
C           PRINT REGULAR Y-LINE INTEGRATION
C
      297 KKK=KK-I+IM+3
      WRITE(6,313) I,X(KK),X6(I),X8(I),H4(I),FSM(I),
     2FSP(I),SUMX3(I),SUMX4(I),X2(I)
      299 CONTINUE
      IF (TRDBG .GT. 68.) TRDBG = TRDBG - 69.
      301 IF (JJ.EQ.1) GO TO 303
C***** EDIT AND Y-LINE ADVANCE *****
C
      KK=KK-IFIX(X(KK+1))+2
      303 YSD1 = YSDUP
      TEMP(9)=Y(JJ)
      JJ=JJ+1
      GO TO 87
      307 FORMAT (38H1Y-LINE PRINT FOR FREQUENCY BAND FROM F9.3, 5H TO F9,
     23, 10X7HINU = I2, 10X6HCYCLE F5.0)
      309 FORMAT (/5H Y = 1PE10.5)
      311 FORMAT (6H JJ = I3,10X7H IAX = I3,10X7H IBX = I3//)
      313 FORMAT (I4,1P4E14.7)
      315 FORMAT (3X1HI,13X1HX,12X2HX6,12X2HX8,
     212X2HH4,11X3HFSM,11X3HFSP,9X5HSUMX3,
     39X5HSUMX4,12X2HX2)
      END
C***** END *****
      TRAN8850
      TRAN8860
      TRAN8870
      TRAN8880
      TRAN8890
      TRAN8900
      TRAN8910
      TRAN8930
      TRAN8950
      TRAN8960
      TRAN8970
      TRAN8940
      TRAN8990
      TRAN9000
      TRAN9010
      *TRAN9020
      *TRAN9030
      *TRAN9040
      TRAN9050
      TRAN9060
      TRAN9080
      TRAN9090
      TRAN9110
      TRAN9130
      TRAN9210

```

QIT FOR STEP/BL, STEP/BL, STEP/BLI  
 SUBROUTINE STEP  
 COMPILED JUNE 22, 1967 WBL

```

C          S P U T I E R   C O M M O N      *T
C*          *T
C*          *T
C*          *T
COMMON  LMDA(37), NR    , NSMLR , IA    , IB    , ICA    , ICB    , T
1   KMAX , BLANK1, BLANK2, BLANK3, IAP1 , IBP1 , ICAP1 , ICBPI , T
2   II    , IG    , NHAD , BLANK4, IAM1 , IBM1 , ICAM1 , ICBM1 , T
3   IIP1 , IGM1 , IALPHA, BLANK5, TH    , TMAX , HLANK6, DELPRT, T
4   FREQ , CNTMAX, AR    , ASMLR , PUSHB , PUSHB , BOILA , BOILB , T
5   CVA  , CVB  , SLUG  , ALPHA , HVA  , HVB  , HCA  , HCB  , T
6   EMINA, EMINH , CA    , CB    , GA    , GB    , GL    , GR    , T
7   RHOL , RHOK , EPIO  , EPSI  , RIA  , RIB  , RDIA  , RDIB  , T
8   RPIA , RPIB , RPDIA , RPDIB , TPRINT, TA    , TB    , TC    , T
COMMON  TU    , TE    , UTH2  , DTH2P , UTH1  , DTRMIN, DTMAX , T
1   *DTMAX1, DTMAX2, DTMAX3, UTR  , SWITCH, CO    , CMIN  , DELTA  , T
2   GAMMA , WCRIT , SIGMAQ, AC    , AC03T4, CNVRT , SUMRA , SUMRB , T
3   ROI1A , ROIAM1, ROIH  , ROIHP1, GMS  , S1    , S2    , S3    , T
4   S4    , S5    , S6    , S7    , S8    , S9    , S10   , S11   , T
5   S12   , S13   , S14   , S15   , S16   , S17   , S18   , S19   , T
6   S20   , EO    , FO    , TAU   , ZERO  , R    (152) , DELTAR(152), T
7   ASQ   (152) , RU    (152) , VO    (152) , RDD   (152) , SMLR (152), T
8   DELR ( 37) , P     (152) , P1    (152) , PB    (152) , PB1  (152) , T
COMMON  P2    (152) , SV    (152) , RHO   (152) , THETA (152), T
1   W    (152) , E     (152) , EI    (152) , EK    (152) , A    (152), T
2   V.   (152) , G     (152) , U     (152) , C     (152) , X2   (152), T
3   X3   (152) , X4    (152) , X5    (152) , X6    (152) , X7   (152), T
4   SMLA (152) , SMLB (152) , SMLC (152) , SMLD (152) , SMLE (152), T
5   EC    (152) , ER    (152) , SMLQ (152) , SMLH (152) , BIGA (152), T
6   BIGB (152) , CV    (152) , BC    (152) , BR    (152) , CHIC (152), T
7   CH1K (152) , CAPAC (152) , CAPAR (152) , CRTC (152) , CRTR (152), T
8   CRTPC (152) , GOFK (152) , FEW   (152) , CAR   (152) , OKLM ( 37), T
COMMON  TELM ( 37) , EKLM ( 37) , ELM  ( 37) , FCLM ( 37), T
1   FRLM ( 37) , WLM  ( 37) , QLM  ( 37) , AMASNO( 37) , CHRNO ( 37), T
2   ZP1  ( 37) , ZP2  ( 37) , SOLID ( 37) , ECHCK ( 37) , RK    (104), T
3   RL-  ( 37) , RHOK (104) , RDK  (104) , THETAK(104) , TEMP ( 16), T
4   HEAD ( 12) , MAXL           , MAXLM           , T
C          T
C*          *T
C*          *****T
COMMON /JIM/ AN, FMU, R1, R2, HD, EST, I1, I2, GMP, A1, A3, FMUS, T
2 FS, LDF, LRI, IZN, TG1, TG2, F2           T
DIMENSION CSQU (1), PR (1), FM (1), H (1), H2 (1), T
1 H3 (1), H4 (1), FMS (1), Q1 (1), TG (1), Q3 (1), T
2 Q37 (1), Q38 (1), SUMX2 (1), SUMX3 (1), SUMX4 (1), X8 (1), T
3 Y (1), Y2 (1), OX (1), FL (1), TR (1), FSM (1), T
4 FSP (1)           T
C          T
EQUIVALENCE (AC03T4,TRDBG), (BC ,SUMX4), (BIGA ,Y ), T 570
1   (BIGB , H ), (BR ,H3 ), (CAR ,Q37 ), (CHIC ,SUMX3), T 580
2   (CHIR ,Q38 ), (CRTR ,SUMX2), (X7 ,PR ), (GOFK ,Q3 ), T 590
3   ( PB ,Q1 ), (S12 ,EDITMF), (EC ,H2 ), ( W ,OX ), T 600

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4      (SMLA ,FMS ), (SMLB ,FL ), (SMLC ,TR ), (SMLH ,H4 ),T
5      (ER ,FM ), (V ,TG ), (SMLD ,FSM ), (CRTC ,CSQD ),T
6.      (X5 ,Y2 ), (X4 ,X8 ), (SMLE ,FSP )          T
C*****OX CONTAINS X FROM THE PREVIOUS Y LINE*****T   660
C                                         *T   670
C                                         *T   680
C                                         *T   690
C                                         *T   700
C                                         *T   710
C                                         *T   740
C                                         *T   760
C                                         *T   770
C                                         *T   830
C                                         *T   840
C                                         *T   850
C                                         *T   860
C                                         *T   870
C                                         *T   880
C                                         *T   890
C                                         *T   910
C                                         *T   940
C*****TRAN 930*****T
C
C      GO TO (20, 25), LUF
C      IF DIFFUSION ZONE, DEFINE STARTING INTENSITY
C
20 F2 = Y2(I1) - TG1
25 IF (HD .GT. 0.01)  GO TO 30
C      THIN ZONE
NN = 0
GO TO 50
C      NORMAL ZONE -- UDEFINE EXP(-DELTATAU)
30 NN = 1
GO TO (35, 40), LRI
35 H4(IZN) = FREXP(-HD)
40 EST = 1. - H4(IZN)**2
50 CALL SCAT
Q = AMAX1(0., 1. - FMUS / H(IZN))
NN1 = NN + 1
GO TO (60, 70), NN1
60 F2 = FS + F2 * (1. - HD - HD) + Q * (((Y2(I1) + Y2(I2)) * 0.5 +
2 X6(IZN)) * HD)
GO TO 100
70 F2 = FS + F2 * H4(IZN)**2 + Q * (Y2(I2) - TG2 + ((TG1 - Y2(I1)) *
2 H4(IZN) + TG2 - TG1) * H4(IZN))
100 IF (F2 .LT. 0.)  F2 = 0.
RETURN
END

```

```

B,T FOR SCAT/PP,SCAT/PP,SCAT/JP          T
      SUBROUTINE SCAT                         T
C   GEOMETRY-INDEPENDENT COMPTON SCATTERING    T
C   COMPILED OCTOBER 9, 1967 WDL             T
C
C+*****+
C
C
COMMON LMDA(37), NR , NSMLR , IA , IB , ICA , ICB , T
1 KMAX , ULANK1, ULANK2, BLANK3, IAP1 , IUP1 , ICAP1 , ICBP1 , T
2 II , 1G , NRAU , BLANK4, IAM1 , IBM1 , ICAM1 , ICBM1 , T
3 ILPI , IGMI , IALPHA, BLANK5, TH , TMAX , BLANK6, DELPRT, T
4 FREQ , CNTMAX, AR , ASMLR , PUSHA , PUSHB , BUILA , BUILB , T
5 CVA , CVU , SLUG , ALPHA , HVA , HVB , HCA , HCB , T
6 EMINA , EMINH , CA , CB , GA , GB , GL , GR , T
7 RHOL , RHOR , EPIO , EPSI , RIA , RIB , RDIA , RDIB , T
8 RPIA , RP1B , RPUIA , RPDIB , TPRINT, TA , TB , TC , T
COMMON ID , TE , UTH2 , DTH2P , DTH1 , DTRMIN, DTMAX , T
1 DTMAX1, DTMAX2, DTMAX3, JTR , SWITCH, CO , CMIN , DELTA , T
2 GAMA , CHI1 , SIGMAJ, AC , AC03T4, CNVRT , SUMRA , SUMRB , T
3 ROLA , ROIAMI , ROIW , KOIBPI, GMS , S1 , S2 , S3 , T
4 S4 , S5 , S6 , S7 , S8 , S9 , S10 , S11 , T
5 S12 , S13 , S14 , S15 , S16 , S17 , S18 , S19 , T
6 S20 , EO , FO , TAU , ZERO , R (152), DELTAR(152), T
7 ASQ (152), RU (152), VD (152), RDD (152), SMLR (152), T
8 DELR (37), P (152), P1 (152), PB (152), PH1 (152) T
COMMON P2 (152), SV (152), RHO (152), THETA (152), T
1 W (152), E (152), EI (152), EK (152), A (152), T
2 V (152), G (152), D (152), C (152), X2 (152), T
3 X3 (152), X4 (152), X5 (152), X6 (152), X7 (152), T
4 SMLA (152), SMLB (152), SMLC (152), SMLD (152), SMLE (152), T
5 EC (152), ER (152), SMLQ (152), SMLH (152), BIGA (152), T
6 BIGD (152), CV (152), UC (152), BR (152), CHIC (152), T
7 CHIH (152), CAPAC (152), CAPAR (152), CRTA (152), CRTR (152), T
8 CRTPC (152), GOFK (152), FEW (152), CAR (152), OKLM (37) T
COMMON TELM (37), EKLM (37), ELM (37), FCLM (37), T
1 FRLM (37), WLM (37), QLM (37), AMASNO(37), CHRMIO (37), T
2 ZP1 (37), ZP2 (37), SOLID (37), ECHCK (37), RK (104), T
3 RL (37), RHOK (104), ROK (104), THETAK(104), TEMP (16), T
4 MEAU (12), MAXL , MAXLM T
C
C+*****+
C
C
COMMON /LINOLY/ HNU,SGML,INHNU,NHNU,HNUP,NT,IM,IN,DHNU,THICK, NY T
C
COMMON /CNTRL/ SCYCLE, JMULT T
COMMON /PALMER/ FIG(152), FI1(152), FI2(152), FI3(152), F00(152), T
2 FQ1(152), FQ2(152), FQ3(152), JURUM T
COMMON /JIM/ NN, FMU, R1, R2, HD, EST, I1, I2, GMP, A1, A3, FMUS, T
2 FS, LUF, LRI, IZN, TG1, TG2, F2 T
DIMENSION CSQD (1), PR (1), FM (1), H (1), H2 (1), T
1 H3 (1), H4 (1), FMS (1), Q1 (1), TG (1), Q3 (1), T
2 Q37 (1), Q38 (1), SUMX2 (1), SUMX3 (1), SUMX4 (1), X8 (1), T
3 Y (1), Y2 (1), OX (1), FL (1), TR (1), FSM (1), T

```



AFWL-TR-67-131, Vol III

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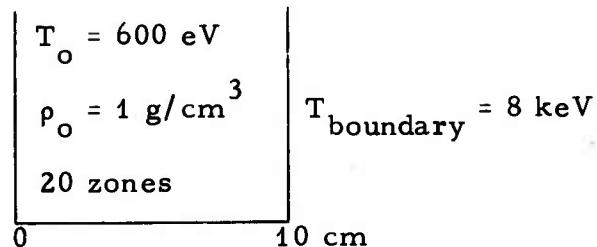
XP = R1 * R1
IF (YSQ .GT. 0.) GO TO 104
SQMU = 1.
FMUX = 1.
GO TO 106
104 SQMU = 0.5 * (XP / (XP + YSQ) + XX / (XX + YSQ))
FMUX = SQRT(SQMU)
106 DX = R2 - R1
110 FMUS = FMS(IZN)
BB = FI0(I1) - FI0(I2) + SQMU * (FI2(I2) - FI2(I1))
IF (ABS(SOLID(36)) .LT. 1.E-20) GO TO 115
C THOMSON SHORT CUT
IF (NN .NE. 0) GO TO 112
FS = .375 * FMUS * DX * ((FI0(I1) + SQMU * FI2(I1)) * (1. - HD) +
2 .5 * BB)
GO TO 25
112 AA = R2 * (FI0(I1) + SQMU*FI2(I1)) - R1 * (FI0(I2) + SQMU*FI2(I2))
FS = .375 * FMUS / (HD + HD) * (AA * EST + BB * (DX + R1 * EST) -
2 BB * DX * EST / (HD + HD))
GO TO 26
C NORMAL COMPTON PATH
115 DD = FI3(I2) - FI3(I1) + SQMU * (FI1(I2) - FI1(I1))
IF (IHNH .EQ. 1) GO TO 5
FF = FQ0(I2) - FQ0(I1) + SQMU * (FQ2(I2) - FQ2(I1))
HH = FQ3(I2) - FQ3(I1) + SQMU * (FQ1(I2) - FQ1(I1))
5 IF (NN .GT. 0) GO TO 10
C THIN RECIPE
FS1 = (1. - A1) * ((1. - HD) * (FI0(I1) + SQMU * FI2(I1)) + .5*BB)
FS2 = -A1 * FMUX* ((1. - HD) * (FI3(I1) + SQMU * FI1(I1)) + .5*DD)
IF (IHNH .EQ. 1) GO TO 9
FS3=A3*((1.-HD)*(FQ0(I1)+SQMU+FQ2(I1))+.5*FF)
FS4 = A3 * FMUX * ((1. - HD) * (FQ3(I1) + SQMU * FQ1(I1)) + .5*HH)
FS = .375 * FMUS * (FS1 + FS2 + FS3 + FS4) * DX
GO TO 25
9 FS = .375 * FMUS * (FS1 + FS2) * DX
GO TO 25
10 AA = R2 * (FI0(I1) + SQMU*FI2(I1)) - R1 * (FI0(I2) + SQMU*FI2(I2))
CC = R2 * (FI3(I1) + SQMU*FI1(I1)) - R1 * (FI3(I2) + SQMU*FI1(I2))
IF (IHNH .EQ. 1) GO TO 15
EE=R2*(FQ0(I1)+SQMU+FQ2(I1))-R1*(FQ0(I2)+SQMU+FQ2(I2))
GG=R2*(FQ3(I1)+SQMU+FQ1(I1))-R1*(FQ3(I2)+SQMU+FQ1(I2))
15 TERM2 = DX + EST * (R1 - DX / (HD + HD))
FS1 = (1. - A1) * (AA * EST + BB * TERM2)
FS2 = -A1 * FMUX * (EST + CC + DD * TERM2)
IF (IHNH .EQ. 1) GO TO 20
FS3=A3*(EE*EST + FF*TERM2)
FS4 = A3 * FMUX * (EST + GG + HH * TERM2)
FS=.375*FMUS/(HD+HD)*(FS1+FS2+FS3+FS4)
GO TO 25
20 FS=.375*FMUS/(HD+HD)*(FS1+FS2)
25 FMUS = FMUS * GMP
26 IF (FS .GT. (-1.E-20)) GO TO 29
IF (ABS(HVB) .LT. 1.E-20) GO TO 128
WRITE (6, 27) FS, IZN, IHNH, FMUX
27 FORMAT (6H FS = 1PE13.6, 7X6HIZN = I3, 7X7HIHNH = I2, 7X6HFMU = 1
2PE13.6)
WHITE (6,28) R1, R2, HD, EST, FS1, FS2, FS3, FS4, AA, BB, CC, DD,
2 EE, FF, GG, HH, FI0(I1), FI0(I2), FI1(I1), FI1(I2), FI2(I1),
3 FI2(I2), FI3(I1), FI3(I2)
28 FORMAT (1P8E15.6)
IF (TRDBG .LT. 68.) TRDBG = TRDBG + 69.
128 FS=0.
IF (ABS(CVB) .LT. 1.E-20) GO TO 29
S1=75.008
CALL UNCLE
29 CALL DVCHK(IMHAD)
GO TO (30,40),IMHAD
30 S1=75.0009
CALL UNCLE
40 RETURN
END

```

## APPENDIX IV

APPLICATIONS OF THE PLANE GEOMETRY CODE

The plane geometry code was applied in conjunction with the standard SPUTTER code to several one-dimensional problems. The problem definition can be summarized as follows:



The only physical process considered was radiation transport, i.e., no hydrodynamics. The equation of state for the material was that of  $\text{CH}_2$ . The opacity used was that of a hypothetical material in which  $\kappa_R = 0.2$  at all frequency groups.

The results of the calculations are shown in figures 8 through 13. Figure 8 shows the flux as a function of  $u = h\nu/kT$  for the Thomson scattering, i.e.,  $1/m_0 c^2 = 0$ . The difference between the curves at 0 and 2 mfp represents the energy deposited into the material. However, in the case of Thomson scattering, no material heating should occur. One should look at this discrepancy as a convergence problem. The results of figure 8 are plotted for cycle 20. Extending the calculation further would improve the accuracy at the cost of further computer time.

Figure 9 is the same calculation with Compton scattering. One should note the spectrum changes at the high-frequency end. This reduction in the spectrum for the Compton case is qualitatively in agreement with theory.

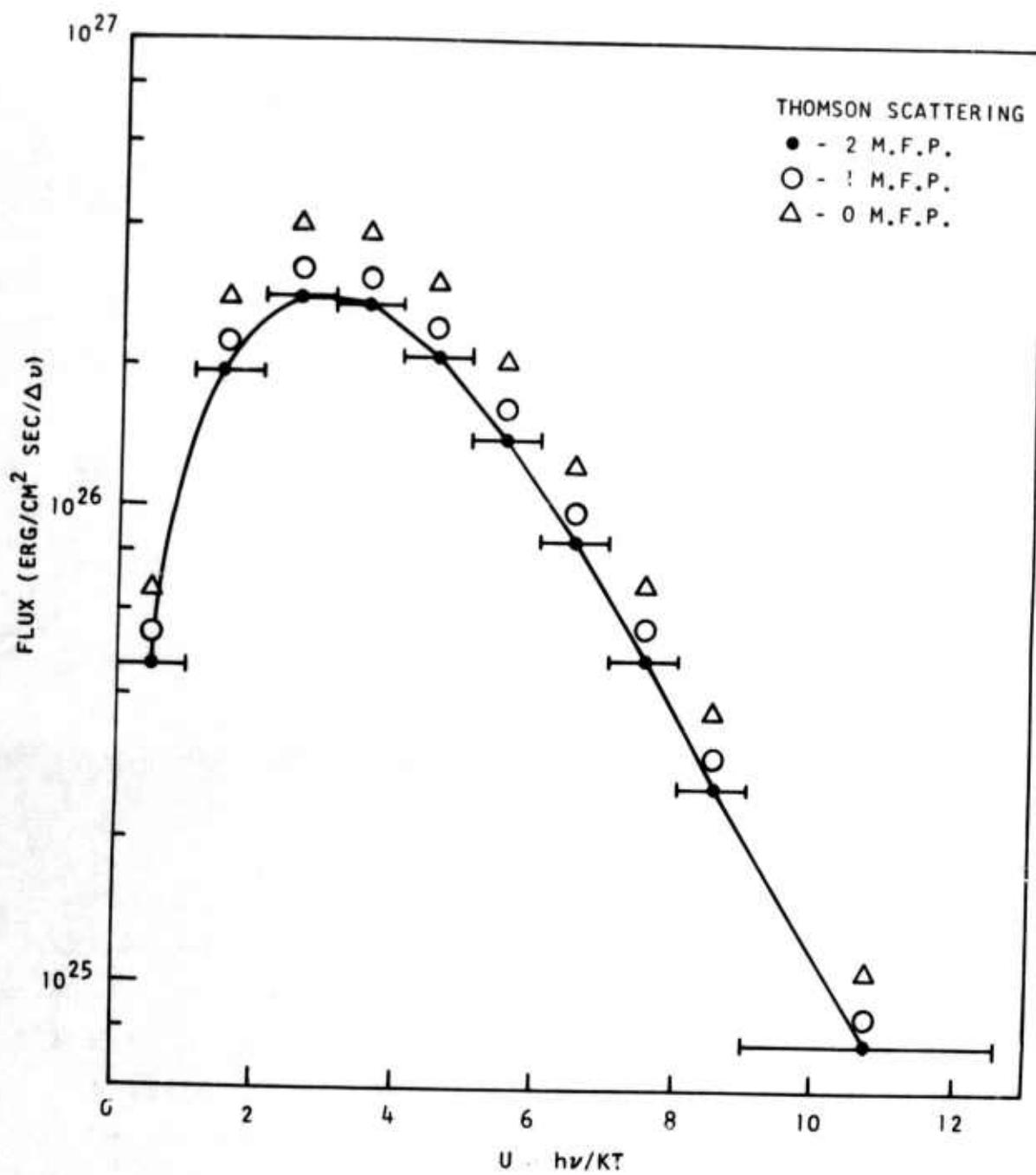


Figure 8. Spectrum at Various Optical Depths Employing Thomson Scattering

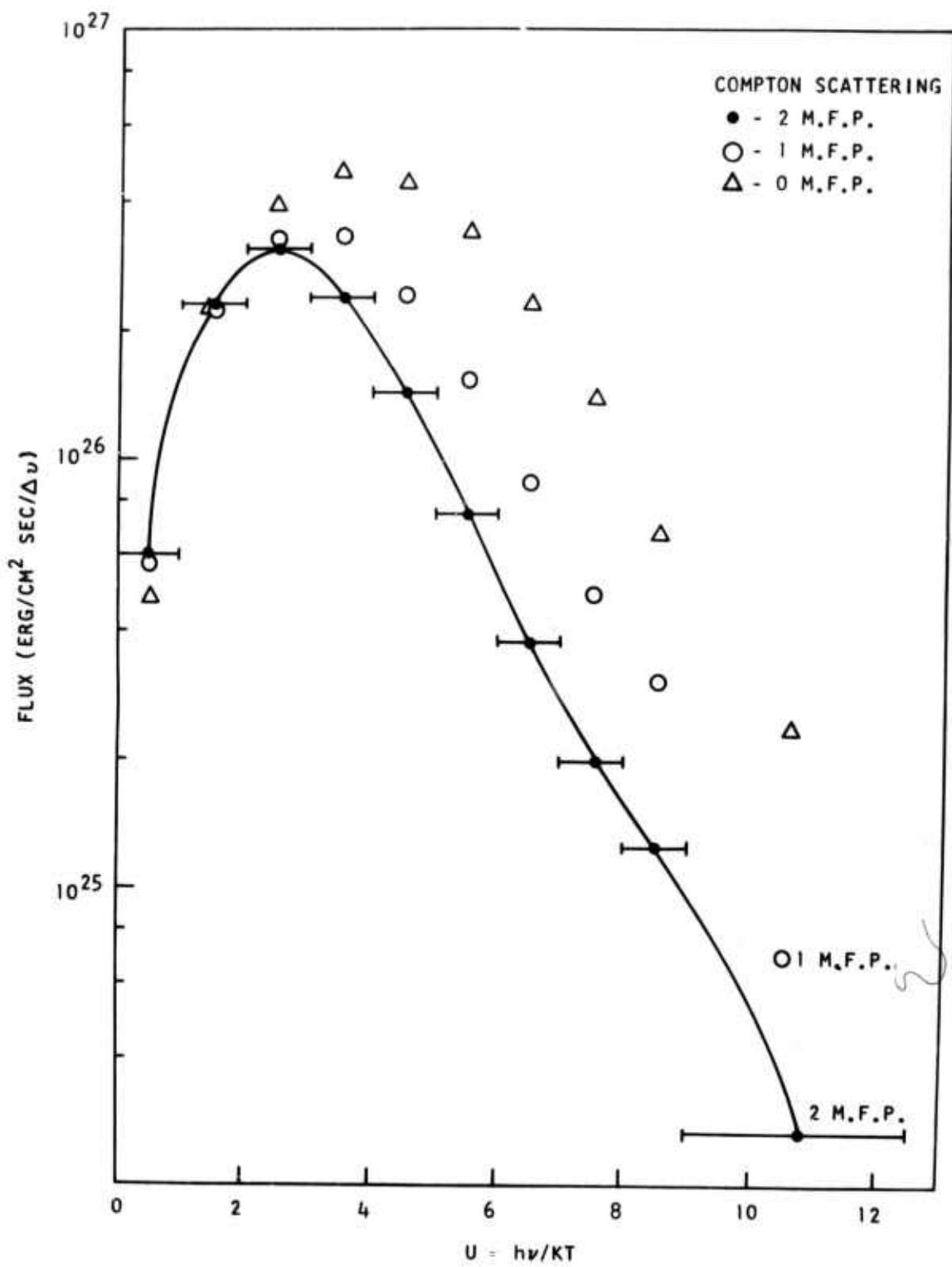


Figure 9. Spectrum at Various Optical Depths Employing Compton Scattering

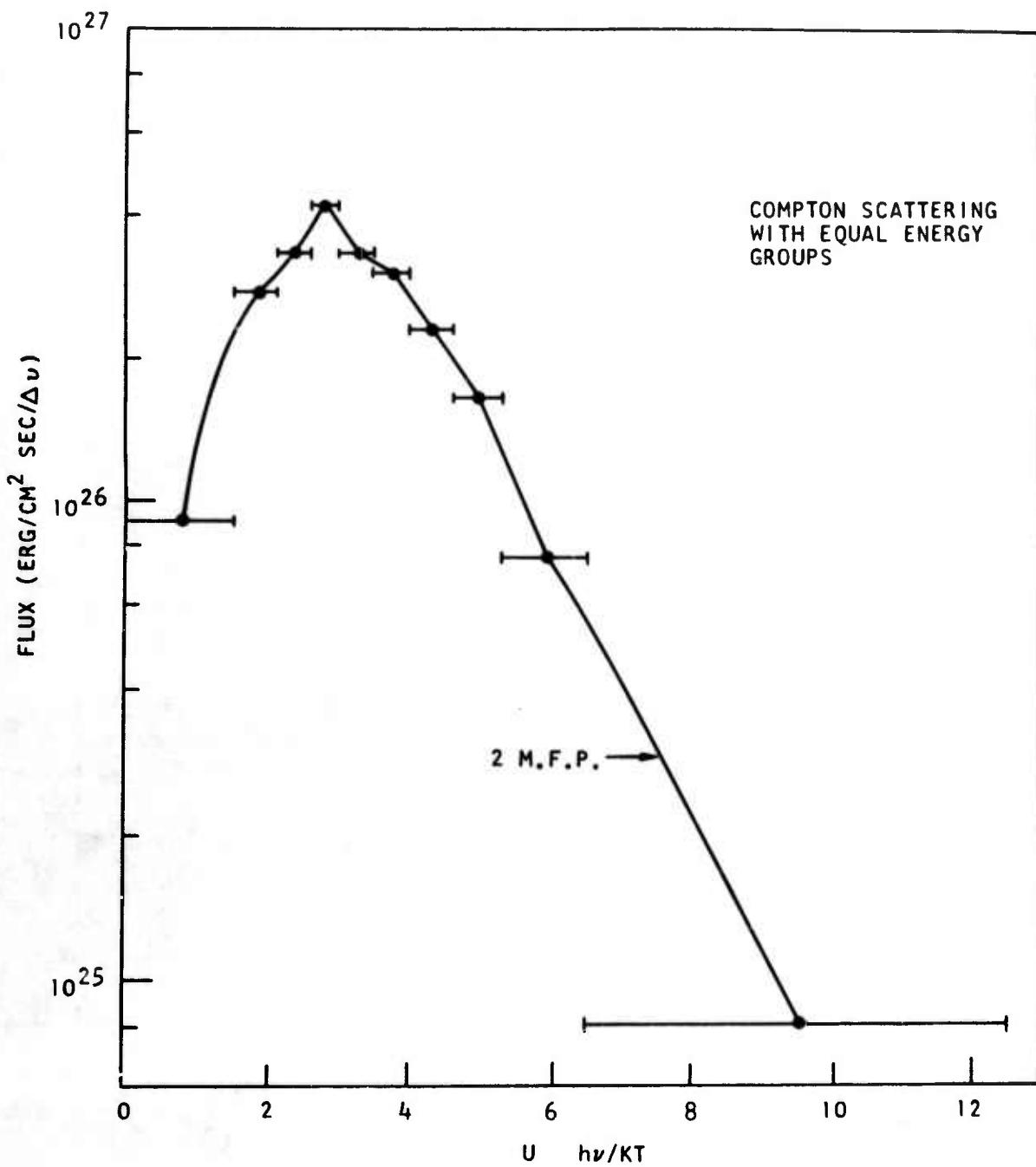


Figure 10. Exit Spectrum for Compton Scattering

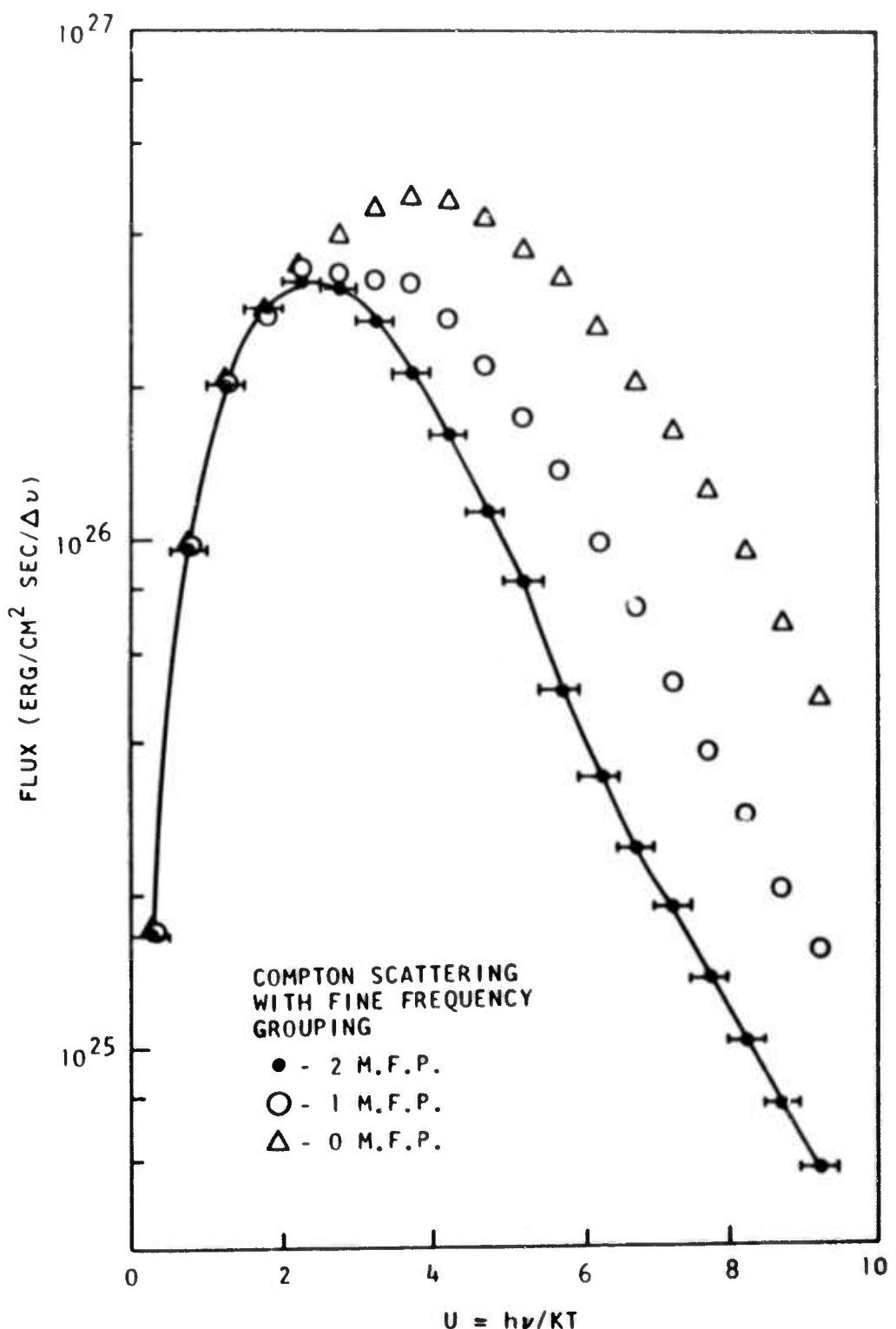


Figure 11. Spectrum using Compton Scattering

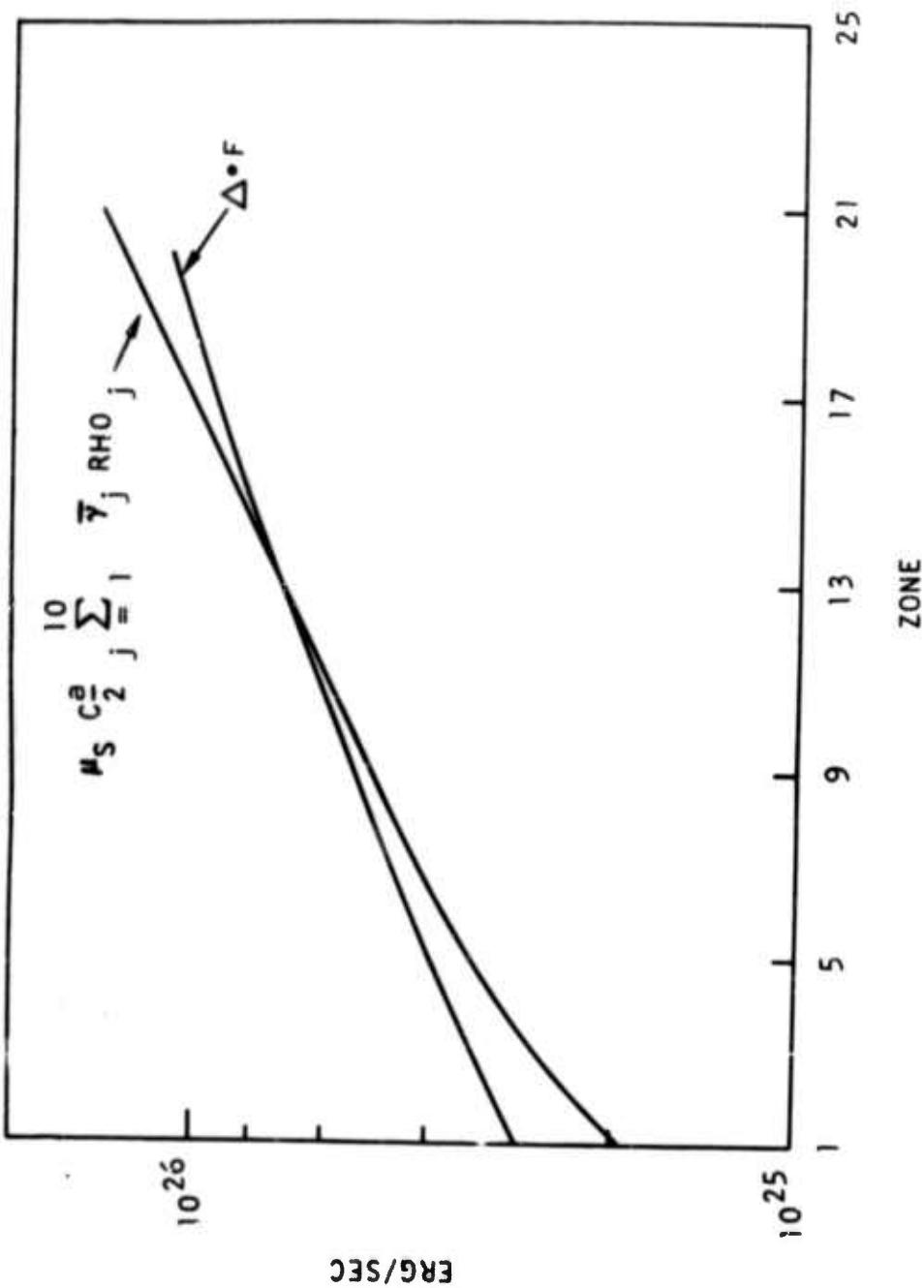


Figure 12. Comparison of Theoretical and Computed Heating Rates

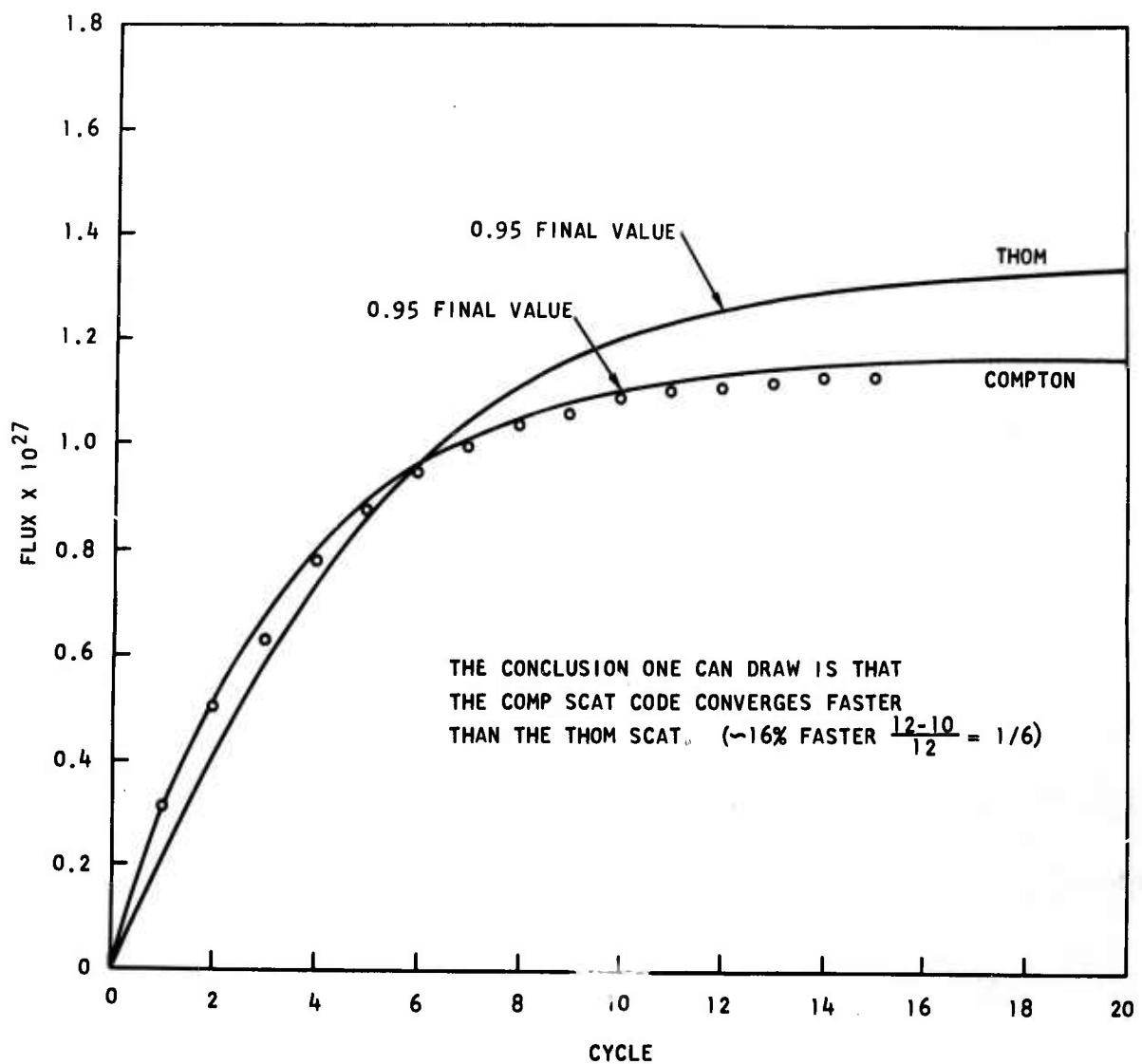


Figure 13. Comparison of Convergence Rates for Thomson and Compton Scattering

In the first two problems, illustrated in figures 8 and 9, the frequency groups were (0 to 8 keV), (8 to 16 keV), (16 to 24 keV), (24 to 32 keV), (32 to 40 keV), (40 to 48 keV), (48 to 56 keV), (56 to 64 keV), (64 to 72 keV), and (72 to 100 keV). The question arises as to how dependent is the solution on the number and definition of these groups. In figures 9 and 10, two different groupings were considered. In figure 10, the frequency groups were chosen to include 10 percent of the energy in an 8-keV blackbody. It is apparent from figure 10 that the spectrum has been distorted by unequal frequency groups. This conclusion was also supported by the fourth problem, i.e., the case in which the number of frequency groups was doubled. Initially, the first nine frequency groups were halved, whereas the (72 to 100 keV) group remained constant. This problem was unsuccessful. The unequal frequency grouping caused discontinuities to be propagated through the spectrum. This difficulty was eliminated by changing the (72 to 100 keV) group to (72 to 76 keV), where all the groups were then 4 keV wide. The results of this problem are shown in figure 11. Notice that figures 9 and 11 have converged to approximately the same solution. With the previous remarks in mind, one should be hesitant about solving problems in which the frequency groupings are unequal.

In figure 9, the difference between the curve at 2 and 0 mfp is the amount of energy used to heat the material. To determine if this is a reasonable estimate, a comparison of the material heating ratio (SMLE) and the theoretical heating rates

$$\left( \mu_s c \sum_{j=1}^N \bar{\gamma}_j E_j \right)$$

as a function of zone is presented in figure 12. The difference between these two curves can be used as a criterion for convergence.

A further point of interest is how fast the solution converges to its steady-state value at the edge of the window away from the source.

Figure 13 graphically displays this convergence.

For comparative purposes, a plot of the flux from an 8-keV black-body as a function of frequency is presented in figure 14.

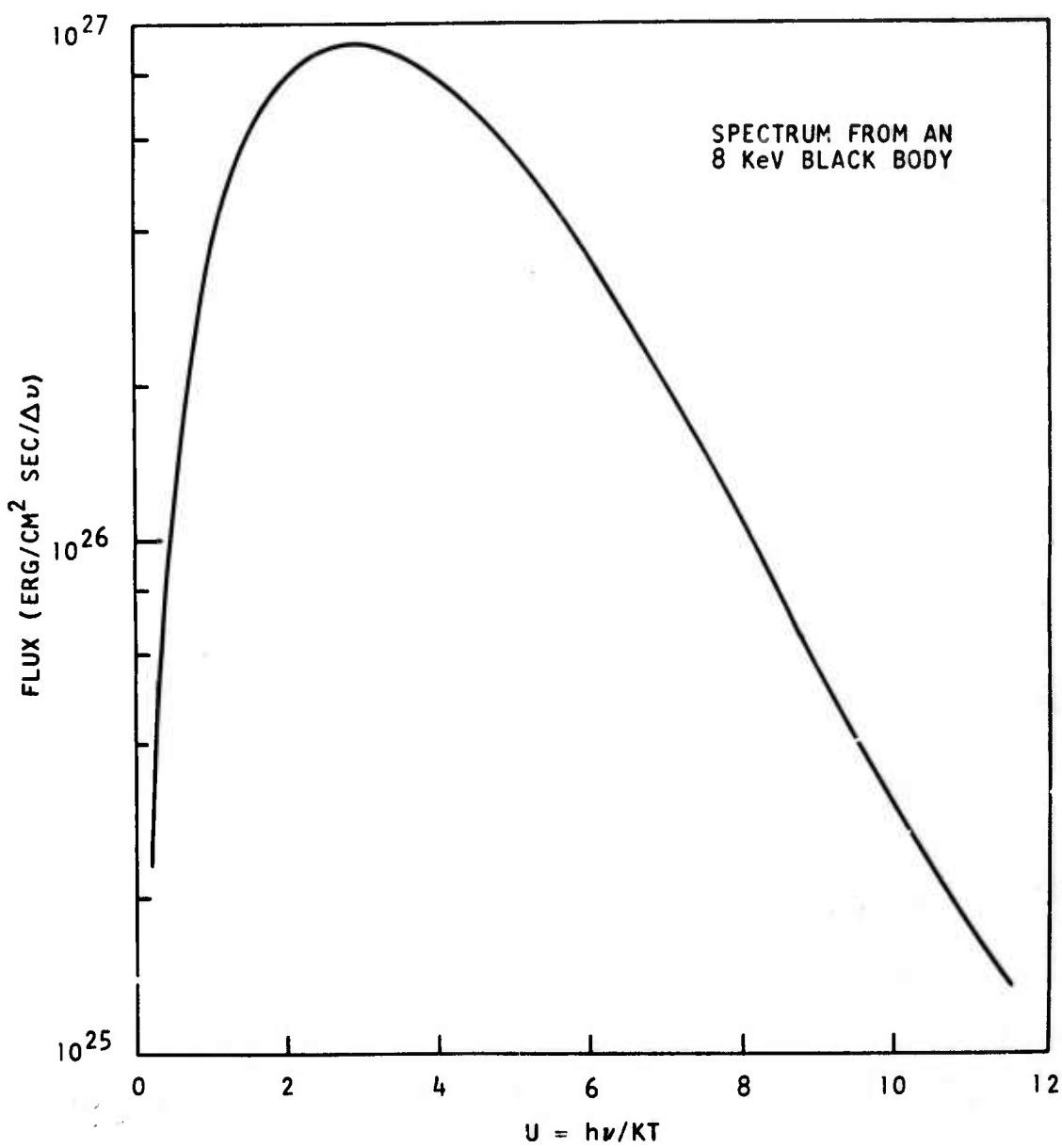


Figure 14. Spectrum from 8-keV Blackbody

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## APPENDIX V

### COMPTON AND INVERSE COMPTON SCATTERING

#### INTRODUCTION

Scattering of photons by free electrons in the Compton scattering process results in a modification of the radiative intensity in angular distribution and spectrum by virtue of the resulting energy and momentum exchange. This rate of change of intensity is described by the equation of radiative transfer, a Boltzmann equation for the photons. The objective here is to derive an approximation to the radiative-transfer equation which is valid when the scattering takes place from electrons in a Maxwellian distribution having temperature  $\theta_e$ . For this temperature to be maintained, it is necessary for electrons to undergo numerous energy-exchanging collisions between each photon scattering event. These collisions are then sufficient to maintain a Maxwellian distribution at the same temperature. Compton collisions with moving electrons may result in gain of photon energy (inverse Compton collisions) as well as energy loss. Both of these processes in the general case are necessary to describe the events which take place in a Compton scattering medium. In particular, the realization of a state of equilibrium between the photons and the electrons requires the inclusion of inverse Compton collisions and the direct Compton collisions to a consistent degree of approximation.

#### SCATTERING EQUATIONS

The equation for the transfer of photons in a scattering medium has been given in general form by Sampson (Ref. 8) and by Bond, Watson, and

Welch (Ref. 9). This equation has been applied by Freeman (Ref. 4) to the Compton scattering of photons by electrons at rest. In that work the equations are derived for photons having energies which are small compared with  $mc^2$ . An expansion is performed of the Klein-Nishina cross section through second order in the quantity  $\gamma = h\nu/mc^2$ . A further Taylor series expansion in frequency of the scattered intensity about the frequency of the photon beam in question is performed to obtain an "agelike" approximation. A similar treatment of the equation of transfer for scattering has been given by Fraser (Ref. 5). He takes into account the Compton scattering to first order in  $\gamma$  and, in addition, includes terms to first order in the quantity  $\alpha_e = \theta_e/mc^2$ , thereby including the inverse Compton effect in first order. The objective in this appendix is to reduce the results of Fraser to a form more suitable for calculation and to test his results in certain known limiting cases. The equation of scattering transfer will be derived first in a form suitable for application to a general geometric situation. The equations are then simplified by specialization to the plane and spherically symmetric geometries, and finally the diffusion approximation is derived.

The equations for the Compton scattering are contained in Fraser's expressions for  $\Gamma_1$ ,  $\Gamma_2$ , and  $\Gamma_3$  as given in his Eqs. (31), (32), and (33). When the drift velocity  $u_e$  is zero, these equations become

$$\Gamma_1 = -N\phi_o I(\nu, \Omega)(1 - 2\gamma)$$

$$\Gamma_2 = -N\phi_o \frac{3c^2}{16\pi h\nu^3} I(\nu, \Omega) \gamma \left(1 - \nu \frac{\partial}{\partial \nu}\right) \int d\Omega' I(\nu, \Omega') (1 - \mu + \mu^2 - \mu^3)$$

$$\begin{aligned} \Gamma_3 = -N\phi_o \frac{3}{16\pi} & \left\{ \gamma \left(1 - \nu \frac{\partial}{\partial \nu}\right) \int d\Omega' I(\nu, \Omega) (1 + \mu^2) (1 - \mu) \right. \\ & \left. - \left[ 1 + \alpha \left( 2 - 2\nu \frac{\partial}{\partial \nu} + \nu^2 \frac{\partial^2}{\partial \nu^2} \right) \right] \int d\Omega' I(\nu, \Omega) + \alpha \left( 4 - 2\nu \frac{\partial}{\partial \nu} + \nu^2 \frac{\partial^2}{\partial \nu^2} \right) \int d\Omega' I(\nu, \Omega') \mu \right. \\ & \left. - \left[ 1 - \alpha \left( 6 + 2\nu \frac{\partial}{\partial \nu} - \nu^2 \frac{\partial^2}{\partial \nu^2} \right) \right] \int d\Omega' I(\nu, \Omega') \mu^2 - \alpha \left( 4 + 2\nu \frac{\partial}{\partial \nu} - \nu^2 \frac{\partial^2}{\partial \nu^2} \right) \int d\Omega' I(\nu, \Omega') \mu^3 \right\} \quad (68) \end{aligned}$$

where  $N$  is the total number density of electrons and  $\phi_0 = 8\pi/3 (r_0^2) = 6.65 \times 10^{-25} \text{ cm}^2$  is the Thomson scattering. In Eq. (68), the values of  $M_0$ ,  $\underline{M}_1$ ,  $\underline{M}_2$ , and  $\underline{M}_3$  have been substituted for as follows:

$$M_0 = \frac{3}{16\pi} \int d\Omega' I(\nu, \Omega')$$

$$\underline{\Omega} \cdot \underline{M}_1 = \frac{3}{16\pi} \int d\Omega' I(\nu, \Omega') \mu$$

$$\underline{\Omega} \cdot \underline{M}_2 = \frac{3}{16\pi} \int d\Omega' I(\nu, \Omega') \mu^2 \quad (69)$$

$$\underline{\Omega} \cdot \underline{M}_3 = \frac{3}{16\pi} \int d\Omega' I(\nu, \Omega') \mu^3$$

$$M_0 - \underline{\Omega} \cdot (\underline{M}_1 - \underline{M}_2 + \underline{M}_3) = \frac{3}{16\pi} \int d\Omega' I(\nu, \Omega') (1 + \mu^2) (1 - \mu)$$

where the cosine of the angle of scattering is  $\underline{\Omega} \cdot \underline{\Omega}' = \mu$ . These terms represent the scattering contributions to the equation of radiative transfer for the intensity  $I(\nu, \Omega)$ . The first term represents the scattering out of the beam proportional to the total Klein-Nishina cross section expanded to first order in  $\gamma$ . The second term represents the contribution of stimulated scattering to both of the processes of scattering from the beam and scattering into the beam. The quantity  $\Gamma_3$  describes the scattering of photons into the beam by all of the electrons contained in the Maxwell distribution.

- A comparison shows that the  $\alpha$ -independent terms are the same as Freeman's first order in  $\gamma$  terms. These terms can be regrouped to display the order of the terms more effectively. Their sum, comprising the scattering contributions to the radiative transfer equation, is

$$\begin{aligned}
& - \mu_s \left[ (1 - 2\gamma) I(\nu, \Omega) - \frac{3}{16\pi} \left\{ \int d\Omega' I(\nu, \Omega') (1 + \mu^2) \right. \right. \\
& \quad - \gamma \left( 1 - \nu \frac{\partial}{\partial \nu} \right) \int d\Omega' I(\nu, \Omega') (1 + \mu^2) (1 - \mu) \\
& \quad + \alpha \left[ 2 \int d\Omega' I(\nu, \Omega') (1 - 2\mu - 3\mu^2 + 2\mu^3) \right. \\
& \quad \left. + \left( \nu^2 \frac{\partial^2}{\partial \nu^2} - 2\nu \frac{\partial}{\partial \nu} \right) \int d\Omega' I(\nu, \Omega') (1 + \mu^2) (1 - \mu) \right] \\
& \quad \left. - \frac{c^2}{h\nu^3} I(\nu, \Omega) \gamma \left( 1 - \nu \frac{\partial}{\partial \nu} \right) \int d\Omega' I(\nu, \Omega') (1 + \mu^2) (1 - \mu) \right] \quad (70)
\end{aligned}$$

where  $\mu_s = \phi_0 N$ .

In order to place the equation in conservative form in anticipation of formation of difference equations for numerical work, it is desirable to reformulate the frequency derivative terms. This reformulation corresponds to isolating those terms which contribute to exchange of energy within the photon distribution from the terms which contribute to exchange of energy between the electrons and photons. The desired substitutions are as follows:

$$\begin{aligned}
& \gamma \left( 1 - \nu \frac{\partial}{\partial \nu} \right) I = \frac{h}{mc^2} \left[ 3\nu I - \frac{\partial(\nu^2 I)}{\partial \nu} \right] \\
& \nu^2 \frac{\partial^2 I}{\partial \nu^2} - 2\nu \frac{\partial I}{\partial \nu} = \frac{\partial^2(\nu^2 I)}{\partial \nu^2} - 6 \frac{\partial(\nu I)}{\partial \nu} + 4I \quad (71)
\end{aligned}$$

For the case of spherical or plane geometry, symmetry of the radiative intensity about the polar direction can be assumed. This symmetry permits the integration of the azimuthal component of the solid angle in the photon scattering terms. If the results of Freeman (Ref. 4, p. 10) are used, the scattering terms become

$$\begin{aligned}
& - \mu_s \left[ (1 - 2\gamma) I(\nu, \Omega) - \frac{3}{16} \left( \int d\mu_3 I_3 [3 - \mu_1^2 + (3\mu_1^2 - 1)\mu_3^2] \right. \right. \\
& - \gamma \left( 1 + \frac{c^2}{h\nu^3} I(\nu, \Omega) \right) \left( 1 - \nu \frac{\partial}{\partial \nu} \right) \int d\mu_3 I_3 f(\mu_1, \mu_3) \\
& \left. \left. + \alpha \left\{ 2 \int d\mu_3 I_3 [(3\mu_1^2 - 1)(1 - 3\mu_3^2) + 2\mu_1(1 - 3\mu_1^2)\mu_3 + 2\mu_1(5\mu_1^3 - 3)\mu_3^3] \right. \right. \right. \\
& \left. \left. \left. + \left( \nu^2 \frac{\partial^2}{\partial \nu^2} - 2\nu \frac{\partial}{\partial \nu} \right) \int d\mu_3 I_3 f(\mu_1, \mu_3) \right\} \right] \right) \quad (72)
\end{aligned}$$

where

$$f(\mu_1, \mu_3) = 3\mu_1^2 - 1 + \mu_1(3\mu_1^2 - 5)\mu_3 + (3\mu_1^2 - 1)\mu_3^2 + \mu_1(3 - 5\mu_1^2)\mu_3^3$$

In Eq. (72), the inclusion of the additional terms for induced emission and inverse Compton scattering is not appreciably more difficult than for the Compton scattering term already considered. The same angular integrals enter as in Freeman's Eq. (18) (Ref. 4). The main difference arises in the presence of the second derivative with respect to frequency in addition to the first derivative occurring in the Compton formulation.

The diffusion approximation can also be carried out in parallel fashion to the derivation of Compton scattering for cold electrons. In this case two equations result, which are obtained from the zeroth and first moments of the equation of transfer. The Compton terms entering in the zeroth-moment equation are

$$\begin{aligned}
& - \mu_s \left[ -\gamma \left( I_o + \nu \frac{\partial I_o}{\partial \nu} \right) + \frac{c^2}{h\nu^3} \gamma \left( I_o^2 - \frac{2}{15} I_1^2 - \nu I_o \frac{\partial I_o}{\partial \nu} + \frac{2}{15} \nu I_1 \frac{\partial I_1}{\partial \nu} \right) \right. \\
& \left. - \alpha \left( \nu^2 \frac{\partial^2 I_o}{\partial \nu^2} - 2\nu \frac{\partial I_o}{\partial \nu} \right) \right] \quad (73)
\end{aligned}$$

and those for the first-moment equation are

$$\begin{aligned} -\mu_s \left[ I_1 \left( 1 - \frac{12}{5} \gamma \right) + \frac{2}{5} \gamma v \frac{\partial I_1}{\partial v} + \frac{c^2}{h\nu^3} \gamma \left( \frac{3}{5} I_0 I_1 + \frac{2}{5} v I_0 \frac{\partial I_1}{\partial v} \right. \right. \\ \left. \left. - v I_1 \frac{\partial I_0}{\partial v} \right) + \frac{2}{5} \alpha \left( I_1 + v^2 \frac{\partial^2 I_1}{\partial v^2} - 2v \frac{\partial I_1}{\partial v} \right) \right] \end{aligned} \quad (74)$$

### TESTS OF THE EQUATIONS

The equations derived above can be subjected to a number of tests to determine whether known results are obtained. Consider a scattering medium containing a weak radiation field for which the stimulated emission terms can be neglected. Integrations can then be carried out over the solid angle and the entire frequency spectrum. The resultant integrated radiative-transfer equation becomes an equation for the rate of change of radiation energy within a particular volume element. The Compton contributions to the rate of change represent the rate of loss or gain of energy from the radiation field. They also correspond to the gain or loss of energy by the electrons. These terms are

$$\begin{aligned} \int dv \int d\Omega (\Gamma_1 + \Gamma_3) = \mu_s c \left[ \int \gamma E_\nu dv + \int \gamma^2 \frac{\partial E_\nu}{\partial \gamma} dv + \alpha \int \left( v^2 \frac{\partial^2 E_\nu}{\partial v^2} - 2v \frac{\partial E_\nu}{\partial v} \right) dv \right] \\ = \frac{\mu_s}{mc} \left( 4\theta_e \int E_\nu dv - \int h\nu E_\nu dv \right) \end{aligned} \quad (75)$$

where  $E_\nu = 1/c \int d\Omega I(v, \Omega)$  is the spectral radiation energy. This is precisely the result given by Grasberger (Ref. 10).

The terms taking account of the stimulated scattering can also be included provided the radiation is isotropic. The contribution to the heating rate is

$$\int d\nu \int d\Omega R_2 = -\mu_s \frac{4\pi}{2m} \int I^2 \frac{d\nu}{\nu^2} = -\frac{\mu_s c^2}{2m 4\pi} \int E_v^2 \frac{d\nu}{\nu^2} \quad (76)$$

If, in addition, the radiation energy is given by a Planck distribution having a temperature  $\theta_r$ , as given by

$$I = B = \frac{2h}{c} \frac{\nu^3}{e^{h\nu/\theta_r} - 1} \quad (77)$$

the frequency integrations can be performed. The total resulting heating rate of the electrons is

$$\frac{dE_e}{dt} = \frac{16\mu_s}{mc} \sigma \theta_r^4 (\theta_r - \theta_e) \quad (78)$$

where  $\sigma$  is the Stefan-Boltzmann constant. This rate is applicable to a blackbody enclosure in which the radiation intensity is somehow maintained in a Planck distribution corresponding to a radiation temperature which is different from the temperature of the electrons within the enclosure. The derived electron heating rate is a well-known result (Ref. 10) which displays the characteristic that the radiation and electron temperatures will approach each other by virtue of this scattering interaction, yielding an equilibrium state in which the temperatures are equal.

These results test all of the Compton terms, including the induced scattering terms. Consequently, they constitute a strong confirmation of the correctness of these equations.

### APPROXIMATIONS

Several approximations have been made in deriving these equations which limit their applicability. First, the neglect of high-order terms in  $\gamma$  and  $\alpha$  restricts validity of the transfer equation to  $h\nu$  and  $\theta_e$  of the order of 100 kV. The inclusion of the terms corresponding to  $\gamma^2$  given in reference 4 would allow application of the equations to somewhat higher photon energies.

The second approximation, resulting from the expansion of the frequency dependence of the intensity and omission of quadratic terms in the frequency difference, implies that  $I_\mu$  must be a smooth function of frequency. Near a photoelectric edge or in the neighborhood of a line profile, these conditions may not be satisfied. An investigation of the validity of this approximation has been carried out by Chandrasekhar (Ref. 11), who evaluated the transmission of line radiation through a cold scattering atmosphere. His result shows that even in the case of a  $\delta$ -function source of radiation, a relatively small error, as measured by the fraction of energy erroneously scattered to higher frequencies (see Fig. 33, Ref. 11, p. 334), results (approximately 15 percent for an atmosphere containing two-thirds of a mean free path).

The neglect of the specific effects of polarization may also be of some consequence. Chandrasekhar (Ref. 12) has compared the diffuse reflection resulting from radiation incident on a semi-infinite scattering medium when the correlation of photon polarization after scattering is followed or neglected. Differences of the order of 5 percent in the scattered intensity are obtained (see Figs. 24 and 25, Ref. 12, pp. 262 and 263).

Scattering in the treatment of this report results from electrons which are free. Effects of binding of the electrons and the localization of the electrons within the atom give rise to modifications of the incoherent scattering Klein-Nishina formula for free-electron scattering and produce additional scattering which is coherent with the incident radiation. These effects are small when the energy of the photon is much greater than the binding energy of the atom and when the temperature of the material is sufficiently high that the probability of bound electrons being present is small.

EQUILIBRIUM SPECTRA

In this subsection, the solutions to the above equations corresponding to scattering of radiation in a homogeneous enclosure are examined. The equation for the spectral energy density is

$$\frac{1}{c} \frac{\partial I_o}{\partial t} = \mu'_a (B - I_o) - \mu_s \left[ -2\gamma I_o + \left( 1 + \frac{c^2}{h\nu^3} I_o \right) \left( I_o - \nu \frac{\partial I_o}{\partial \nu} \right) - \alpha \left( \nu^2 \frac{\partial^2 I_o}{\partial \nu^2} - 2\nu \frac{\partial I_o}{\partial \nu} \right) \right] \quad (79)$$

The steady-state solution of these equations should admit a Planck function for the radiation intensity corresponding to the temperature of the electrons. Substitution of the Planck function into Eq. (79) shows that such is the case. For the Planck function to satisfy this equation, however, it is necessary that the correct number density of photons be present. This condition will not in general be satisfied, however, if only the scattering interaction is present. More generally, the equilibrium will be achieved with too few or too many photons for the thermodynamic equilibrium solution to apply. If too few photons are present, the induced scattering-terms quadratic in  $I_o$  will be negligible, giving the steady-state equation

$$\nu^2 \frac{\partial^2 I_o}{\partial \nu^2} + \left( \frac{\gamma}{\alpha} - 2 \right) \nu \frac{\partial I_o}{\partial \nu} + \frac{\gamma}{\alpha} I_o = 0 . \quad (80)$$

This equation is satisfied by the function  $I_o = A\nu^3 e^{-h\nu/\theta}$ , the Wien approximation to the Planck function, but it has an arbitrary normalization, depending on the number of photons present.

If, on the other hand, there are too many photons for the thermodynamic equilibrium solution, only the quadratic terms need be retained,

$$-\gamma \frac{c^2}{h\nu} I_o v \frac{\partial I_o}{\partial v} + \gamma \frac{c^2}{h\nu} I_o^2 = 0 \quad \text{or} \quad I_o = v \frac{\partial I_o}{\partial v} \quad (81)$$

In this case, the solution is  $I_o = Bv$ , where  $B$  depends on the photon number density. However, this solution is valid only for frequencies such that  $v^2 \leq (c^2/h) B$ .

### CRITERIA FOR COMPTON HEATING

The inclusion of the Compton and inverse Compton scattering contributions to the heating rate is required only when the heating by pure absorption is negligible by comparison. In order to estimate the conditions where the Compton terms must be included, an estimate of the absorption contribution to the heating is derived and is then compared with the corresponding Compton scattering formula. As shown above, if the radiation intensity can be characterized by a Planck distribution with  $\theta_r$ , the electron heating rate from Compton scattering is

$$\left. \frac{dE_e}{dt} \right)_c = 2.1 \times 10^{-17} N_e \theta_r^4 (\theta_r - \theta_e) \quad (\text{ergs/cm}^3 \text{ sec}) \quad (82)$$

when temperatures are measured in eV and  $N_e$  in  $\text{cm}^{-3}$ .

At high temperatures, the pure absorption mechanism giving the largest contribution will be that from the free-free absorption. By using the absorption coefficient for the free-free process (Ref. 13), it is possible to derive a formula for the rate at which electrons exchange energy with the photons:

$$\left. \frac{dE_e}{dt} \right)_{ff} = 1.5 \times 10^{-25} Z^2 N_e N_i \theta_e^{-1/2} (\theta_r - \theta_e) \quad (\text{ergs/cm}^3 \text{ sec}) \quad (83)$$

where  $Z^2$  is the effective square of the ionic charge and  $N_i$  is the ionic number density.

The ratio  $R$  of the absorption heating rate to the scattering heating rate is given by

$$R = \frac{7.2 \times 10^{-9} Z^2 N_i}{\theta_e^{1/2} \theta_r^4} \quad (84)$$

A large value of  $R$  implies that scattering can be neglected in the electron heating rate. As an illustration of the application of this formula, one can obtain the temperature and density conditions at which heating due to scattering and absorption are equal. Figure 15 displays the temperature-density dependence of this condition,  $R = 1$ , for several typical materials, assuming that  $\theta_e = \theta_r$ . In applying this result, it should be kept in mind that deviations of the radiation spectrum from Planckian will strongly affect the answer.

The above criterion for the electron heating rate does not directly apply to the relative importance of absorption and scattering on the spectrum itself, since the frequency dependence of the absorption and scattering terms is quite different. If  $R \geq 1$ , the low-frequency portion of the spectrum will be dominated completely by absorption. The high-frequency portion of the spectrum will be modified predominantly by direct Compton scattering if  $\theta_r \gg \theta_e$  and by inverse Compton scattering when  $\theta_r \ll \theta_e$ .

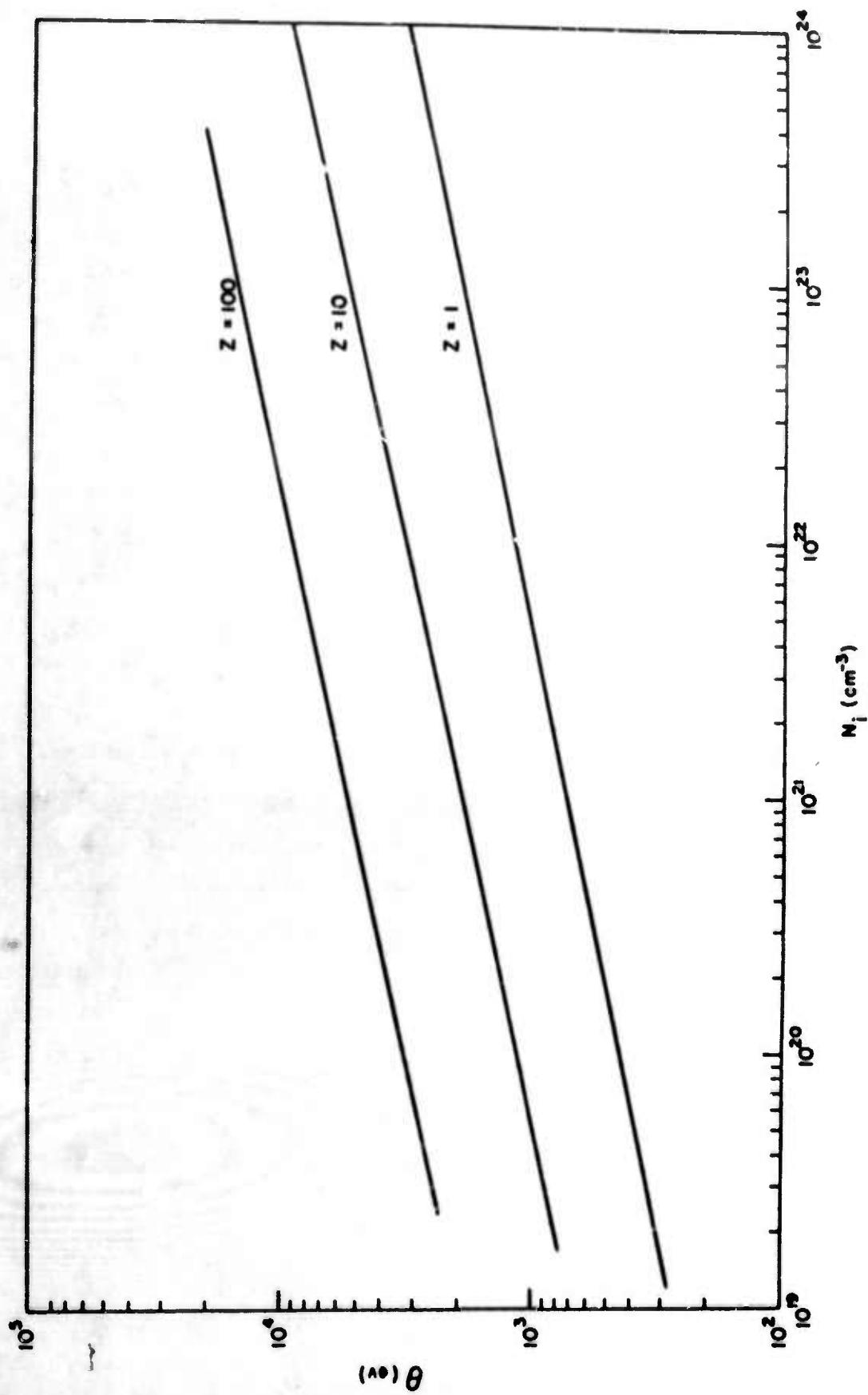


Figure 15. Temperature and Density at Which Heating Rates Due to Absorption and Scattering are Equal

## APPENDIX VI

A MODIFIED METHOD OF CHARACTERISTICS  
FOR RADIATIVE TRANSFER

In references 14 and 15, quantities termed "transmission functions" are proposed for use in the analysis of radiative transfer by the method of characteristics. These functions are mean values of the monochromatic attenuation factor  $\exp(-\rho \kappa_{\nu} \Delta x)$  for a range of frequencies  $\Delta\nu$ , where  $\Delta x$  is an interval along a ray within which the density  $\rho$  and opacity  $\kappa_{\nu}$  are assumed to be uniform. The frequency average of this exponential is taken with two different weighting functions  $B_{\nu}$  and  $dB_{\nu}/d\theta$ . In an idealized problem with uniform temperature only the first of these averages, denoted by  $S(\Delta x)$ , is needed. This function, defined by

$$S_{ij}(\Delta x_i) = \frac{\int_{\nu_j}^{\nu_{j+1}} B(\nu, \theta_i) \exp[-\rho_i \kappa_i(\nu, \rho_i, \theta_i) \Delta x_i] d\nu}{\int_{\nu_j}^{\nu_{j+1}} B(\nu, \theta_i) d\nu} \quad (85)$$

has a limiting form for very small  $\Delta x_i$  given by

$$S_{ij}(\Delta x_i) \sim \exp(-\rho \kappa_{ij}^P \Delta x_i) \quad (86)$$

where  $\kappa_{ij}^P$  is the usual Planck mean. For larger  $\Delta x$  (in most cases still rather small) the function flattens out rapidly; i.e.,  $-(1/\rho \Delta x) \log S$  is a decreasing function of  $\Delta x$ , eventually limiting to the minimum value of  $\kappa_{\nu}$  in the frequency group at sufficiently large  $\Delta x$ , plus terms of order  $\Delta x^{-1}$ .

In a medium with uniform composition, temperature, and density, the transport equation for the intensity in group  $j$  at a point  $x_n$  along a ray  $x_o < x < x_n$  due to radiation moving in the positive  $x$  direction is given in references 14 and 15 as

$$I_j(x_n) = B_j(x_n) + [I(x_o) - B(x_o)] S_j(x_n - x_o) \quad (87)$$

If the interval  $(x_o, x_n)$  is now subdivided into  $n$  zones of width  $\Delta x = (x_n - x_o)/n$ , the intensity at each of the subdivision points is

$$\begin{aligned} I_j(x_i) &= B_j(x_i) + [I(x_{i-1}) - B(x_{i-1})] S_j(\Delta x) \\ &= B_j(x_i) + [I(x_{i-2}) - B(x_{i-2})] S_j^2(\Delta x), \text{ etc.} \end{aligned} \quad (88)$$

By repeated application of Eq. (88), one readily obtains

$$I_j(x_n) = B_j(x_n) + [I(x_o) - B(x_o)] S_j^n(\Delta x) \quad (89)$$

which, with Eq. (87), implies that

$$S_j(n\Delta x) = S_j^n(\Delta x) \quad (90)$$

a condition which is satisfied if  $S$  is an exponential function, but is unfortunately violated by functions of the type described in references 14 and 15. With sufficiently fine subdivision of the medium in this example, the attenuation would in fact be that given by the Planck opacity, an incorrect result apart from the limiting case in which  $\rho\kappa^P(x_n - x_o) \ll 1$ .

This example exhibits the limitations resulting from one of the assumptions made in references 14 and 15, and in Eq. (87), namely that the frequency dependence of  $I(v)$  within each group is proportional to that of  $B(v)$  at every point along the ray. Actually, the frequency dependence of  $I(v)$  depends inherently upon conditions along the entire ray, and no method of frequency averaging based only upon locally evaluated temperature and density can be expected to be valid.

A second and unrelated difficulty occurs in spherical geometry, namely that consistent and simple models for the spatial dependence of the source function do not reduce to the correct diffusion limit, and other models which limit correctly are somewhat inconsistent and have a tendency to yield negative values for the intensity (Ref. 2, Vol III).

In this appendix a formulation is proposed which employs (1) nonlocal "transmission functions" of exponential form and (2) an apparently consistent model for the source function which has the correct diffusion limit. The method utilizes only the currently available Planck and Rosseland group mean opacities, but is capable of generalization to a formulation which, by describing the frequency dependence of opacity within each group in more detail, may allow use of fewer groups to attain the needed accuracy.

For monochromatic radiation, the intensity along a ray at optical depth  $\tau_n$  is given by

$$I(\tau_n) = I(\tau_o) e^{-(\tau_n - \tau_o)} + \int_{\tau_o}^{\tau_n} B e^{-(\tau_n - \tau)} d\tau \quad (91)$$

with integration by parts,

$$I(\tau_n) = B(\tau_n^+) + [I(\tau_o) - B(\tau_o^+)] e^{-(\tau_n - \tau_o)} - \int_{\tau_o}^{\tau_n} \frac{dB}{d\tau} e^{-(\tau_n - \tau)} d\tau \quad (92)$$

The last term, representing source gradient contributions to the intensity, may be represented as a sum of individual zone contributions:

$$\int_{\tau_o}^{\tau_n} \frac{dB}{d\tau} e^{-(\tau_n - \tau)} d\tau = \sum_{k=1}^n Q_{kn} \quad (93)$$

where

$$\begin{aligned} Q_{kn} &= \int_{\tau_{k-1}}^{\tau_k} \frac{dB}{d\tau} e^{-(\tau_n - \tau)} d\tau + (B_k^+ - B_k^-) e^{-(\tau_n - \tau_k)} \\ &= e^{-(\tau_n - \tau_k)} \left[ \int_{s_{k-1}}^{s_k} \frac{dB}{ds} e^{-\sigma_k(s_k - s)} ds + B_k^+ - B_k^- \right] \end{aligned} \quad (94)$$

Here  $s$  is the geometric distance along the ray,  $\sigma_k = \rho_k c_k$  is assumed constant in each interval  $s_{k-1} < s < s_k$ , and the possibility of discontinuous  $B(s)$  at interval boundaries is allowed for.

It is now necessary to represent the source gradient  $dB/ds$  in an explicit manner. Let the direction cosines of a ray with respect to Cartesian axes  $x$ ,  $y$ , and  $z$  be denoted respectively by  $\xi$ ,  $\eta$ , and  $\zeta$ ; then

$$ds = \frac{1}{\xi} dx = \frac{1}{\eta} dy = \frac{1}{\zeta} dz \quad (95)$$

The quadratic form

$$\begin{aligned} B(x, y, z) &= B(x_k, y_k, z_k) + \alpha_k(z - z_k) + \beta_k(x^2 + y^2 - x_k^2 - y_k^2) \\ &\quad + \gamma_k(x^2 + y^2 + z^2 - x_k^2 - y_k^2 - z_k^2) \\ &\quad x_{k-1} < x < x_k \\ &\quad y_{k-1} < y < y_k \\ &\quad z_{k-1} < z < z_k \end{aligned} \quad (96)$$

has plane symmetry if  $\beta_k = \gamma_k = 0$ , cylindrical symmetry if  $\alpha_k = \gamma_k = 0$ , spherical symmetry if  $\alpha_k = \beta_k = 0$ , and two-dimensional (finite) cylindrical symmetry if just  $\gamma_k = 0$ . This form may be simplified somewhat if the Cartesian axes can be chosen independently for each ray, so that, for example, the  $x$ -axis is parallel to the direction of the ray in spherical

geometry, or to the projection of the ray on the plane  $z = 0$  in cylindrical geometry.

From Eqs. (95) and (96) the source gradient along the ray is

$$\frac{dB}{ds} = \alpha_k \xi + 2\beta_k (\xi x + \eta y) + 2\gamma_k (\xi x + \eta y + \zeta z) \quad (97)$$

so that

$$Q_{kn} = \left[ \frac{\alpha_k}{\sigma_k} \xi + 2 \frac{\beta_k}{\sigma_k} (\xi \sigma_k x_k + \eta \sigma_k y_k - \xi^2 - \eta^2) + 2 \frac{\gamma_k}{\sigma_k} (\xi \sigma_k x_k + \eta \sigma_k y_k \right. \\ \left. + \zeta \sigma_k z_k - 1) + B_k^+ - B_k^- \right] e^{-(\tau_n - \tau_k)} \\ - \left[ \frac{\alpha_k}{\sigma_k} \xi + 2 \frac{\beta_k}{\sigma_k} (\xi \sigma_k x_{k-1} + \eta \sigma_k y_{k-1} - \xi^2 - \eta^2) + 2 \frac{\gamma_k}{\sigma_k} (\xi \sigma_k x_{k-1} \right. \\ \left. + \eta \sigma_k y_{k-1} + \zeta \sigma_k z_{k-1} - 1) \right] e^{-(\tau_n - \tau_{k-1})} \quad (98)$$

On comparing Eqs. (97) and (98), it is apparent that when  $\tau_n - \tau_{n-1} \gg 1$  the intensity limits to the value

$$I(\tau_n) = B(\tau_n) - \frac{dB}{d\tau} \Big|_{\tau = \tau'_n} \quad (99)$$

where for plane geometry  $\tau'_n = \tau_n$  and otherwise  $\tau'_n = \tau_n - 1$ . This is essentially consistent with the diffusion approximation.

So far, the analysis has been restricted to the single-frequency case. Equation (92) contains several types of terms, each with a different frequency dependence, corresponding to a variety of different materials, temperatures, and densities. Since frequency averaging destroys the separability of these contributions (as discussed above), approximation methods of a nonlocal character are needed.

The transmission functions of reference 15 can easily be generalized to nonlocal form. For example, the Planck transmission function defined

for a single zone by Eq. (85) is to be replaced by the kernel

$$S_j(s_i, s_n) = \int_{v_j}^{v_{j+1}} \sigma_i B(\theta_i, v) \exp \sum_{k=i}^{n-1} [-\sigma_k(v, \rho_k, \theta_k) \Delta s_k] dv / \int_{v_i}^{v_{j+1}} \sigma_i B(\theta_i, v) dv \quad (100)$$

which describes the attenuation of radiation emitted at  $s_i$  between  $s_i$  and a field point  $s_n$ . Of course, the direct evaluation of the terms in Eq. (100) is scarcely feasible in practice. Instead, a representation of the entire function is required which can be readily evaluated and which preserves some of the characteristics of Eq. (100), in particular, the thin limit

$$\lim_{s_n \rightarrow s_i} S_j(s_i, s_n) = 1 - (s_n - s_i) \sigma_j^P(s_n) \quad (101)$$

where  $\sigma_j^P(s_n)$  is the group Planck mean. A necessary condition for the diffusion limit is that for opacity independent of position

$$\lim_{s_i \rightarrow -\infty} \sigma_j^R(s_n) \int_{s_i}^{s_n} S_j(s, s_n) ds = 1 \quad (102)$$

where  $\sigma_j^R(s_n)$  is the group Rosseland mean. In general,  $S_j$  as a function of  $s_n$  should be uniformly positive, while the first derivative should be uniformly negative; also, the function should be short-range in the sense that for

$$s_i \rightarrow -\infty, \quad S_j(s_i, s_n) \sim (s_n - s_i)^{-m}, \quad m > 1 \quad (103)$$

where if  $m$  is not infinite, as for an exponential, it is at least large enough to restrict the significant contributions to the integral in Eq. (102) to source positions within a very few Rosseland mean free paths of the field point in all except pathologically nongrey cases.

The following construction is a nonlocal "picket-fence" (Refs. 16 and 17) transmission function of this type. It is assumed that the group

Planck and Rosseland means are known at each point, and that a parameter  $b_{kj}$  characteristic of the opacity distribution within group  $j$  at point  $s_k$  can be chosen. (For brevity, the group index  $j$  will be suppressed in the following formulas.) The first step is to construct the following composite opacities for each point:

$$\sigma^A = \frac{\sigma^P + b\sigma^R - \text{sgn}(1 - b^2) [(\sigma^P - \sigma^R)(\sigma^P - b^2\sigma^R)]^{1/2}}{1 + b} \quad (104)$$

$$\sigma^B = \frac{\sigma^P - b\sigma^R + \text{sgn}(1 - b^2) [(\sigma^P - \sigma^R)(\sigma^P - b^2\sigma^R)]^{1/2}}{1 - b} \quad (105)$$

where  $b$  is to be chosen so that  $\sigma^A$  and  $\sigma^B$  are real, and

$$0 < \sigma^A \leq \sigma^B \quad (106)$$

These conditions are satisfied if  $\sigma^P = \sigma^R$ , or if  $\text{sgn}(1 - b^2) = \text{sgn}(\sigma^P - \sigma^R)$ , i.e.,

$$\begin{aligned} b^2 &< 1 < \sigma^P / \sigma^R \\ \text{or} \end{aligned} \quad (107)$$

$$b^2 > 1 > \sigma^P / \sigma^R \quad (108)$$

Equation (108) only applies to cases which are nearly grey, or which have substantial scattering contributions in the Rosseland mean. Values of  $\sigma^A$  and  $\sigma^B$  for two ratios  $\sigma^P / \sigma^R$  are shown in figure 16.

The second step is to define the optical distances

$$\tau^A(s_i, s_n) = \int_{s_i}^{s_n} \sigma^A(s) ds \quad (109)$$

$$\tau^B(s_i, s_n) = \int_{s_i}^{s_n} \sigma^B(s) ds \quad (110)$$

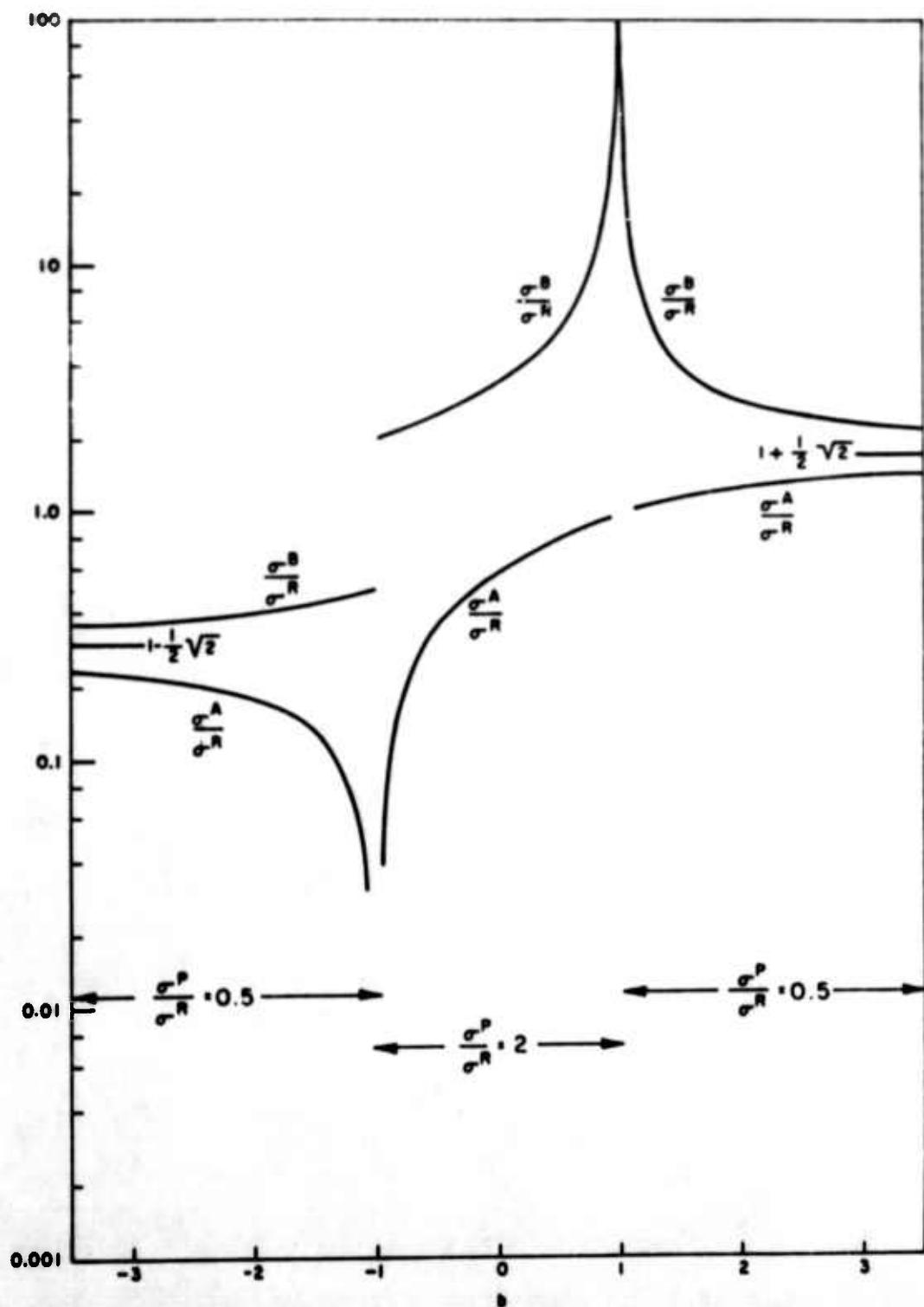


Figure 16. Dependence of Component Opacities on the Parameter  $b$

Next, the transmission function is represented by

$$S(s_i, s_n) = \frac{1}{2} (1 + b_n) e^{-\tau^A(s_i, s_n)} + \frac{1}{2} (1 - b_n) e^{-\tau^B(s_i, s_n)} \quad (111)$$

and, finally, the "nonlocal opacity" by

$$\begin{aligned} \sigma(s_i, s_n) &= -\frac{a}{ds_n} \log S(s_i, s_n) \\ &= \frac{1}{S(s_i, s_n)} \left[ \frac{1}{2} (1 + b_n) \sigma^A(s_n) e^{-\tau^A(s_i, s_n)} + \frac{1}{2} (1 - b_n) \sigma^B(s_n) e^{-\tau^B(s_i, s_n)} \right] \end{aligned} \quad (112)$$

The significance of the parameter  $b$  is now clearer. Assuming  $\sigma^P$  to be considerably larger than  $\sigma^R$  along the ray, two limiting cases can be distinguished. The nongrey character may be due to the presence within the group of a few strong line components, with the absorption minima, or windows, occupying most of the group width. The appropriate values of  $b$  will then be those near 1, so that according to Eq. (105) the "line" opacity  $\sigma^B$  is large compared with  $\sigma^P$ . The transmission, as shown in figure 17, attenuates at the Planck rate for a very short distance, with most of the contributions attenuated at a rate  $\sigma^A$  which is slightly less than the Rosseland mean. At the other extreme, there is the less probable situation in which the absorption maxima are broad, with the nongrey character due to the presence of isolated narrow but deep windows. For  $b \rightarrow -1$ ,  $\sigma^A \rightarrow 0$ , and  $\sigma^B \rightarrow \sigma^P$ . The transmission attenuates at nearly the Planck mean rate for perhaps several Planck mean free paths, but then levels out in such a way that small contributions are transmitted with little attenuation from remote sources. For intermediate cases, values of  $b$  near zero may be used. In particular, for  $b = 0$ ,

$$\left. \begin{array}{l} \sigma^A \\ \sigma^B \end{array} \right\} = \sigma^P \pm [\sigma^P(\sigma^P - \sigma^R)]^{1/2} \quad (113)$$

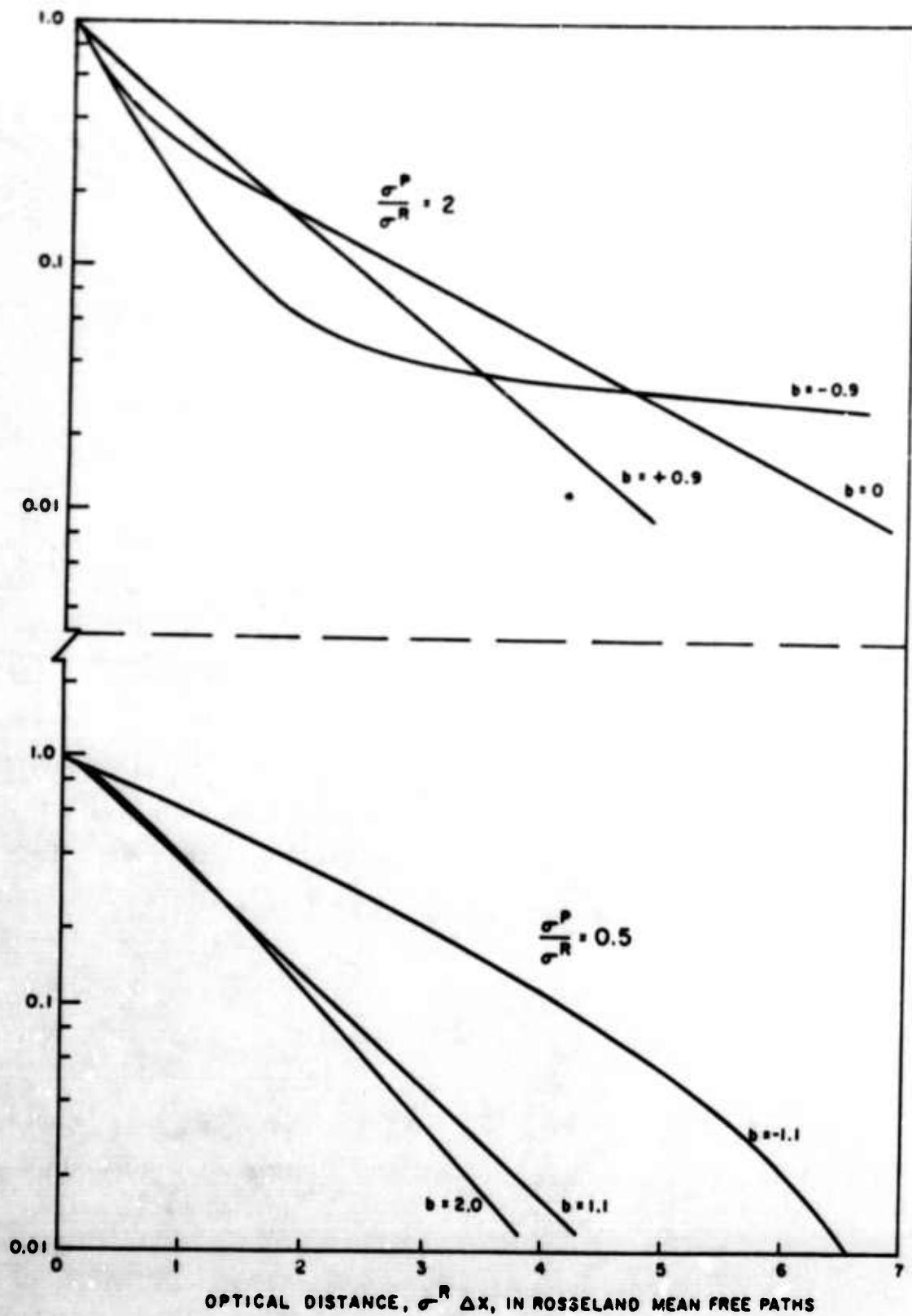


Figure 17. Transmission Function  $S(\Delta x)$

and just half of the emitted radiation attenuates by each of the rates  $\sigma^A$  and  $\sigma^B$ . For  $\sigma^P < \sigma^R$ , in order that the S function defined by Eq. (111) be positive at large distances, b must be restricted to positive values  $b > 1$ . The transmission functions also have a more appropriate shape near the source for  $b > 1$ , as shown in figure 17.

The evaluation of the parameter b will be discussed in a subsequent report, in terms of a distribution function for mean free paths at each point (Ref. 17). In the absence of such information, a "defined" value such as 0 may be used. It should be noted that the Milne-Eddington model is not explicitly assumed; however, b is not treated as a function of position in the definitions of the derivative and integral of the transmission function, Eqs. (112) and (102).

The quantities of the form  $e^{-(\tau_n - \tau_k)}$  and  $\sigma_k$  appearing in Eqs. (92) and (98) may then be evaluated by the expressions given in Eqs. (111) and (112), respectively, to obtain a computationally feasible frequency-averaged form.

The source gradient coefficients  $\alpha_k$ ,  $\beta_k$ ,  $\gamma_k$  in Eqs. (96), (97), and (98) are readily defined by Rosseland optical depth interpolation. The index k may be regarded as ordering the intersections of the ray with successive surfaces, which are alternately boundary surfaces and midsurfaces for each cell; furthermore, k and s are assumed to increase in the direction of the ray. Thus, if k indexes a boundary point and k-1 a zone midpoint, for one-dimensional cases

$$\alpha_k = \frac{B_k - B_{k-1}}{z_k - z_{k-1}}, \quad \beta_k = \frac{B_k - B_{k-1}}{\frac{r_k^2 - r_{k-1}^2}{2}}, \quad \gamma_k = \frac{B_k - B_{k-1}}{\frac{R_k^2 - R_{k-1}^2}{2}} \quad (114)$$

where  $r^2$  and  $R^2$  are, respectively,  $x^2 + y^2$  and  $x^2 + y^2 + z^2$ , and

$$B_k^- = B_{k-1} + \epsilon_k \frac{(B_{k+1} - B_{k-1}) \sigma_k^R (z_k - z_{k-1})}{\sigma_k^R (z_k - z_{k-1}) + \sigma_{k+1}^R (z_{k+1} - z_k)} \quad (115)$$

for plane geometry, and

$$B_k^- = B_{k-1} + \epsilon_k \frac{(B_{k+1} - B_{k-1}) (\sigma_k^R)^2 (R_k^2 - R_{k-1}^2)}{(\sigma_k^R)^2 (R_k^2 - R_{k-1}^2) + (\sigma_{k+1}^R)^2 (R_{k+1}^2 - R_k^2)} \quad (116)$$

for spherical geometry. One-dimensional cylindrical geometry may employ an interpolation formula like Eq. (116) with  $R$  replaced by  $r$ . The coefficient  $\epsilon_k$  is 1 unless the step model is to be used, in which case  $\epsilon_k = 0$  and consequently  $\alpha_k = \rho_k = \gamma_k = 0$ . In two-dimensional cylindrical geometry, the corresponding formulae for bilinear Rosseland optical depth interpolation in the variables  $z$  and  $r^2$  are left as an exercise for the reader. If  $k$  indexes a midpoint, similar formulae apply, with

$$\alpha_k = \frac{B_k - B_{k-1}^+}{z_k - z_{k-1}}, \text{ etc.} \quad (117)$$

A final remark is that for the plane case, the integration over ray orientation should be performed analytically. The exponentials in Eq. (111) are then to be replaced by E-functions. The detailed evaluation of this approach is left for a subsequent report.

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**UNCLASSIFIED**

Security Classification

**DOCUMENT CONTROL DATA - R & D**

(Security classification of title, body of abstract and indexing annotation must be entered when the overall report is classified)

1. ORIGINATING ACTIVITY (Corporate author) Gulf General Atomic Incorporated San Diego, California 92112		2a. REPORT SECURITY CLASSIFICATION <b>UNCLASSIFIED</b>
		2b. GROUP
3. REPORT TITLE <b>NUCLEAR EXPLOSION INTERACTION STUDIES, Vol III, The OUTPUT Code</b>		
4. DESCRIPTIVE NOTES (Type of report and inclusive dates) <b>29 September 1966-27 October 1967</b>		
5. AUTHOR(S) (First name, middle initial, last name) <b>J. R. Triplett et al.</b>		
6. REPORT DATE <b>April 1968</b>	7a. TOTAL NO. OF PAGES <b>206</b>	7b. NO. OF REFS <b>17</b>
8a. CONTRACT OR GRANT NO. <b>F29601-67-C-0014</b>	8b. ORIGINATOR'S REPORT NUMBER(S) <b>AFWL-TR-67-131, Vol III</b>	
8c. PROJECT NO. <b>5710</b>	8d. OTHER REPORT NO(S) (Any other numbers that may be assigned this report) <b>Contractor's report No. GA-7764, Vol III</b>	
c. Subtask No.      07.017 ARPA Order No.      313		
10. DISTRIBUTION STATEMENT This document is subject to special export controls and each transmittal to foreign governments or foreign nationals may be made only with prior approval of AFWL (WLRT), Kirtland AFB, NM, 87117. Distribution is limited because of the technology discussed in the report.		
11. SUPPLEMENTARY NOTES	12. SPONSORING MILITARY ACTIVITY <b>AFWL (WLRT) Kirtland AFB, NM 87117</b>	

13. ABSTRACT

(Distribution Limitation Statement No. 2)

The OUTPUT code is designed for the analysis of early-time nuclear explosions. The equations for radiative transfer (characteristic method) and conservation of total (fluid and radiation) momentum and energy are solved in one-dimensional (plane or spherical) geometry. The radiation equations include first-order Compton scattering, and the hydrodynamic equations are treated in explicit Lagrangian form. The code is undergoing continuing development; the formulation, flow charts, glossary, and listings presented represent its status as of 27 October 1967.

DD FORM NOV 68 1473

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14 KEY WORDS	LINK A		LINK B		LINK C	
	ROLE	WT	ROLE	WT	ROLE	WT
Radiative hydrodynamics codes Compton scattering Early-time nuclear explosions						

APSC (KAFB NM)

**UNCLASSIFIED**

Security Classification